

Beyond Brownian motion: Long-range correlation

—Pohang, 27th April 2023—



The founding fathers of diffusion theory: $\langle r^2(t) \rangle \propto Kt$



Paul Langevin (1872-1946)

Albert Einstein (1879-1955)



Marian Smoluchowski
(1872-1917)



William Sutherland
(1859-1911)

Einstein's conditions on Brownian motion

Paul Lévy [Processus stochastiques & mouvement brownien (1948)]: «The stochastic process, that we will call linear Brownian motion, is a schematisation that well represents the properties of real Brownian motion, observable on a sufficiently small but not infinitely small scale, and which assumes that the same properties exist across the scales.»

Einstein's postulate: a stochastic process describes normal diffusion if

- (i) $\exists$ finite correlation time beyond which displacements are independent
- (ii) displacements are identically distributed
- (iii) displacements have a finite second moment

Main features: linear mean squared displacement & Gaussian probability density (CLT)

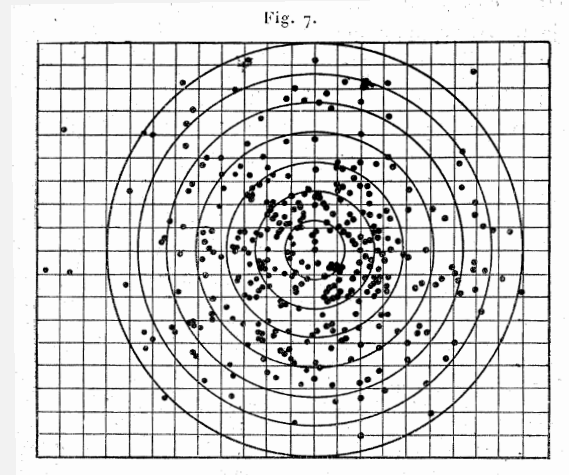
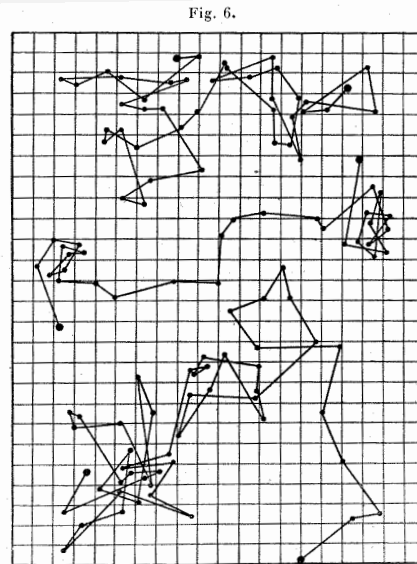
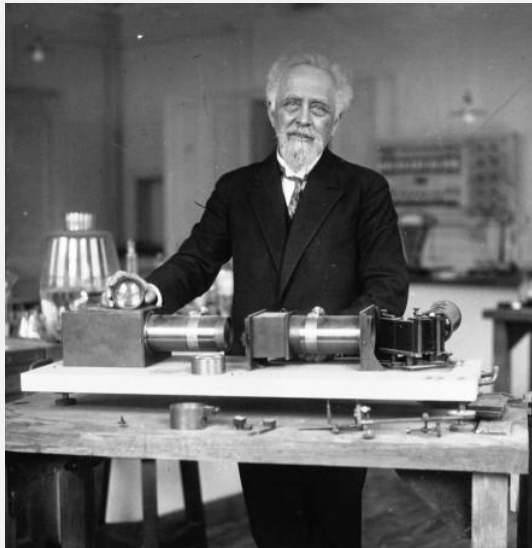
$$\langle \mathbf{r}^2(t) \rangle = 2dK_1t, \quad P(\mathbf{r}, t) = (4\pi K_1t)^{-d/2} \exp\left(-\frac{\mathbf{r}^2}{4K_1t}\right)$$

Confining potential $V(x)$: Boltzmannian equilibrium PDF: $P_{\text{eq}}(x) = \mathcal{N} \exp(-\beta V(x))$

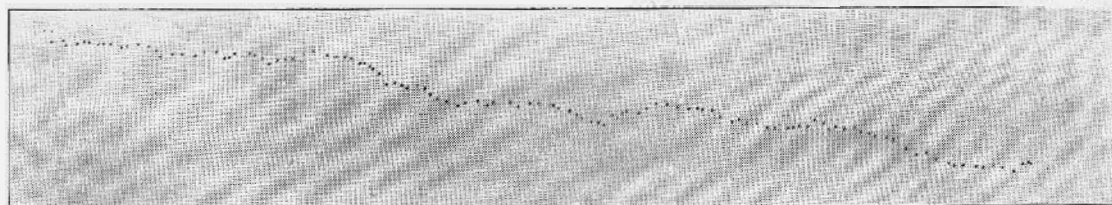
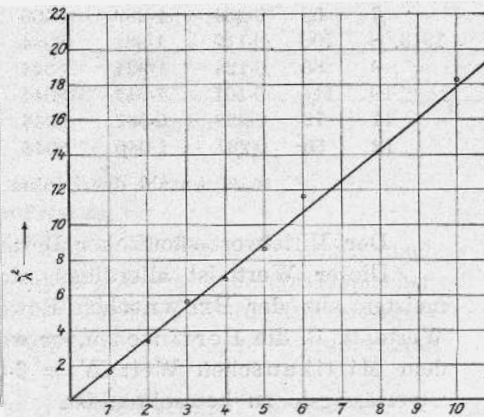
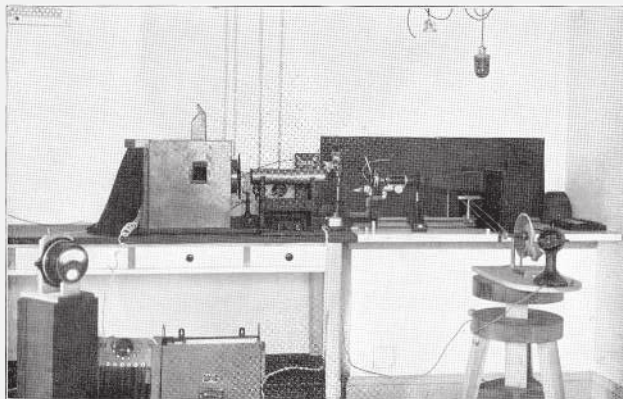
Ergodic in Birkhoff-Boltzmann-Khinchin sense (time average converges to ensemble)

Violation of these conditions leads to anomalous diffusion with MSD $\langle \mathbf{r}^2(t) \rangle \simeq K_\alpha t^\alpha$ and/or non-Gaussian PDF $\leadsto$ Breakdown of universality

Les atomes: Brownian motion and Avogadro's number



$$K = \frac{k_B T}{m \eta} = \frac{(R/N_A) T}{m \eta}$$

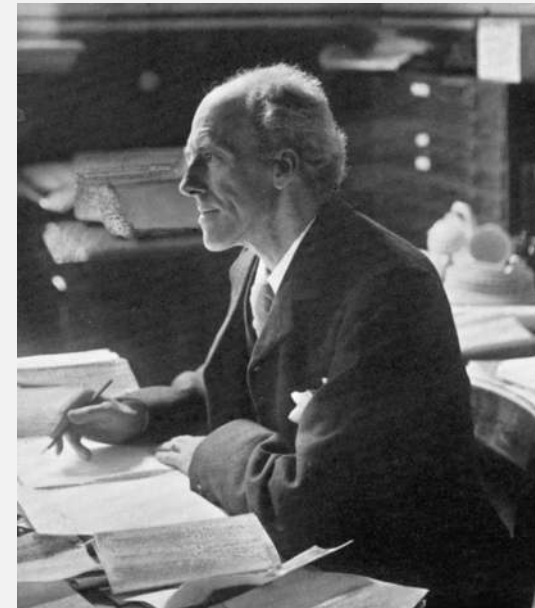


J Perrin, Comptes Rendus (Paris) 146 (1908) 967: $N_A = 70.5 \times 10^{22}$; I Nordlund, Z Physik (1914): $N_A = 5.91 \times 10^{23}$ 4

Brownian motion & random walks applied . . .



Louis Bachelier, 1900:
Brownian motion of
financial markets



Karl Pearson, 1905:
Spreading of mosquitoes
in rain forests

Marian v Smoluchowski, 1916:
Molecular reaction kinetics



Elliott Montroll, 1943:
Manhattan project,
chain reaction of
neutrons

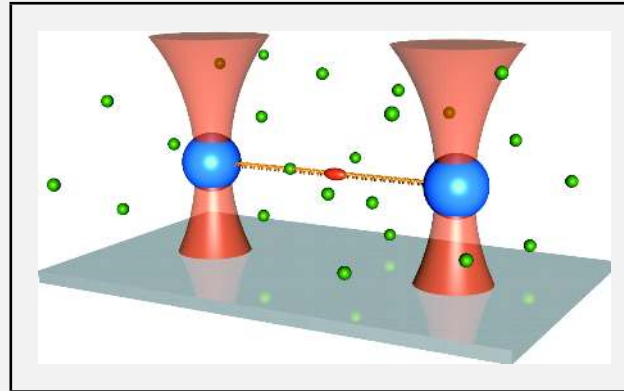
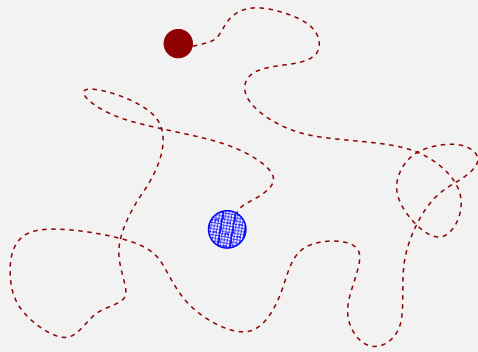


Scanned at the American
Institute of Physics

Molecular reaction times: macro vs micro

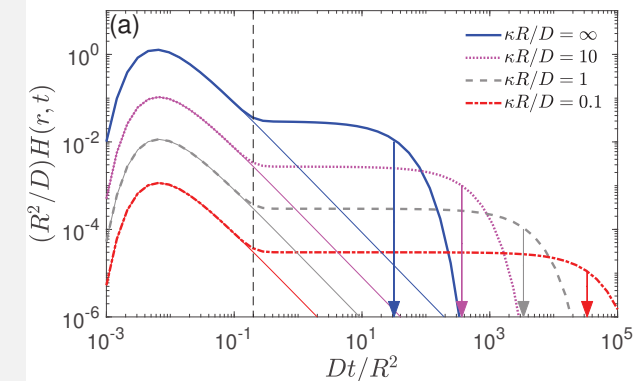
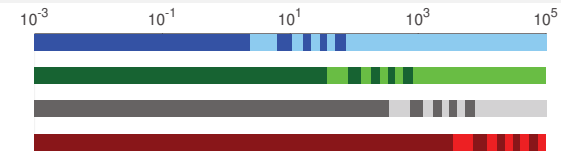
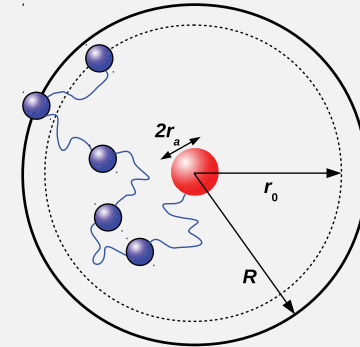
Search rate for particle with diffusivity D_{3d} to find an immobile target of radius a (assuming immediate binding):

$$k_{\text{on}} = 4\pi D_{3d}a$$

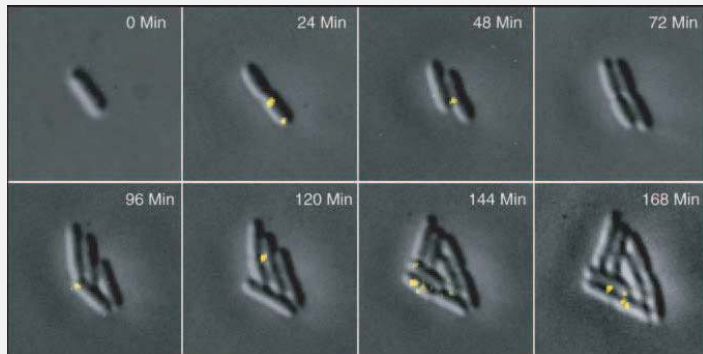


G Wuite, RM et al, PNAS (2008,2009)

Strong defocusing of reaction times
Geometry- vs reaction-control



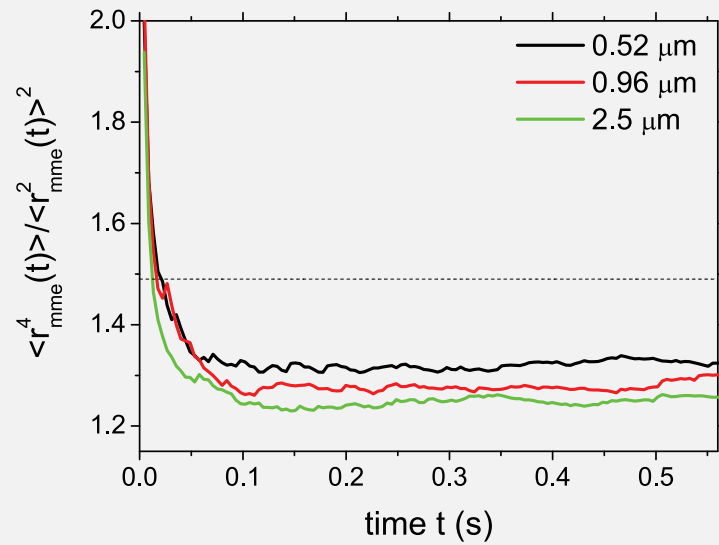
Yu et al, Science (2006)



M v Smoluchowski, Physikal Zeitschr (1916)

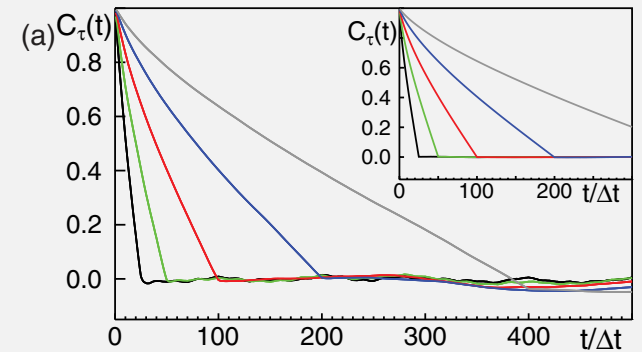
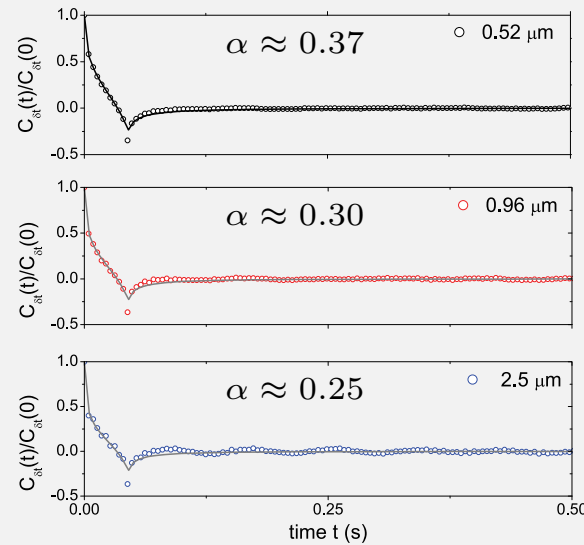
A Godec & RM, PRX (2016); D Grebenkov, RM & G Oshanin (2018-) 6

