

For TA MSD we find:

$$\begin{aligned}\langle \delta^2(\Delta; t_a, T) \rangle &= \frac{1}{T-\Delta} \int_{t_a}^{t_a+T-\Delta} (x(t+\Delta) - x(t))^2 dt \\ &= \frac{\Delta_\alpha(t_a/T)}{\Gamma(1+\alpha)} \frac{g(\Delta)}{T^{1-\alpha}}\end{aligned}$$

$g(\Delta) = \Delta$ for free motion, $g(\Delta) \approx \Delta^{1-\alpha}$ under confinement

universal prefactor $\Delta_\alpha(z) = (1+z)^\alpha - z^\alpha$ for entire range of Δ !

"Death of linear response":

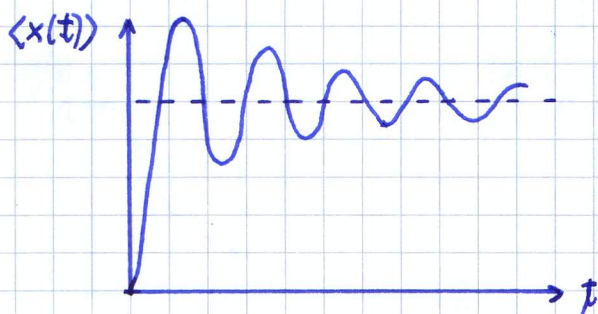
In presence of a time-dependent external force $F(t)$, the fractional Fokker-Planck equation becomes (Solomon & Klafter, PRL 97, 140602 (2006)):

$$\frac{\partial}{\partial t} P(x, t) = \left(-\frac{F(t)}{m\eta_\alpha} \frac{\partial}{\partial x} + K_\alpha \frac{\partial^2}{\partial x^2} \right) \mathcal{D}_t^{1-\alpha} P(x, t).$$

Consider an harmonic forcing of the form $F(t) = F_0 \sin \omega t$.

The first moment becomes

$$\langle x(t) \rangle = \frac{F_0}{m\eta_\alpha} \sqrt{2\pi} S\left(\sqrt{\frac{2t}{\pi}}\right) \text{ in terms of the Fresnel integral } S(z) = \int_0^z \sin\left(\frac{\pi}{2} t^2\right) dt$$



Thus the particle experiences an initial shift. Due to longer & longer waiting times, it shows oscillations of decreasing amplitude.

VII. FIRST PASSAGE PHENOMENA.

In many cases it is of interest when a stochastic variable crosses a given amount for the first time. Examples include chemical reactions, stock market dynamics, or search for a target by animals.

We determine the PDF of first arrival, $g(t)$, from the following boundary condition: $P(x_t, t) = 0$, where x_t is the threshold value.

Consider a normal diffusion process starting @ $x_0 > 0$ with $x_t = 0$: $P_0(x) = \delta(x - x_0)$.

In Laplace space the diffusion equation is

$$uP(x, u) - \delta(x - x_0) = K \frac{\partial^2}{\partial x^2} P(x, u)$$

$$\Rightarrow P(x, u) = c_1 e^{+x\sqrt{u/K}} + c_2 e^{-x\sqrt{u/K}} + \frac{\delta(x - x_0)}{2\sqrt{Ku}} \left(e^{+(x-x_0)\sqrt{u/K}} + e^{-(x-x_0)\sqrt{u/K}} \right)$$

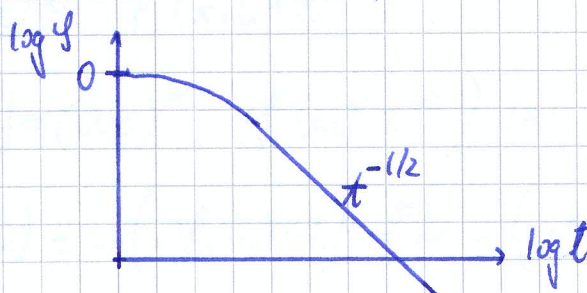
$$P(x=0, u) = 0 \Rightarrow c_1 = -c_2$$