Passive motion of submicron tracers in cells is viscoelastic I

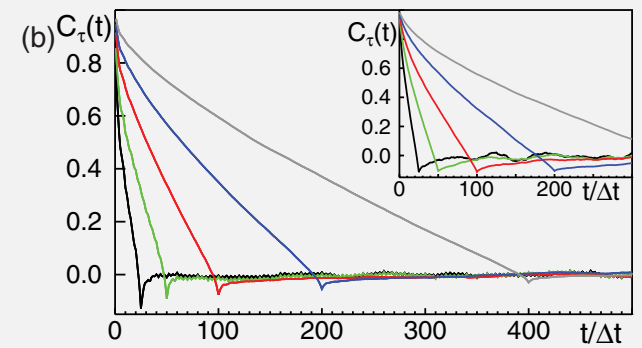


Beads in worm-like micellar solution

JH Jeon, N Leijnse, L Oddershede & RM, NJP (2013)

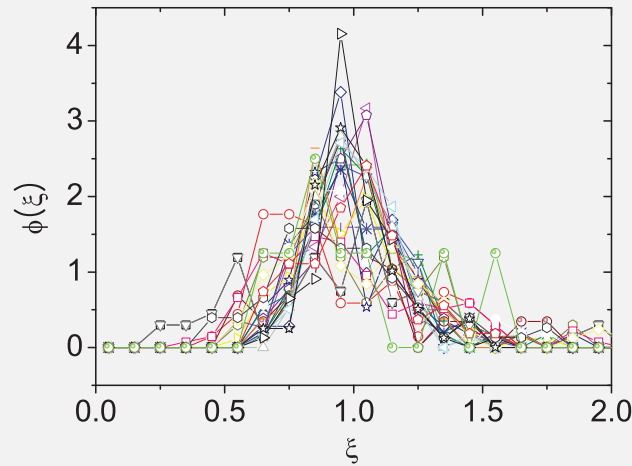
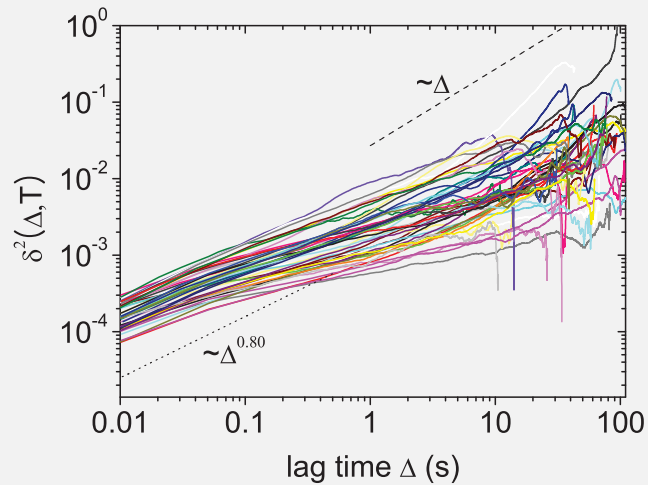


Sucrose solution: $\alpha = 1$



Dextrane solution: $\alpha \approx 0.8$

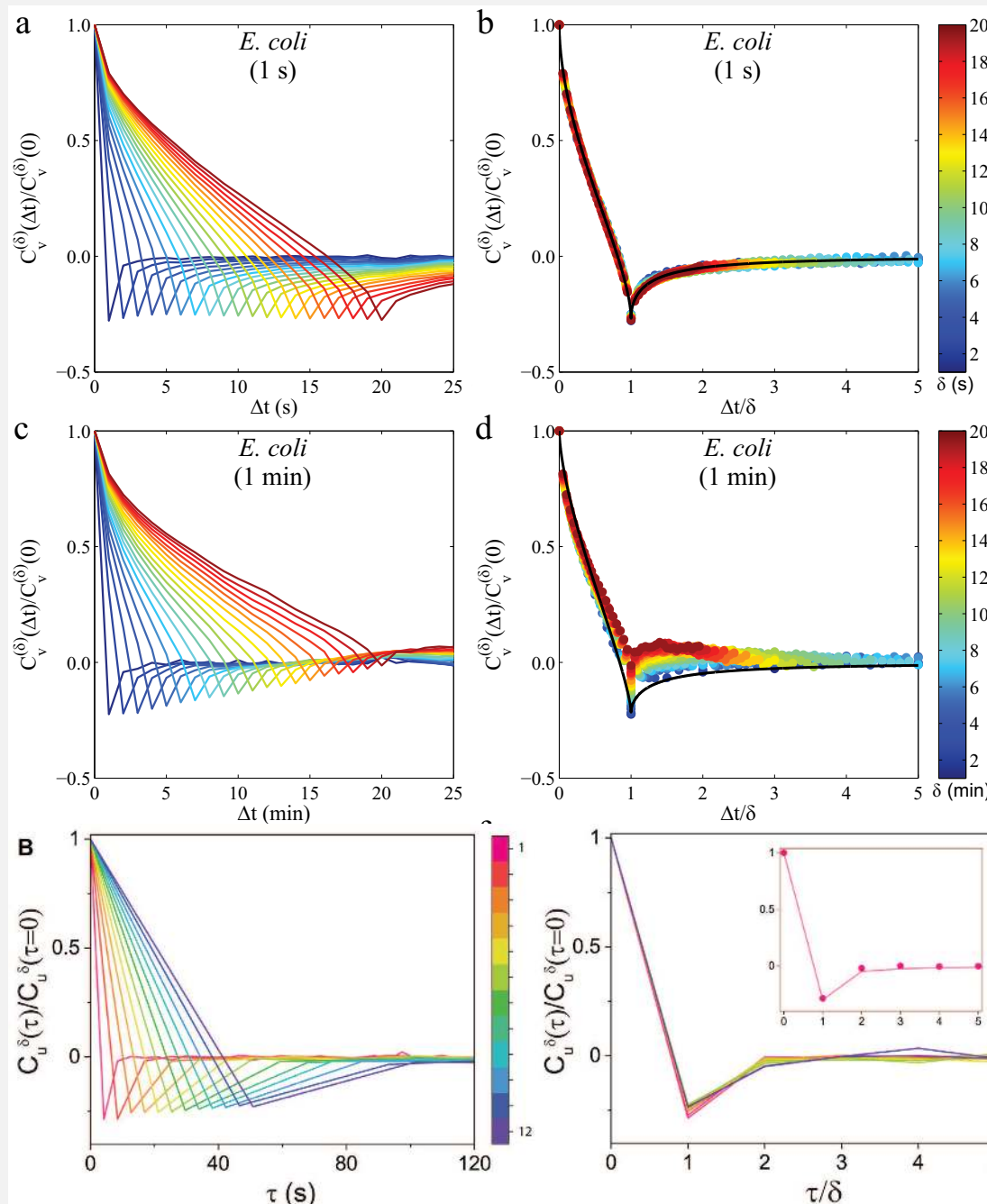
M Weiss, PRE (2013)



Endogeneous lipid granules
in living yeast cells

J-H Jeon, V Tejedor, S Burov, E Barkai, C Selhuber-Unkel, K Berg-Sørensen, L Oddershede & RM, PRL (2011)

Passive motion of submicron tracers in cells is viscoelastic II



RNA molecules in living *E. coli* cells

TJ Lampo, S Stylianidou, MP

Backlund, PA Wiggins & AJ

Spakowitz, BPJ (2017)

AttB DNA-binding proteins in living *Bacillus subtilis* cells

R Arbel-Goren, SA McKeithen-

Mead, D Voglmaier, I Afremov, G

Teza, AD Grossman & J Stavans,

NAR (2023)

Fractional Brownian motion

FLE: fluctuation-dissipation $\leadsto$ asymptotic thermal equilibrium [books by Zwanzig, Kubo]

FBM: “external noise” for non-equilibrium systems or “open systems” [Klimontovich, Statistical physics of open systems] $\leadsto$ biological cells, animals, weather, finance

Mandelbrot-van Ness smoothed FBM [SIAM Rev (1968)]:

$$\frac{dx(t)}{dt} = \sqrt{2D(t)}\xi_H(t)$$

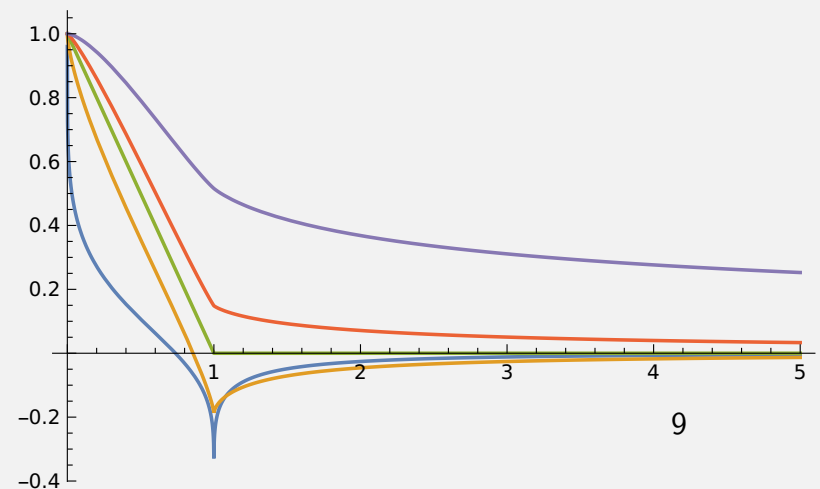
$\xi_H(t)$ is fractional Gaussian noise, understood as the derivative of smoothed FBM:

$$\langle \xi_H(t)\xi_H(t+\tau) \rangle = \frac{1}{2\delta} \left(|t+\delta|^{2H} - 2|\tau|^{2H} + |t-\delta|^{2H} \right) \sim H(2H-1)\tau^{2H-2}$$

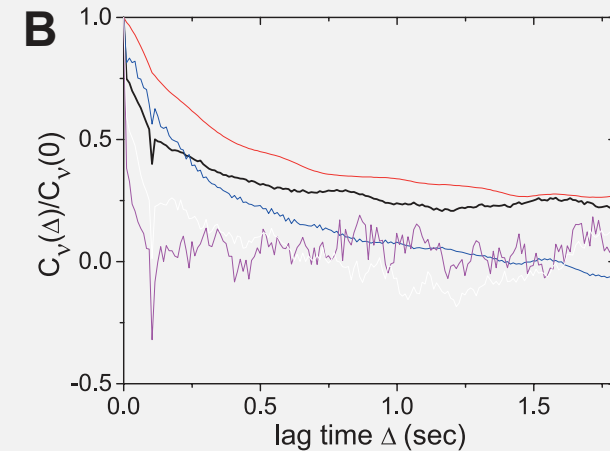
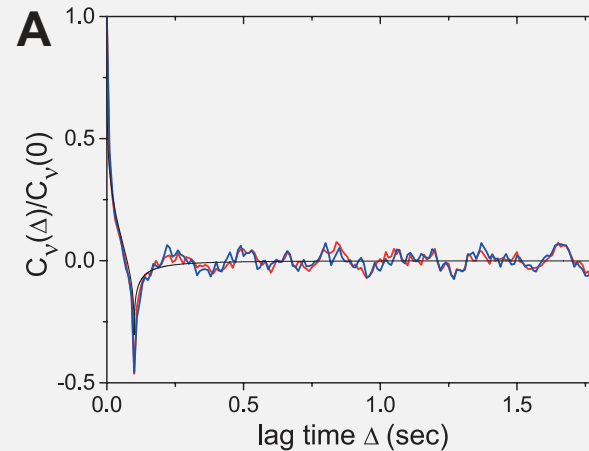
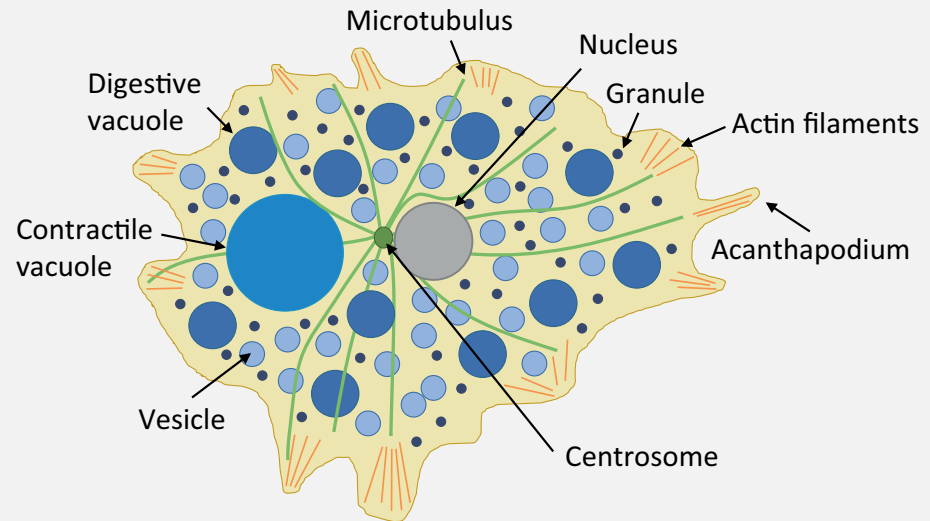
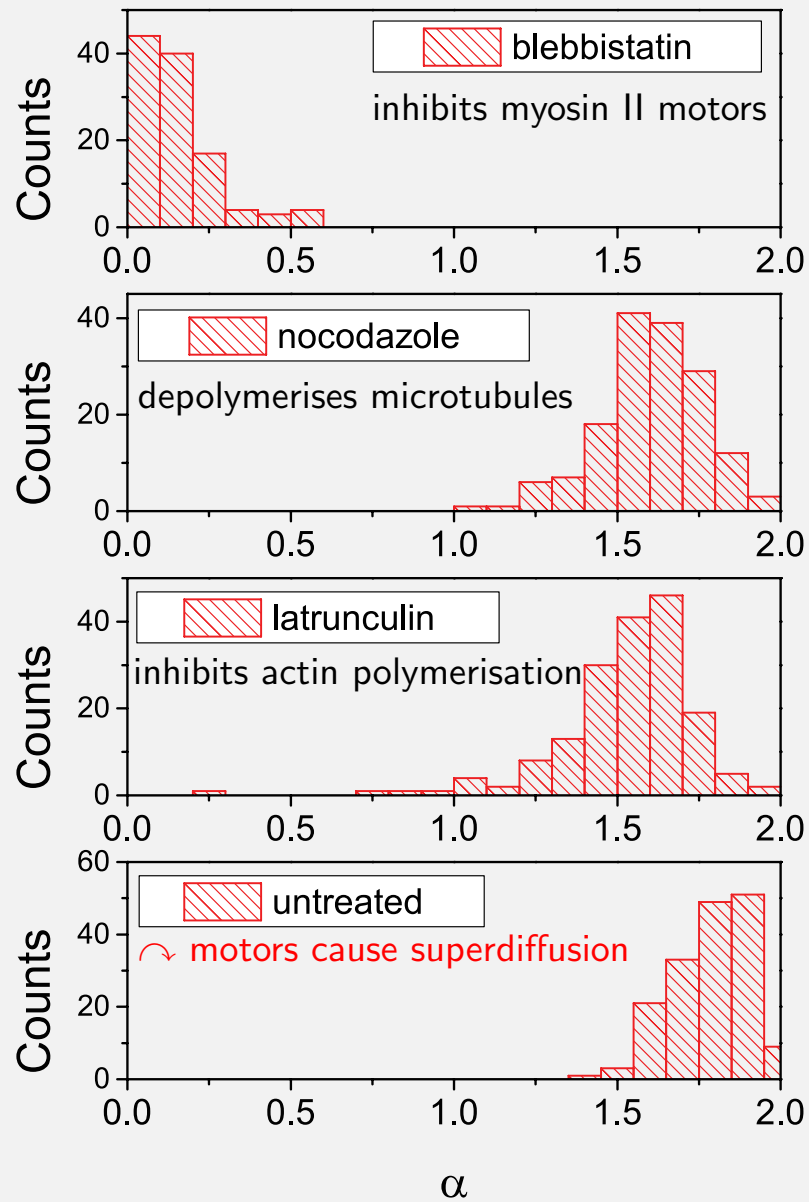
Displacement correlator:

$$C_{\delta t}(t) = \frac{\langle [\mathbf{r}(t+\delta t) - \mathbf{r}(t)] \cdot \mathbf{r}(\delta t) - \mathbf{r}(0) \rangle}{\delta t^2}$$

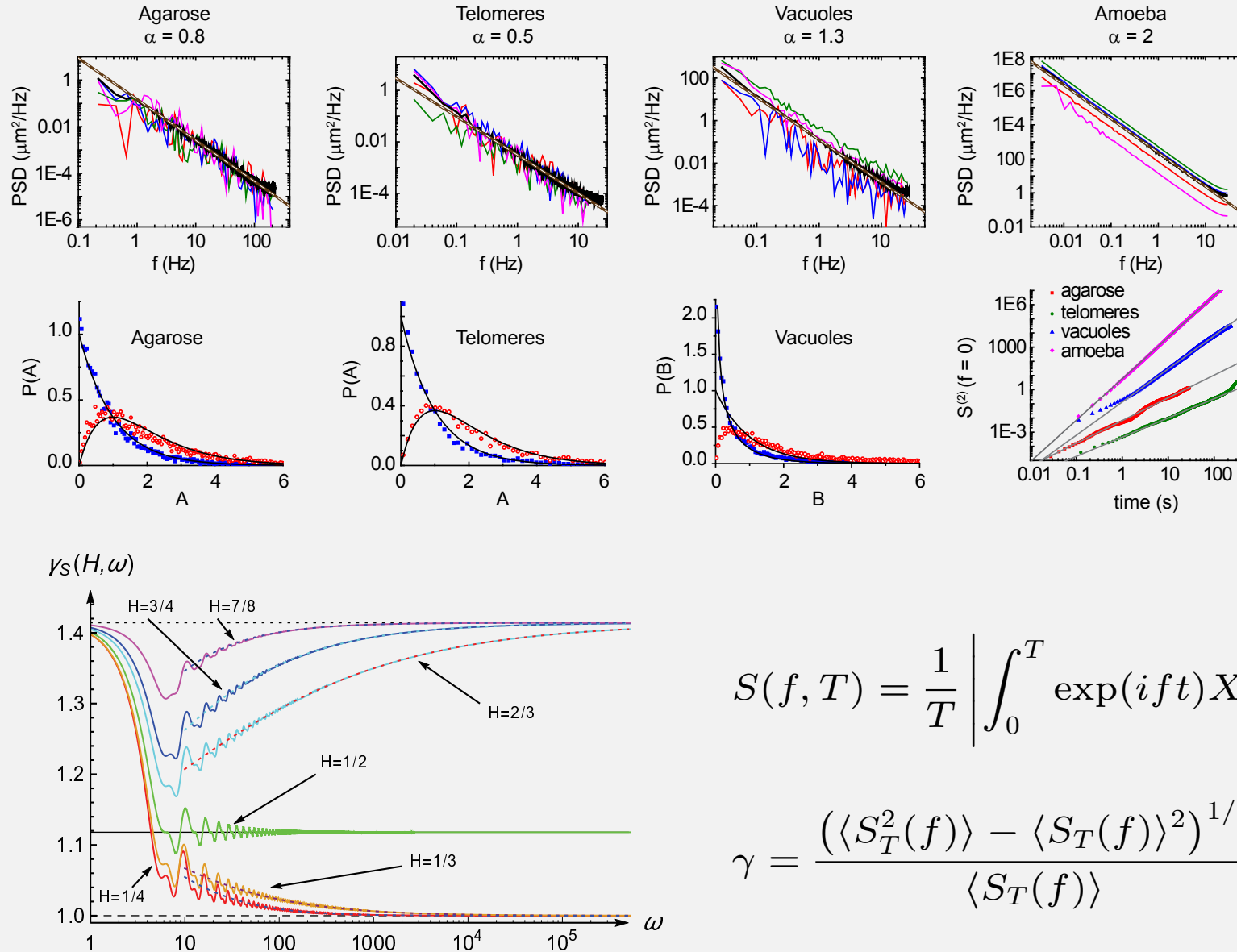
$$\frac{C_{\delta t}(t)}{C_{\delta t}(0)} = \frac{(t+\delta t)^{2H} - 2t^{2H} + |t-\delta t|^{2H}}{2\delta t^{2H}}$$



Superdiffusion in supercrowded *Acanthamoeba castellani*



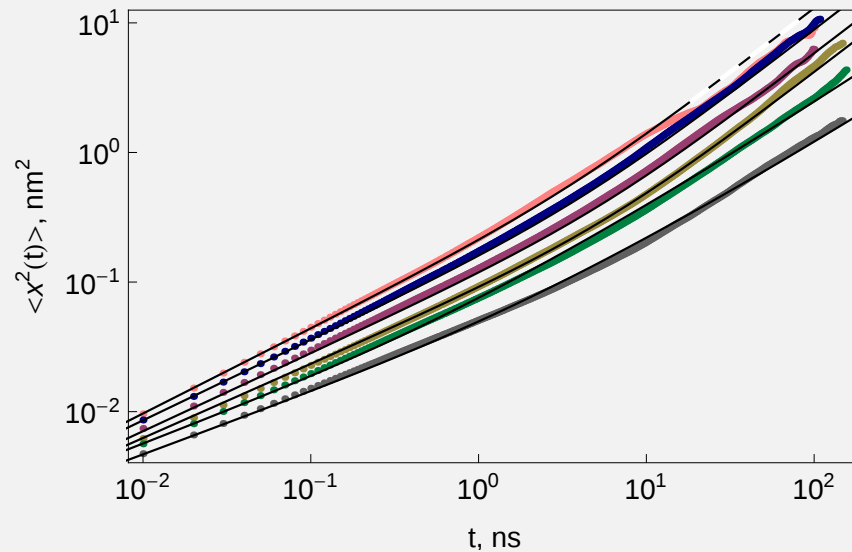
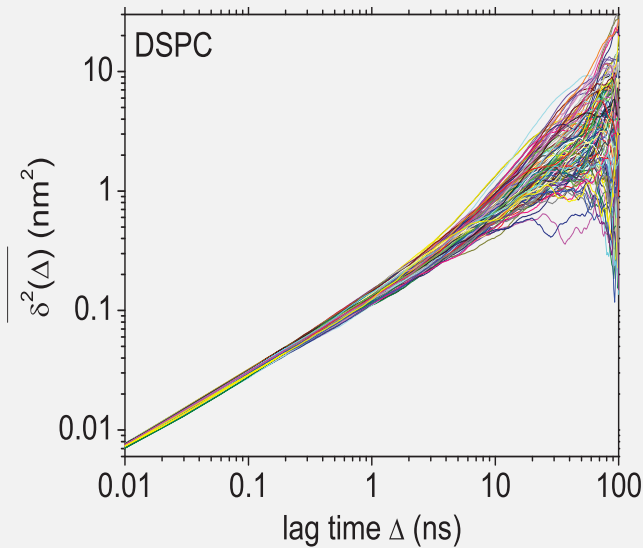
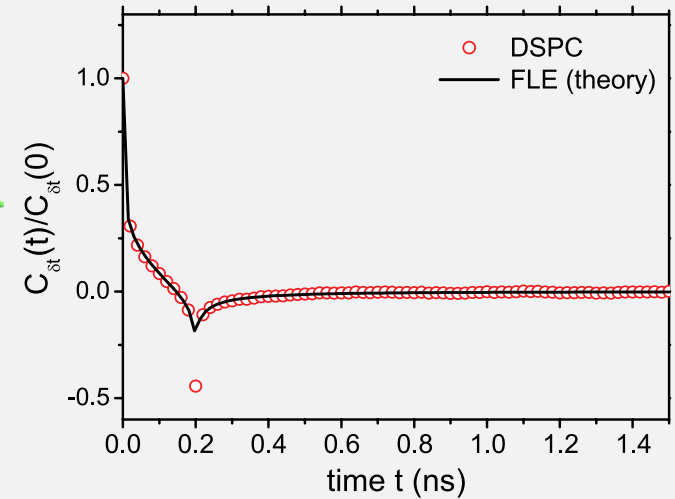
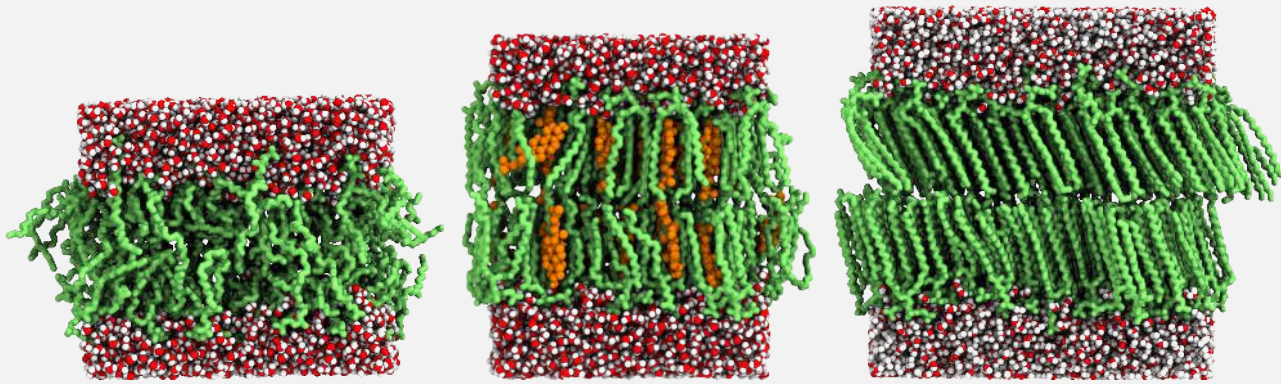
Power spectral density of a single FBM trajectory



$$S(f, T) = \frac{1}{T} \left| \int_0^T \exp(i f t) X_t dt \right|^2$$

$$\gamma = \frac{(\langle S_T^2(f) \rangle - \langle S_T(f) \rangle^2)^{1/2}}{\langle S_T(f) \rangle}$$

Single lipid motion in bilayer membranes



Pure bilayer:

$$t^{2/3} \longrightarrow t^1$$

/w cholesterol/gel:

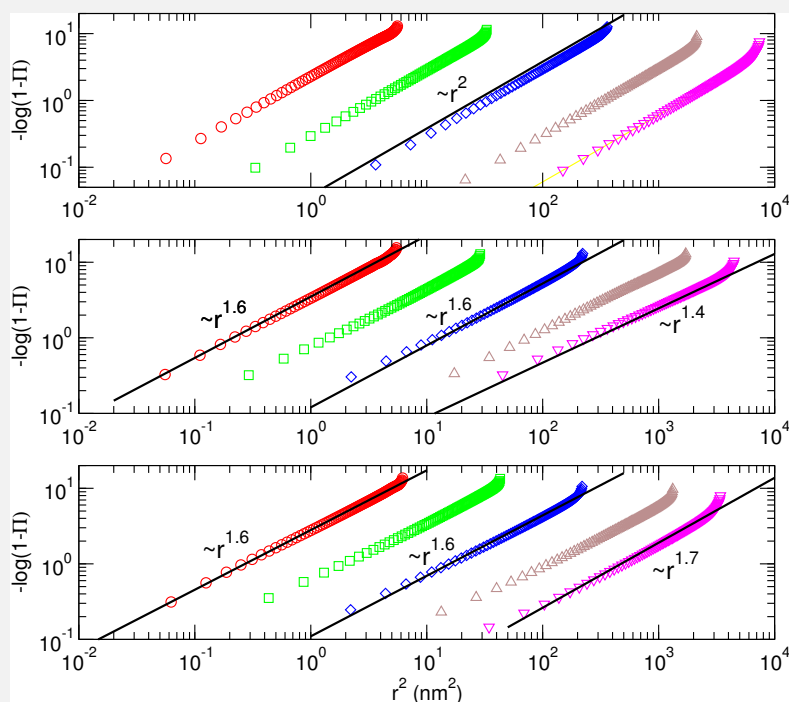
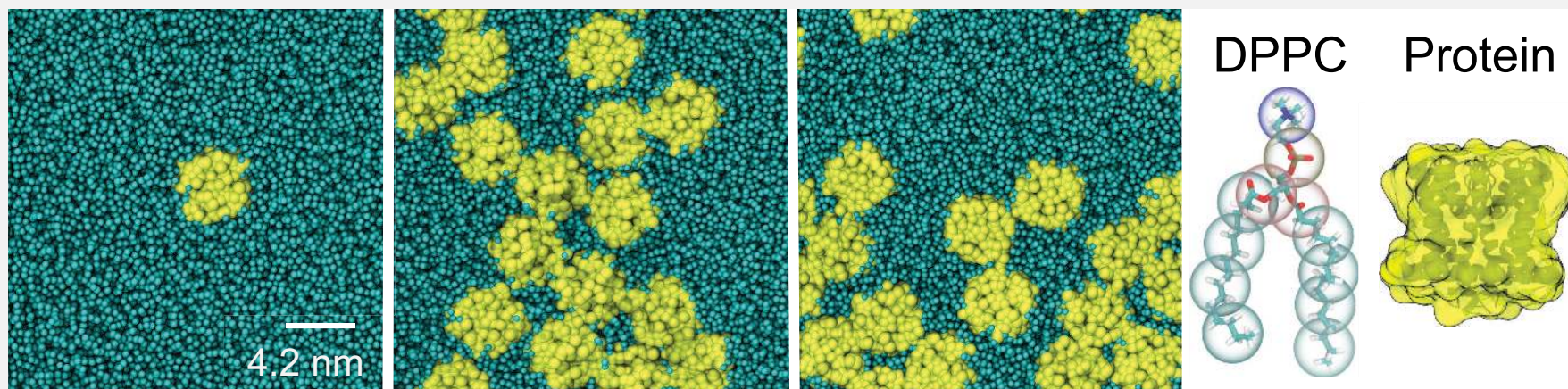
$$t^\alpha \longrightarrow t^{\alpha' > \alpha}$$