$$\lim_{x \rightarrow \infty} P(x, u) = 0 \Rightarrow c_1 e^{x\sqrt{u/K}} - \frac{1}{2\sqrt{Ku}} e^{x\sqrt{u/K} - x_0\sqrt{u/K}} = 0$$

$$\Rightarrow c_1 = \frac{1}{2\sqrt{Ku}} e^{-x_0\sqrt{u/K}}$$

The survival probability becomes

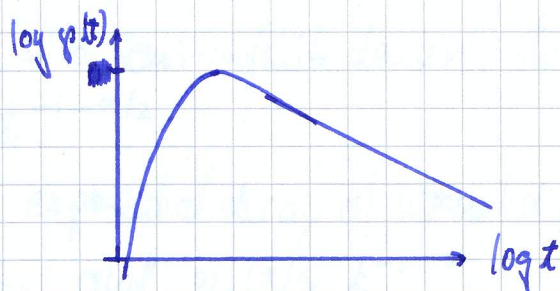
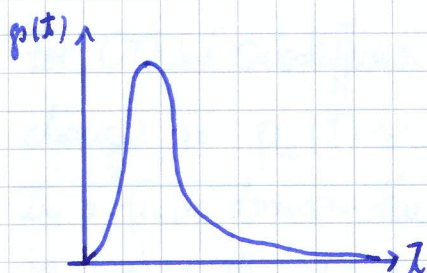
$$S(t) = \int_0^\infty P(x, u) dx = 1 - \operatorname{erfc}\left(\frac{x_0}{\sqrt{Kt}}\right) \sim \frac{x_0}{\sqrt{\pi Kt}}$$



And we obtain the first passage PDF

$$g(t) = -\frac{d}{dt} S(t) = \frac{x_0}{\sqrt{4\pi Kt^3}} e^{-\frac{x_0^2}{4Kt}} \sim \frac{x_0}{\sqrt{4\pi Kt^3}} \quad \text{Lévy-Smirnov PDF}$$

$$\langle t \rangle = \int_0^\infty t g(t) dt = \infty \quad \text{diverging mean first passage time}$$



The same result would immediately follow from the image solution

$Q(x, t) = G(x - x_0, t) - G(x + x_0, t)$ where $G(x, t)$ is the solution of the diffusion equation for initial condition $G(x, 0) = \delta(x)$ (Green's function)

For random walk processes with local jump length distribution the first passage is identical to the first arrival to a point, which we describe by a sink term:

$$\frac{\partial}{\partial t} P(x, t) = \kappa \frac{\partial^2}{\partial x^2} P(x, t) - g_{fa}(t) \delta(x)$$

such that $g_{fa}(t) = p(t)$. The first arrival PDF g_{fa} fulfils the chain rule

$$P(-x_0, t) = \int_0^t g_{fa}(t') P(0, t - t') dt' \Rightarrow P(-x_0, u) = g_{fa}(u) P(0, u)$$

This relation is useful when calculating $g(t)$.

For CTRW subdiffusion we obtain a subordination relation for the first passage PDF $g_\alpha(t)$:

$$\text{We had } P_\alpha(x, u) = (u\tau)^{\alpha-1} P_1(x, u^\alpha \tau^{\alpha-1})$$

$$\text{From integration: } \psi_\alpha(u) = (u\tau)^{\alpha-1} \psi_1(u^\alpha \tau^{\alpha-1}) \quad (*)$$

$$\text{In Laplace space: } g_\alpha(u) = - (u \psi_\alpha(u) - 1) \stackrel{(*)}{=} - (u^\alpha \tau^{\alpha-1} \psi_1(u^\alpha \tau^{\alpha-1}) - 1)$$

With $g_1(u) = - (u \psi_1(u) - 1)$ we see that

$$\underline{g_\alpha(u) = g_1(u^\alpha \tau^{\alpha-1})}$$

$$g_\alpha(t) = \int_0^\infty \tilde{F}_\alpha(s, t) g_1(s) ds \therefore \tilde{F}_\alpha(s, u) = \exp(s u^\alpha \tau^{\alpha-1}) \quad \text{characteristic f.t. of a one-sided Lévy stable density}$$