Tempered GLE

D Molina-Garcia, T Sandev, H Safdari, G Pagnini, AV Chechkin & RM, NJP (2018)

J-H Jeon, H Martinez-Seara Monne, M Javanainen & RM, PRL (2012)

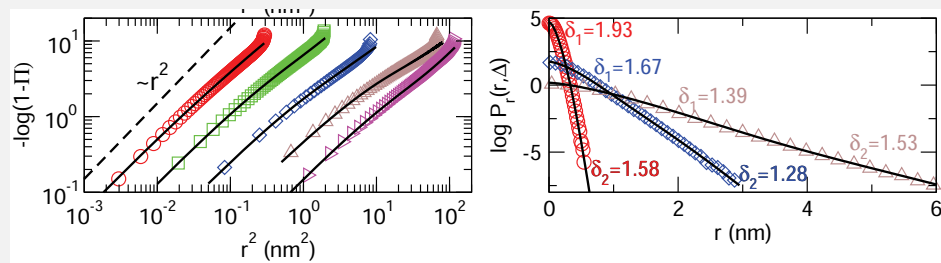
Crowding in membranes: non-Gaussian lipid/protein diffusion



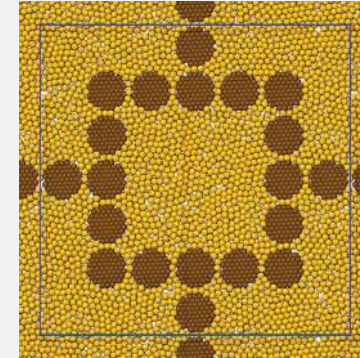
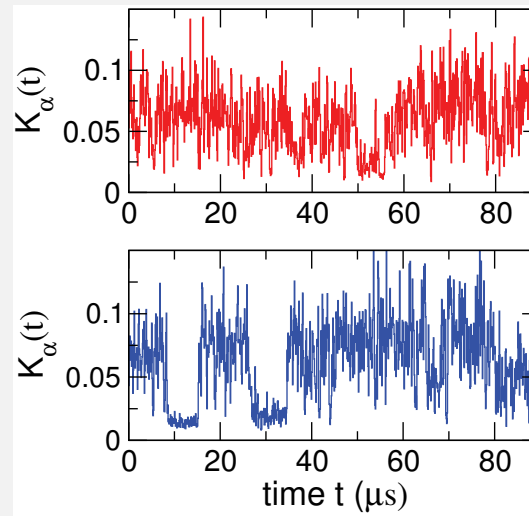
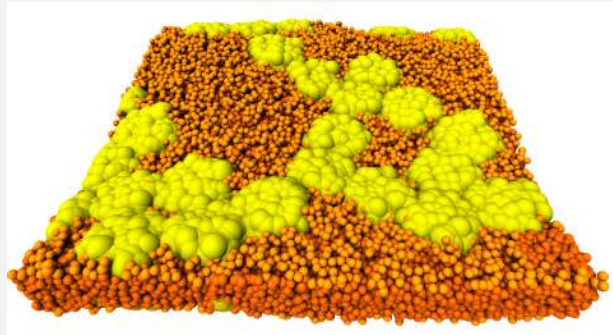
Dilute membrane: $P(r, t)$ Gauss

Crowded membrane ($\delta \approx 1.3 \dots 1.7$):

$$P(r, t) \propto \exp \left(- \left[\frac{r}{ct^{\alpha/2}} \right]^{\delta} \right)$$



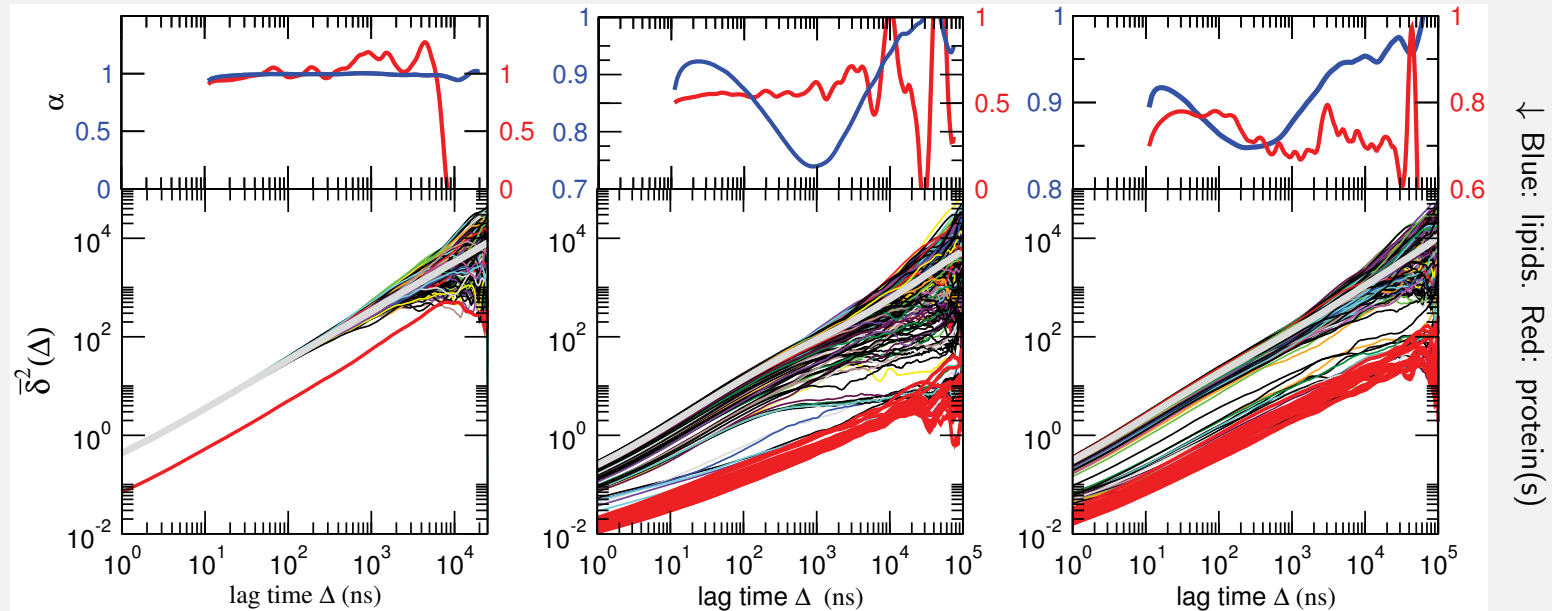
Intermittent lipid diffusion in protein-crowded membranes



Single NaK channel

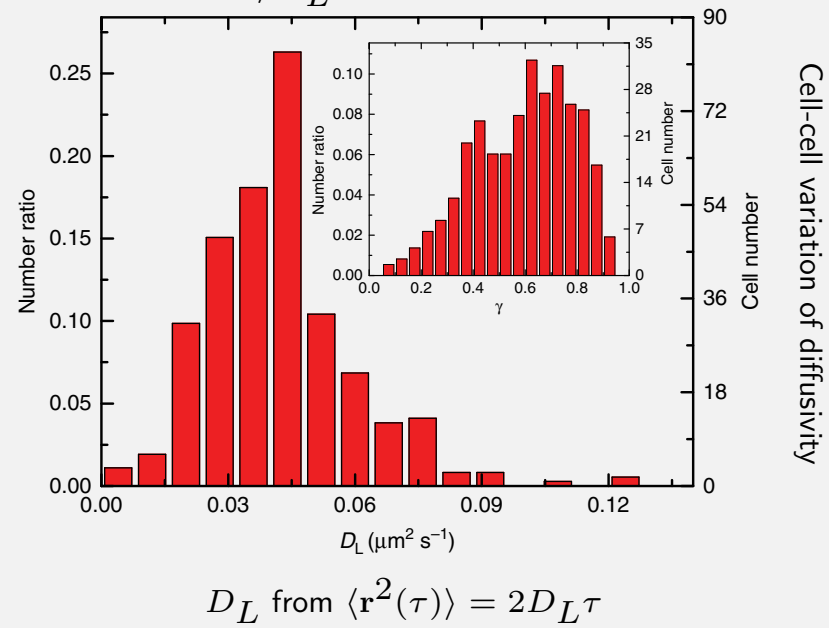
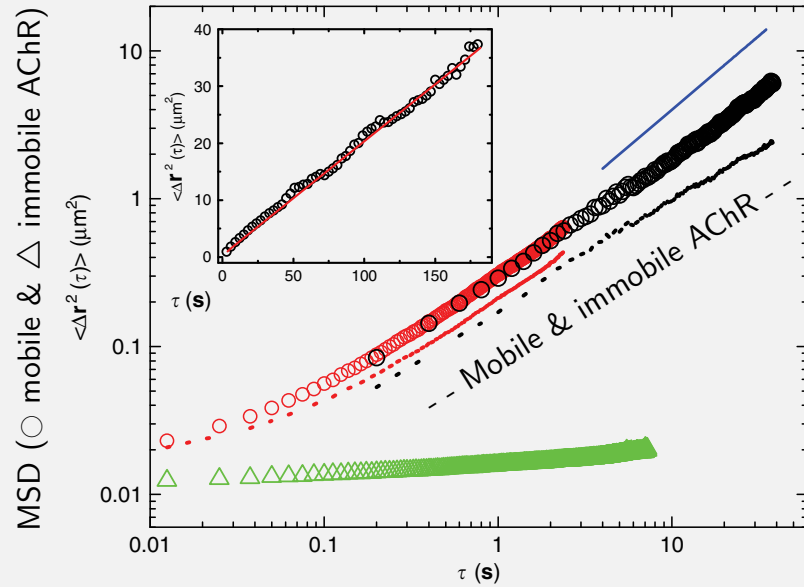
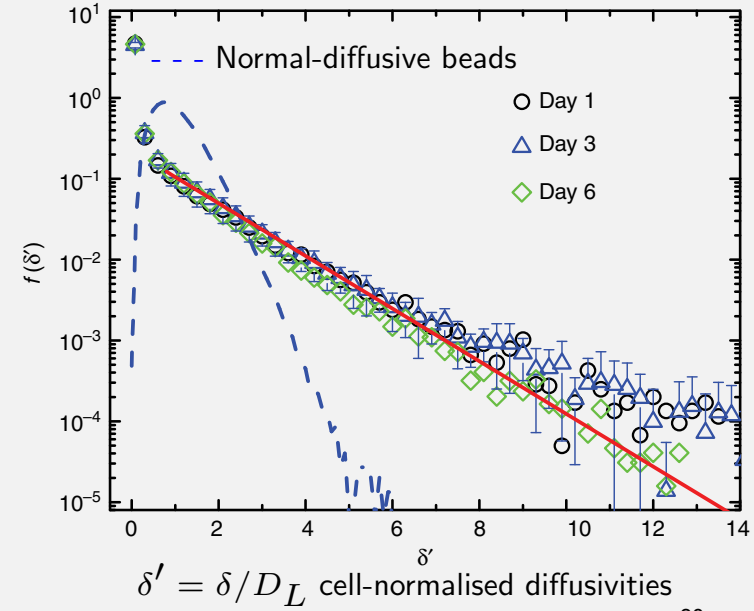
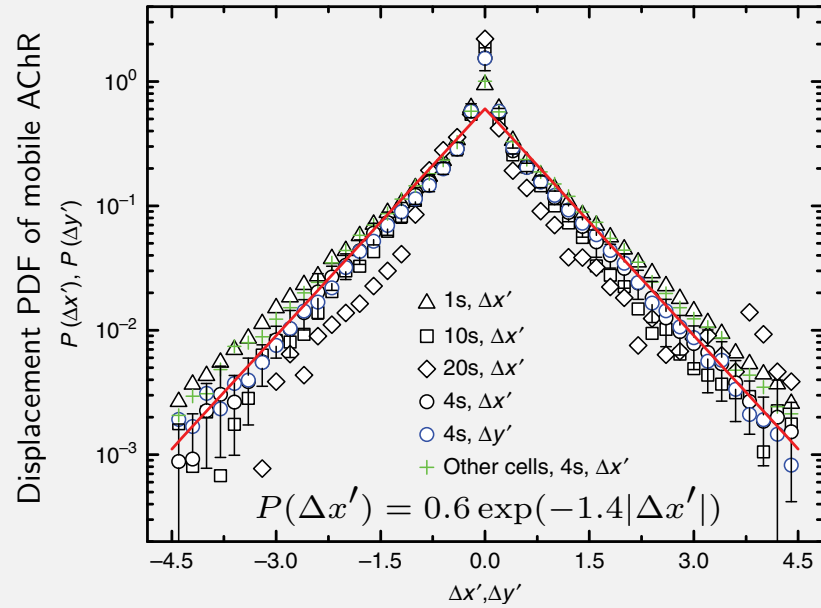
DLPC (non-aggregating)

DPPC (aggregating)

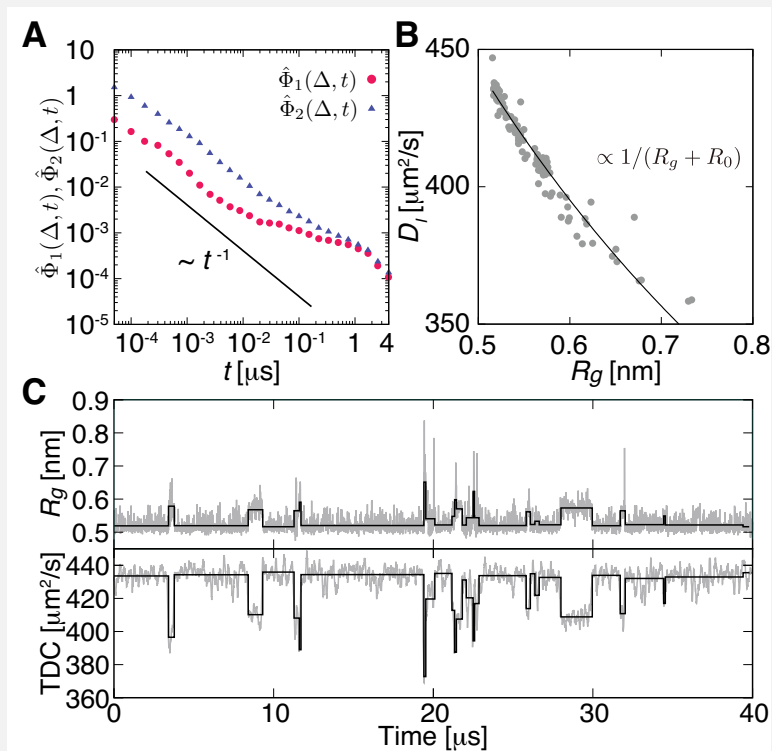
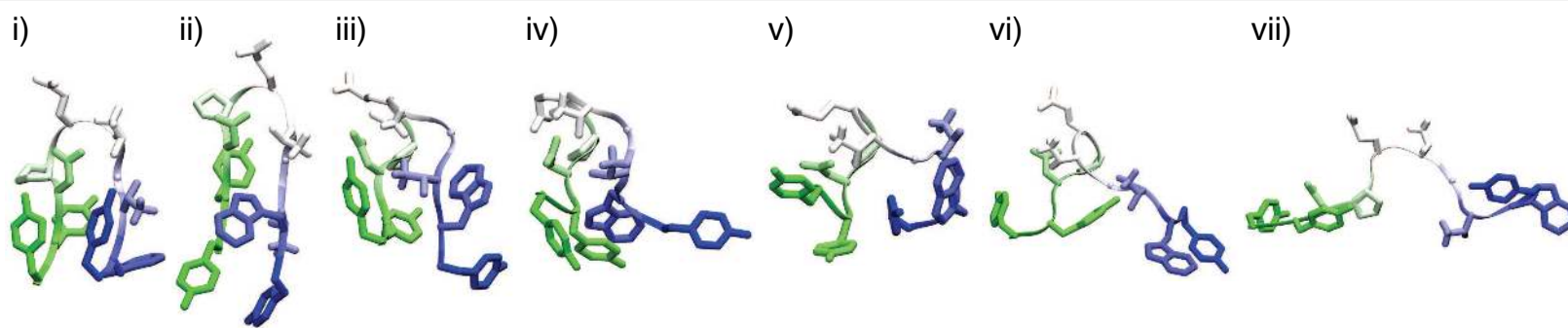


J-H Jeon, M Javanainen, H Martinez-Seara, RM & I Vattulainen, PRX (2016); compare Faraday Disc (2013)

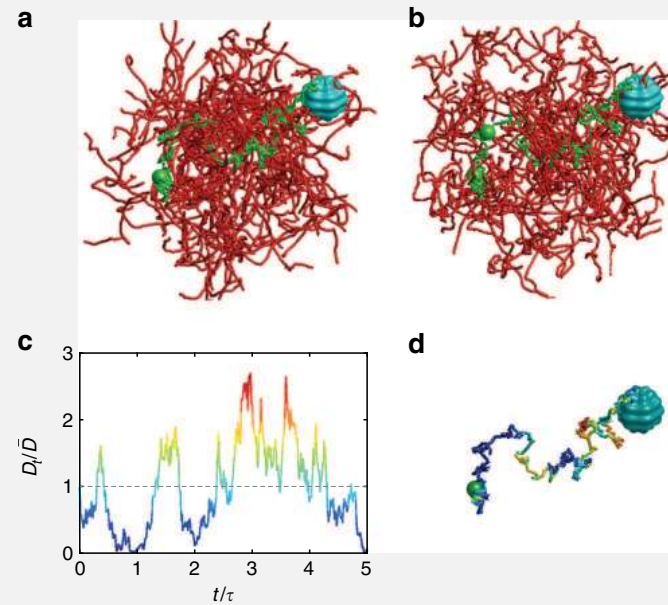
Non-Gaussianity of acetylcholine receptors in *Xenopus* cells



Instantaneous diffusivity of shape-fluctuating proteins



Annealed diffusing-diffusivity in rearranging heterogeneous environments



Y Lanoiselée, N Moutal,
& D Grebenkov, Nat Comm (2018)

E Yamamoto, T Akimoto, A Mitsutake & RM, PRL (2021)

Viscoelastic diffusing-diffusivity model

DD model for Brownian yet non-Gaussian motion:

MV Chubinsky & G Slater, PRL (2014); R Jain & KL Sebastian, JPC B (2016), AV Chechkin, F Seno, RM & IM Sokolov, PRX (2017); V Sposini, AV Chechkin, G Pagnini, F Seno & RM, NJP (2018)

FBM-generalised diffusing-diffusivity model:

$$\frac{dx(t)}{dt} = \sqrt{2D(t)}\xi_H(t)$$

with $D(t)$ as squared Ornstein-Uhlenbeck process $Y(t)$:

$$D(t) = Y^2(t), \quad \frac{dY(t)}{dt} = -Y + \eta(t)$$

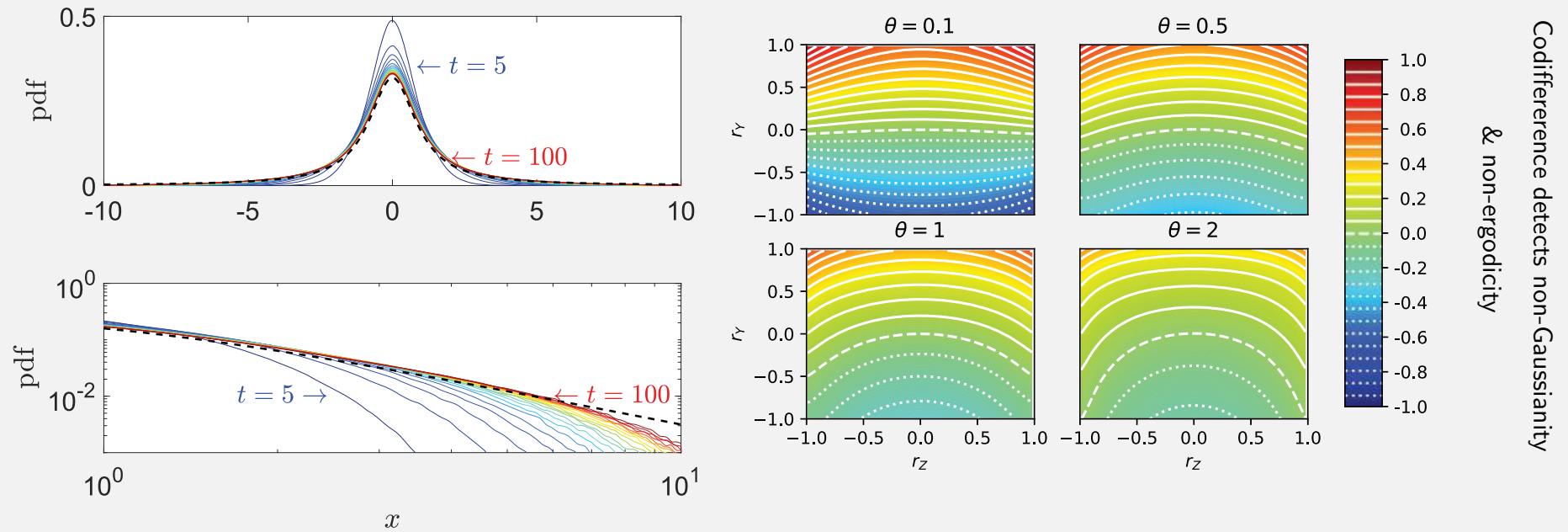
Here, $\eta(t)$ is white Gaussian noise with zero mean & unit variance

$\xi_H(t)$ is fractional Gaussian noise, understood as the derivative of the Mandelbrot-van Ness smoothed FBM [SIAM Rev (1968)]:

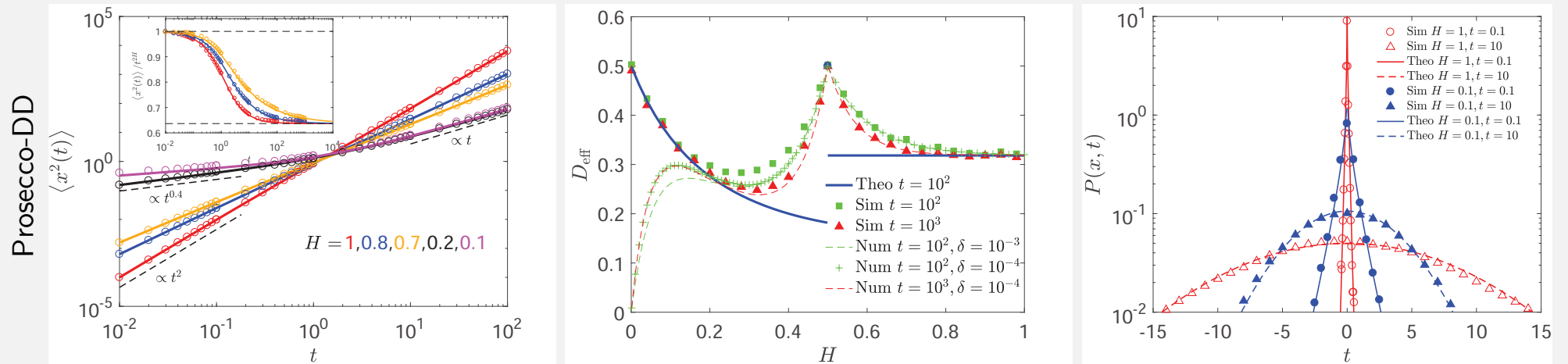
$$\langle \xi_H(t)\xi_H(t + \tau) \rangle = \frac{1}{2\delta} \left(|t + \delta|^{2H} - 2|\tau|^{2H} + |t - \delta|^{2H} \right) \sim H(2H - 1)\tau^{2H-2}$$

W Wang, AV Chechkin, F Seno, IM Sokolov & RM, NJP (2020); W Wang, AG Cherstvy, AV Chechkin, S Thapa, F Seno, X Liu & RM, JPA (2020)

Non-Gaussian dynamics in the presence of correlated noise

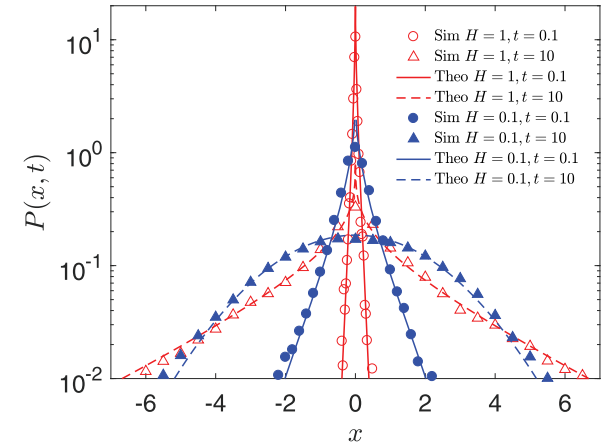
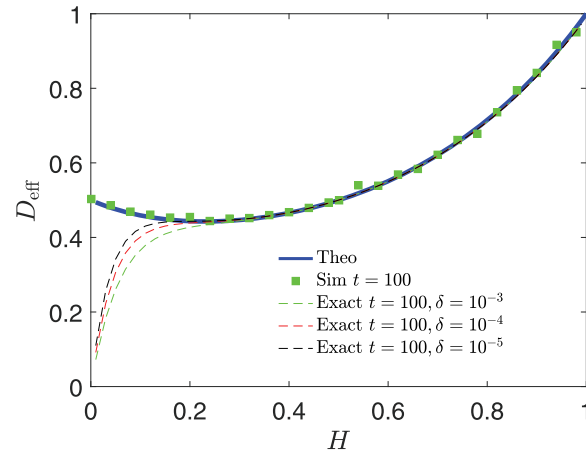
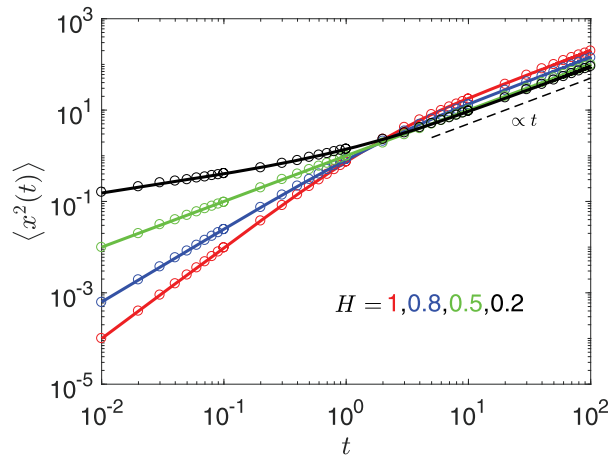


Viscoelastic diffusing-diffusivity model

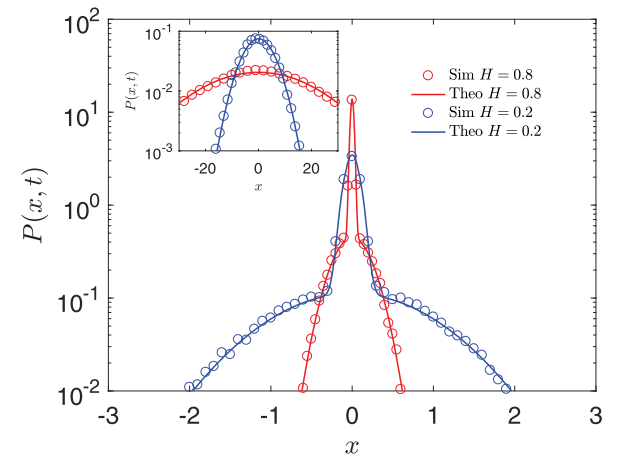
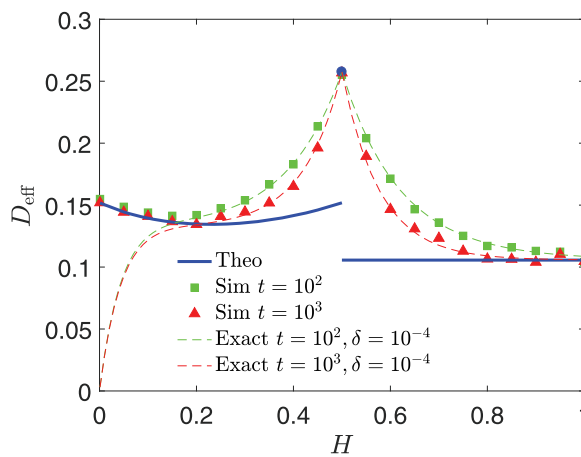
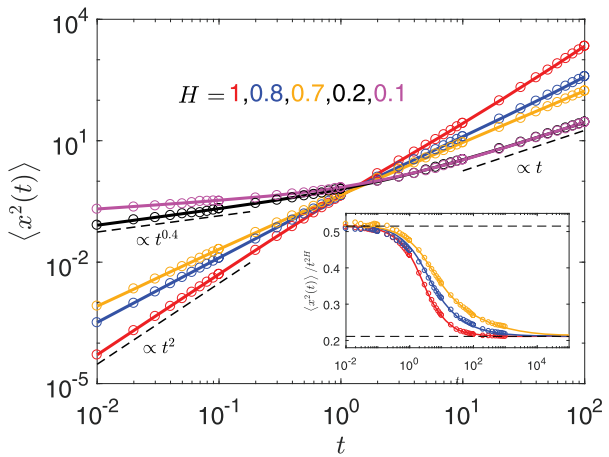


Viscoelastic diffusing-diffusivity model

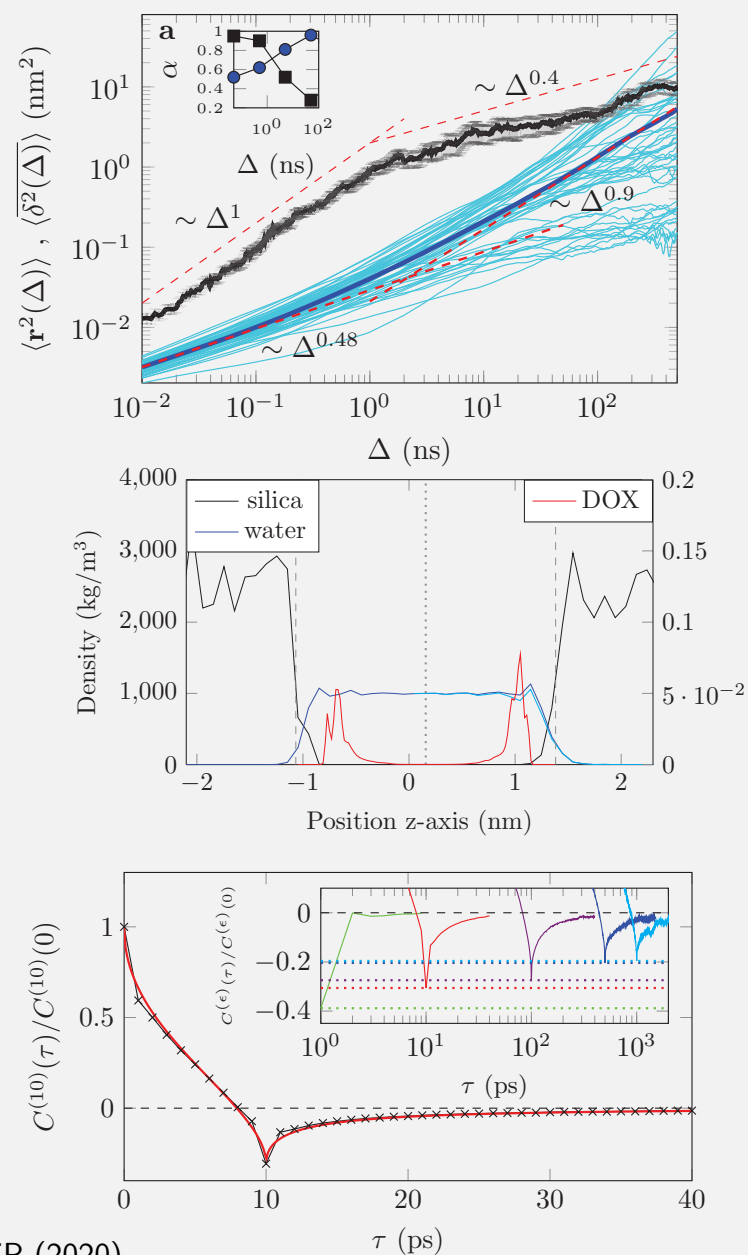
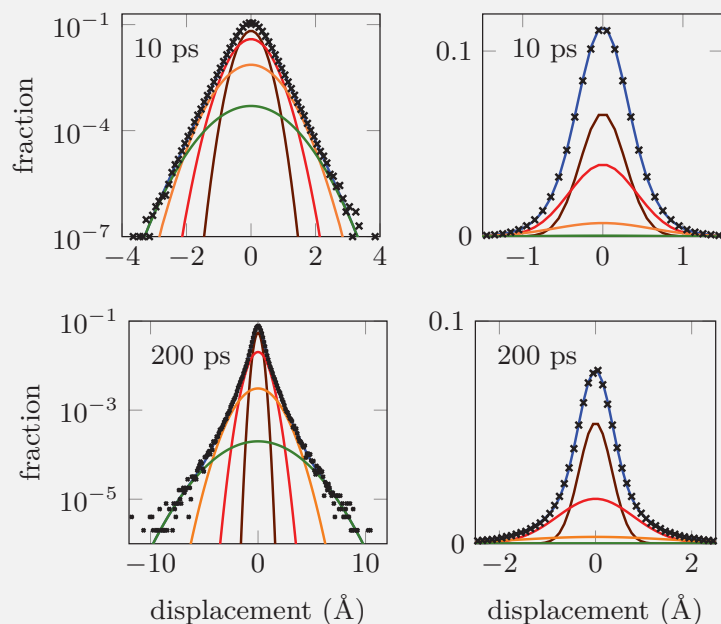
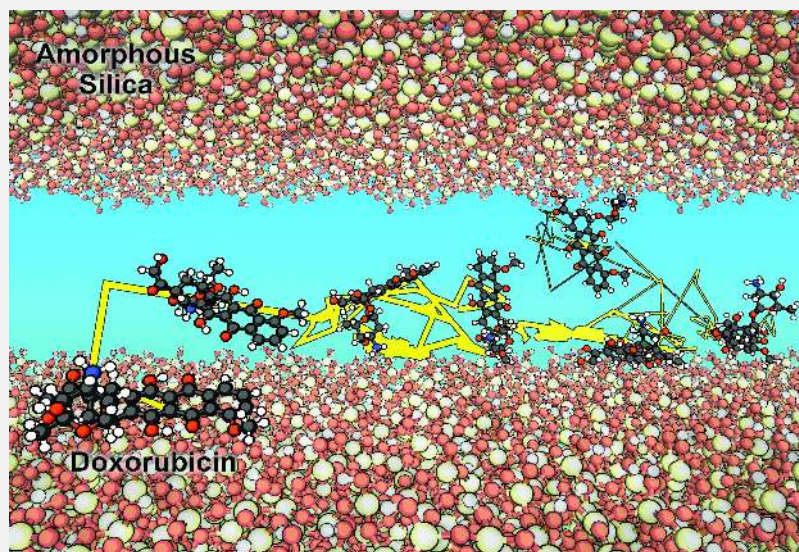
Tyagi model



Switching model



Doxorubicin drug molecule diffusion in silica nanochannels

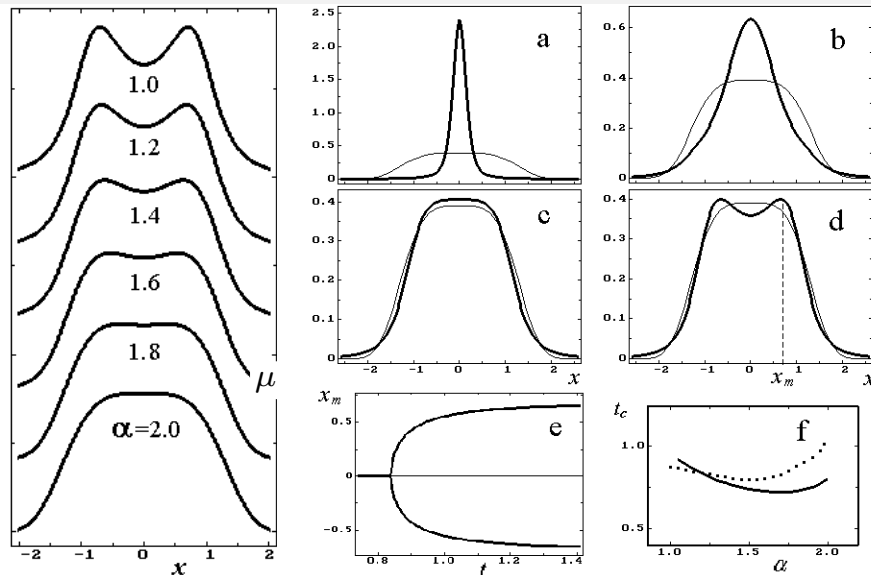


Lévy flights ($\lambda(x) \simeq |x|^{-1-\mu}$) in superharmonic potentials

$$\frac{\partial}{\partial t} P(x, t) = \frac{\partial}{\partial x} \frac{V'(x)}{m\eta} P(x, t) + K^\mu \frac{\partial^\mu}{\partial |x|^\mu} P(x, t), \quad V(x) \simeq |x|^c$$

Harmonic potential: $P_{\text{st}}(x) = L_\mu(x)$ [S Jespersen, RM & HC Fogedby, PRE (1999)]

Quartic potential $V(x) = \frac{1}{4}x^4$



Box-like potentials: K Capała, A Padash, AV Chechkin, B Shokri, RM & B Dybiec, Chaos (2020)

Subharmonic potentials:
Convergence to stationary state
only if:

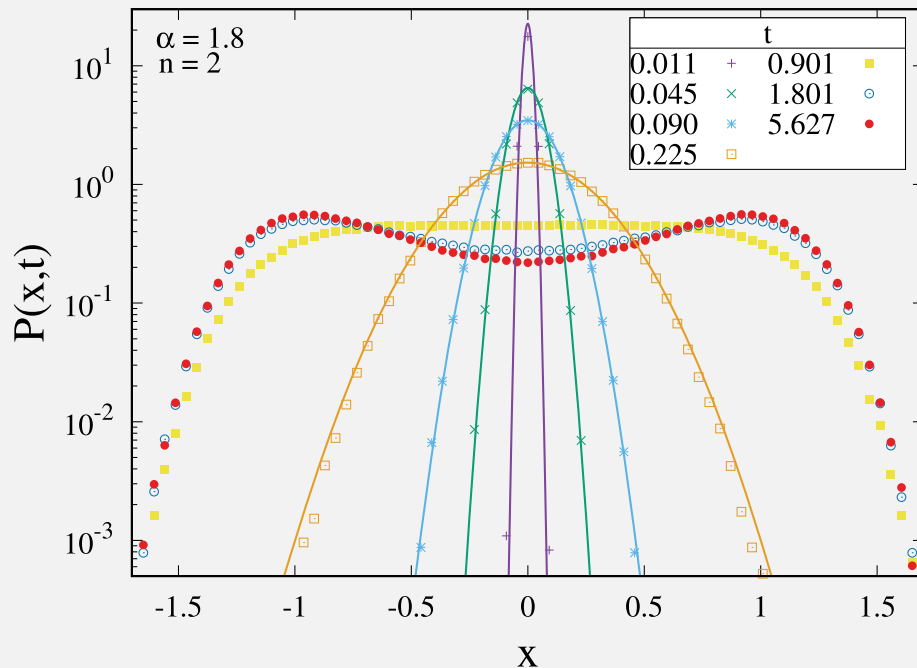
$$c > 2 - \mu, \quad 0 < \mu \leq 2$$

Otherwise motion is unconfined B Dybiec, AV Chechkin & IM Sokolov, JS-TAT (2010)

Fractional BM in super- and subharmonic potentials

$$\frac{d}{dt}X(t) = -V'(x)X(t) + \xi_\alpha(t) \quad \therefore \quad \langle \xi_\alpha(t)\xi_\alpha(t+\tau) \rangle \sim \alpha(\alpha-1)K_\alpha\tau^{\alpha-2}$$

Quartic potential



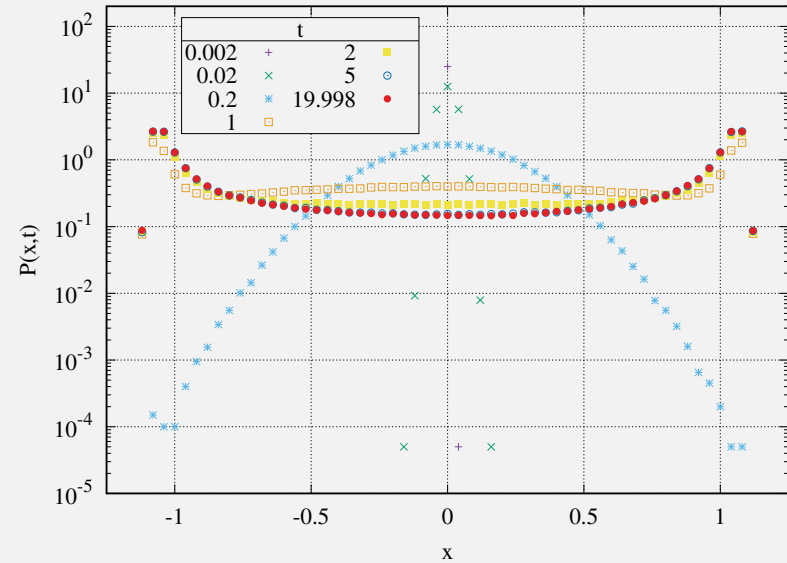
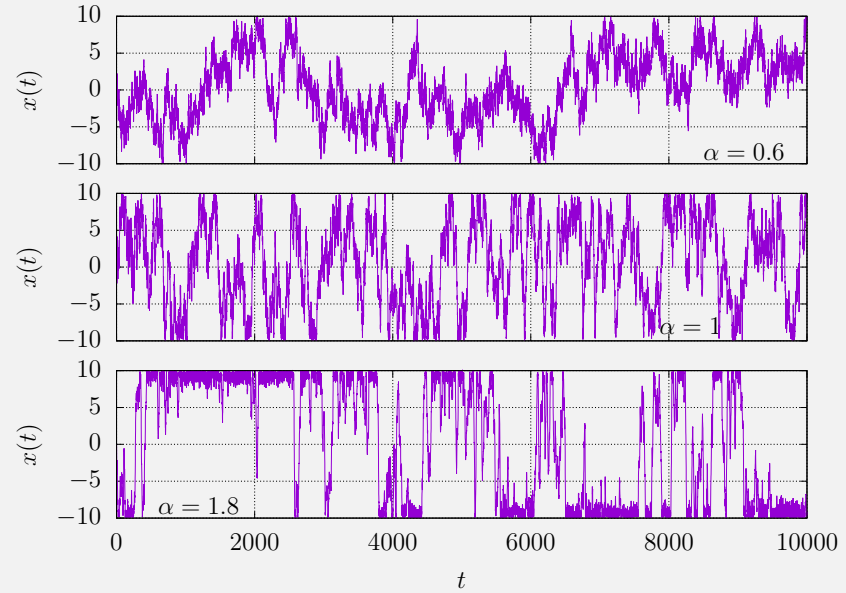
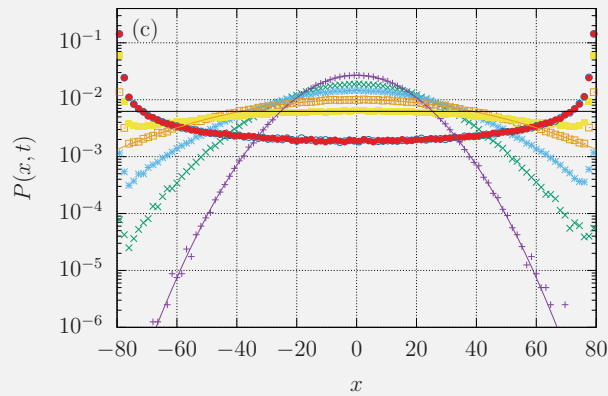
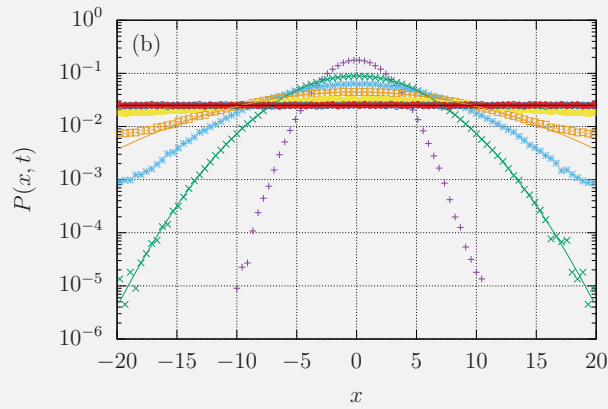
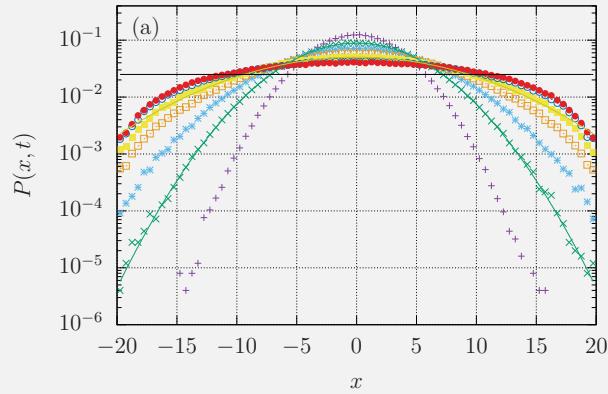
Convergence to stationarity only if
(free motion $\langle x^2(t) \rangle_{\text{FBM}} \simeq t^\alpha$,
 $\langle |x|^\delta(t) \rangle_{\text{LF}}^{2/\delta} \simeq t^{2/\mu}$):

$c > 2 - \frac{2}{\alpha}$ self-similarity index $H = \frac{\alpha}{2}$

$c > 2 - 1/H$

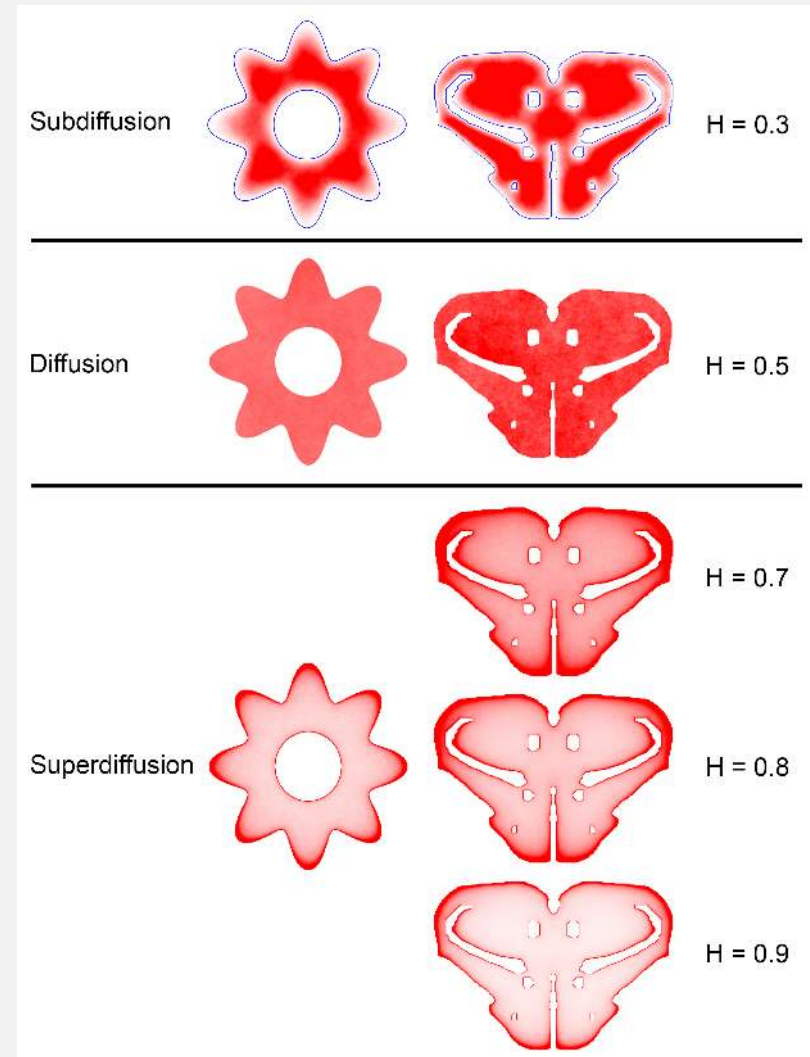
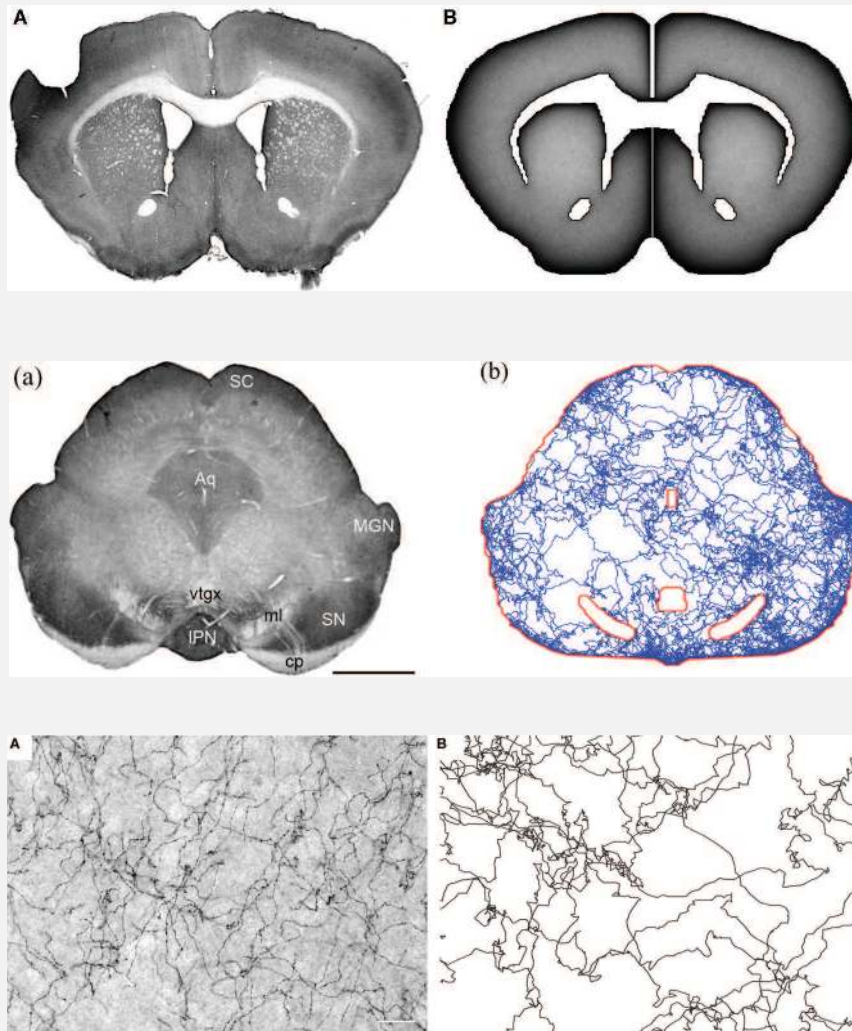
For Lévy flight $H = 1/\mu \Rightarrow$ same criterion for existence for stationary states

FBM: accretion & depletion effects near hard boundaries



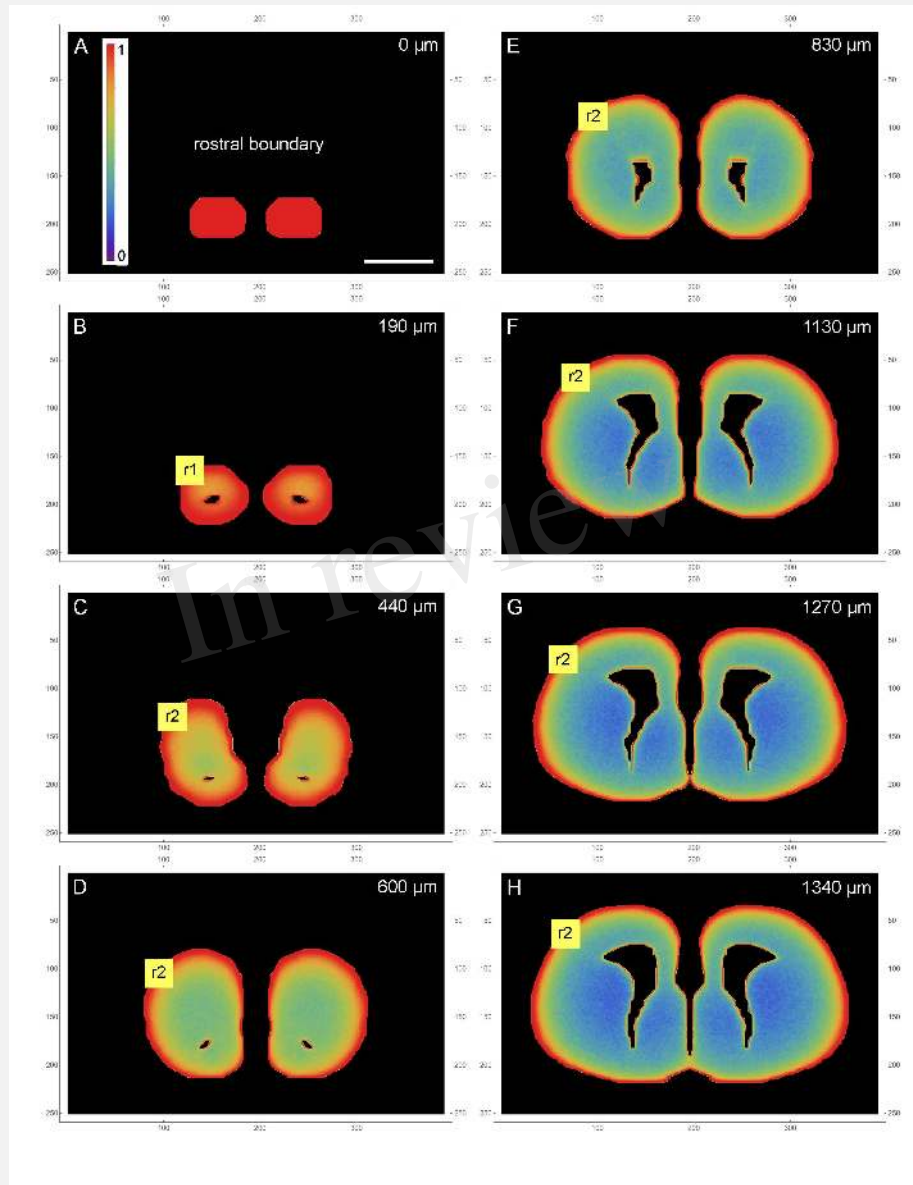
Brain serotonergic axons as FBM paths

Correlated fluctuations effect non-flat profile



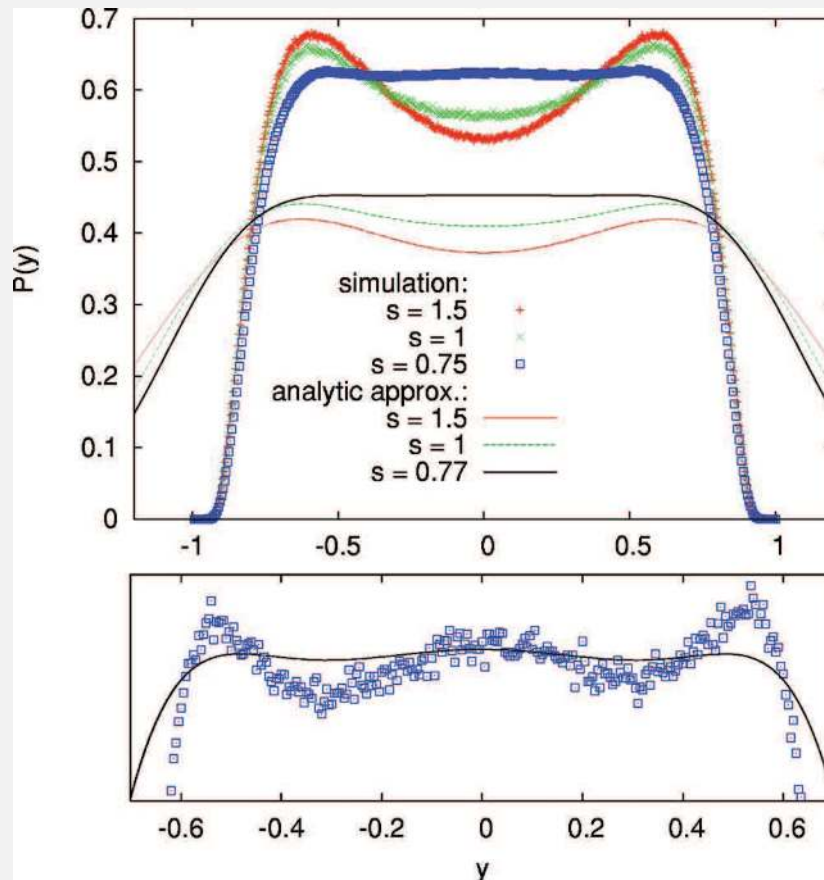
Brain serotonergic axons as FBM paths

Figure 2.TIF

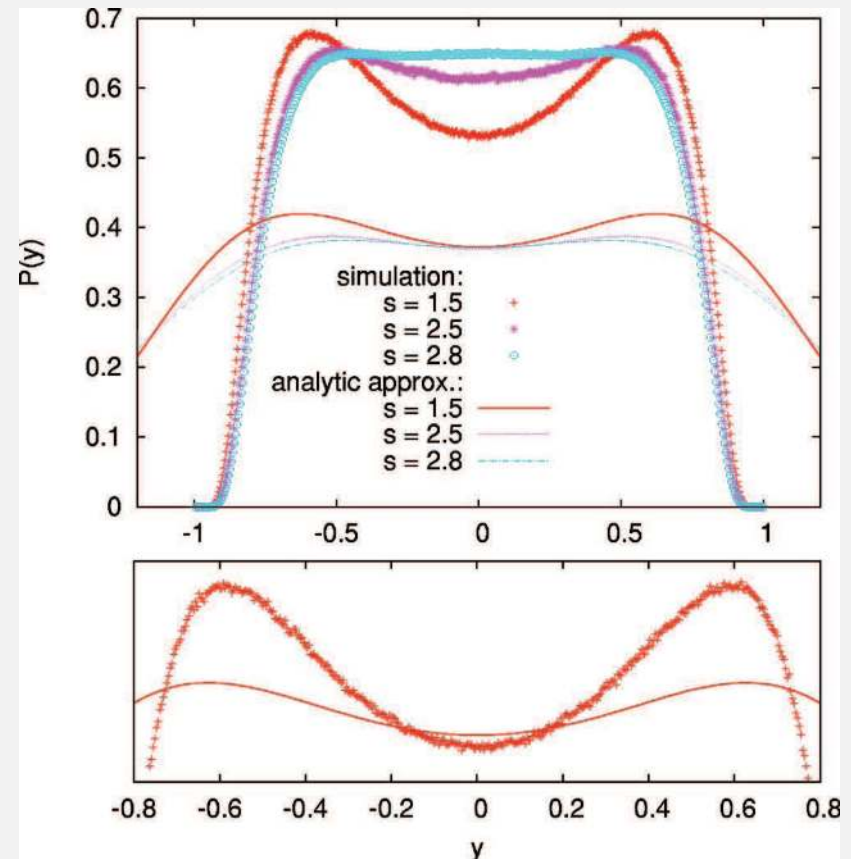


S Janušonis, JH Haiman, RM & T Vojta, Front Comp Neurosc (2022)

Bimodality in the transverse fluctuations of a grafted semiflexible polymer



$s = L/\ell_p$: stiff $\rightarrow$ semi-flexible



semi-flexible $\rightarrow$ flexible

Modelling in terms of worm-like chain model

Doubly-stochastic processes

“Roughness” in, i.a., financial data, is modelled by subdiffusive, antipersistent FBM

Roughness changing over time is often modelled by doubly-stochastic models, e.g.,

$$X(t) = \sqrt{\alpha(t)} \int_0^t (t-s)^{(\alpha(t)-1)/2} dB(s),$$

in Paul Lévy’s Riemann-Liouville formulation. $B(s)$ is the Wiener process $\rightsquigarrow$ locally varying roughness according to $\alpha(t)$ [e.g., A Ayache & J Lévy-Vehel (2000)]:

$$\langle X^2(t) \rangle \simeq t^{\alpha(t)}$$

$\rightsquigarrow$ where is the memory? $H(t)$, e.g., in cells during life cycle (density variation), in cooling membranes, in gel during cross-linking process etc.

Imagine a Brownian motion changing from K_1 to K_2 at $t = \tau$, described as

$$\tilde{B}(t) = \int_0^t \sqrt{K(s)} dB(s), \quad \rightsquigarrow \quad \langle X^2(t) \rangle = \begin{cases} 2Kt_1, & t \leq \tau \\ 2K_2(t - \tau) + 2K_1\tau, & t > \tau \end{cases}$$

whereas the above formulation would produce $\langle X^2(t) \rangle = 2K(t)t$

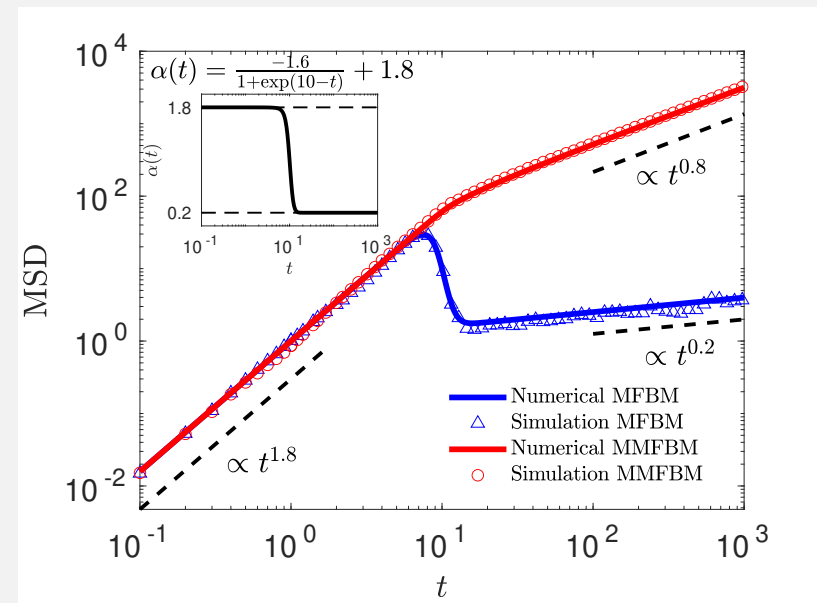
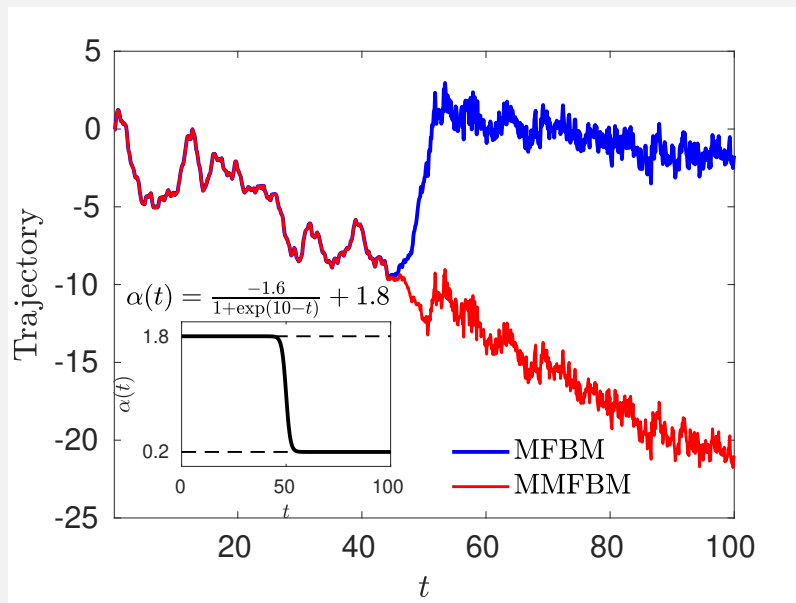
Doubly-stochastic processes

Solution is time-local scaling exponent (memory-multi-FBM):

$$X(t) = \int_0^t \sqrt{\alpha(s)}(t-s)^{(\alpha(s)-1)/2} dB(s),$$

Example of switching α , ($\alpha = \alpha_1$ for $t \leq \tau$, α_2 for $t > \tau$):

$$\langle X^2(t) \rangle = \begin{cases} t^{\alpha_1}, & t \leq \tau \\ t^{\alpha_1} - (t - \tau)^{\alpha_1} + (t - \tau)^{\alpha_2}, & t > \tau \end{cases} \sim (\alpha_1 \tau) t^{\alpha_1-1} + t^{\alpha_2}$$



Doubly-stochastic processes

Response function. Define the increments $X^\delta(\tau) = X(\tau + \delta) - X(\tau)$

Then consider a short α -perturbation:

$$\alpha(t) = \begin{cases} 1, & t \leq \tau \\ \alpha \neq 1, & \tau < t \leq \tau + \delta \\ 1, & t > \tau + \delta \end{cases}$$

Response function:

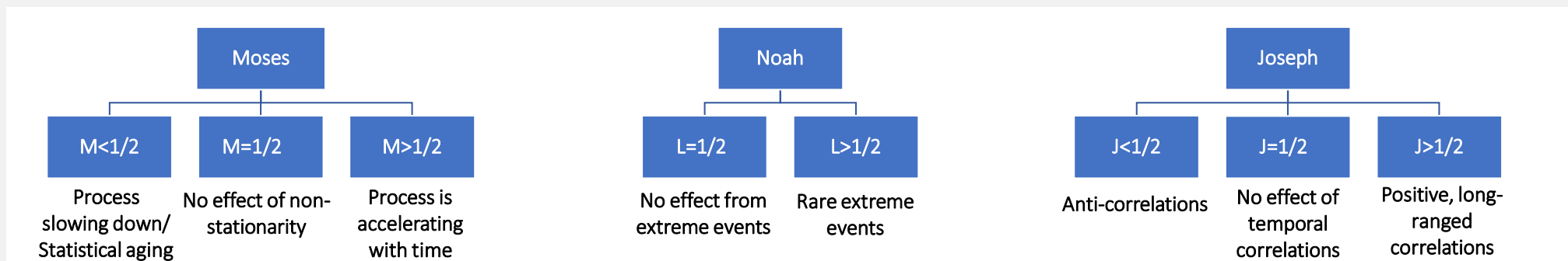
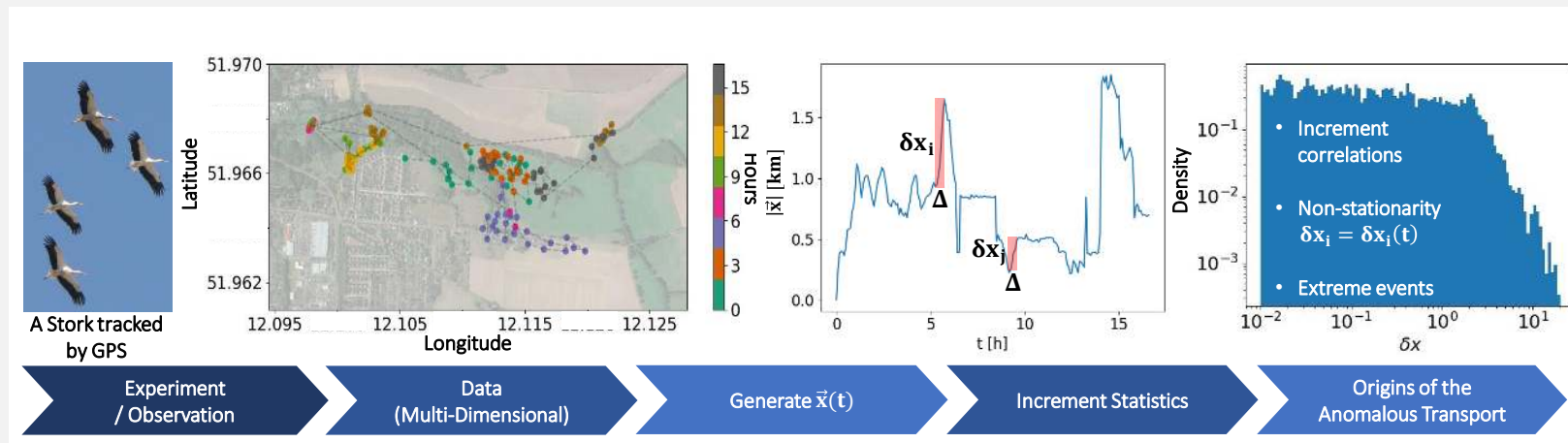
$$\langle X^\delta(\tau) X^\delta(\tau + T) \rangle = \alpha \frac{\alpha - 1}{2T^{1-\alpha}} \delta \mathbb{B} \left(\frac{\delta}{T}; \frac{\alpha + 1}{2}, 1 - \alpha \right)$$

for $\delta \rightarrow 0$, at time T after τ . When $T \rightarrow \infty$

$$\langle X^\delta(\tau) X^\delta(\tau + T) \rangle \sim \alpha [(\alpha - 1)/(\alpha + 1)] \delta^{\alpha/2+3/2} T^{\alpha/2-3/2}.$$

Even after a long period T a short perturbation still influences the process, with a sign depending on whether $\alpha \gtrless 1$

Scaling analysis of anomalous diffusion: molecules $\rightarrow$ animals



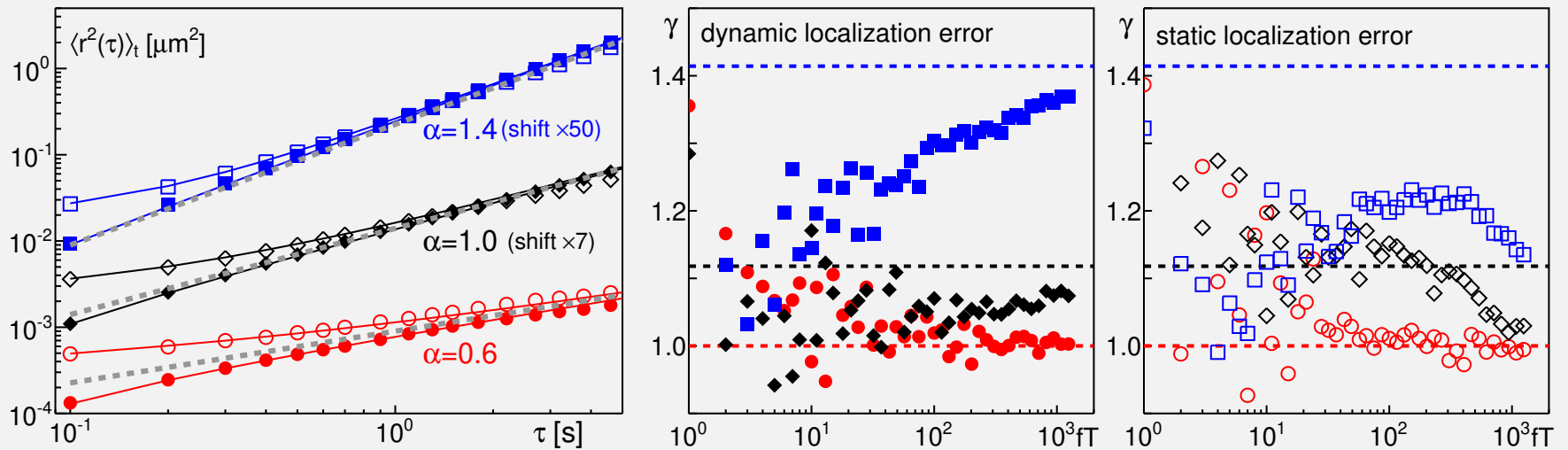
Robust criterion for anomalous diffusion

Two sources of noise: static (finite photon count, additive OU) & dynamic (finite measurement time). Static noise effect on γ :

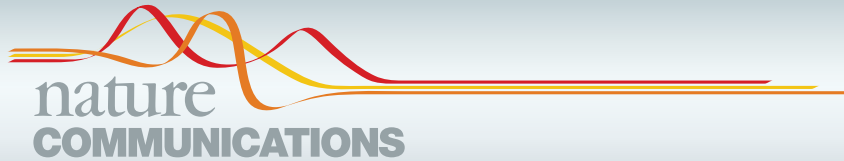
$$H < 1/2 : \quad \gamma_Y^2(f, T) \sim 1; \quad \sigma_e = 0 : 1$$

$$H = 1/2 : \quad \gamma_Y^2(f, T) \sim 1 + \frac{1}{4} \left(1 + \frac{\sigma_e^2}{2u^2 D} f^2 \right)^{-2}; \quad \sigma_e = 0 : \frac{\sqrt{5}}{2}$$

$$H > 1/2 : \quad \gamma_Y^2(f, T) \sim 1 + \left(1 + \frac{\sigma_e^2}{u^2 D} T^{1-2H} f^2 \right)^{-2}; \quad \sigma_e = 0 : \sqrt{2}$$



The ANomalous Diffusion challenge



ARTICLE



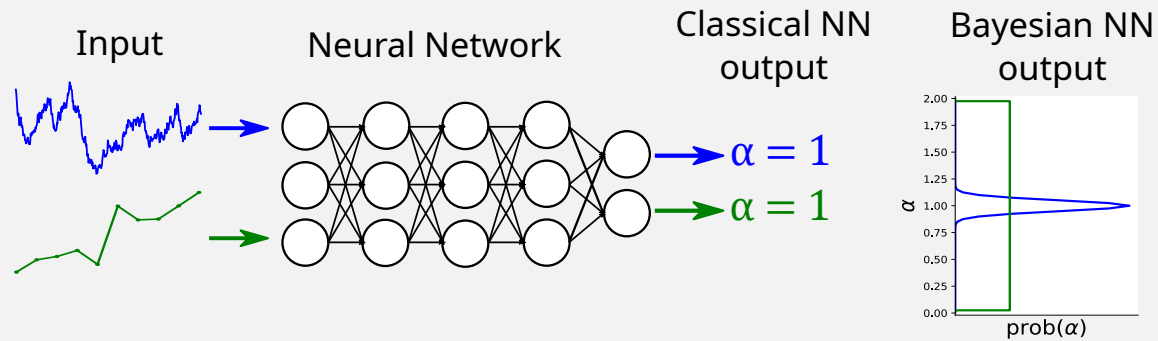
<https://doi.org/10.1038/s41467-021-26320-w>

OPEN

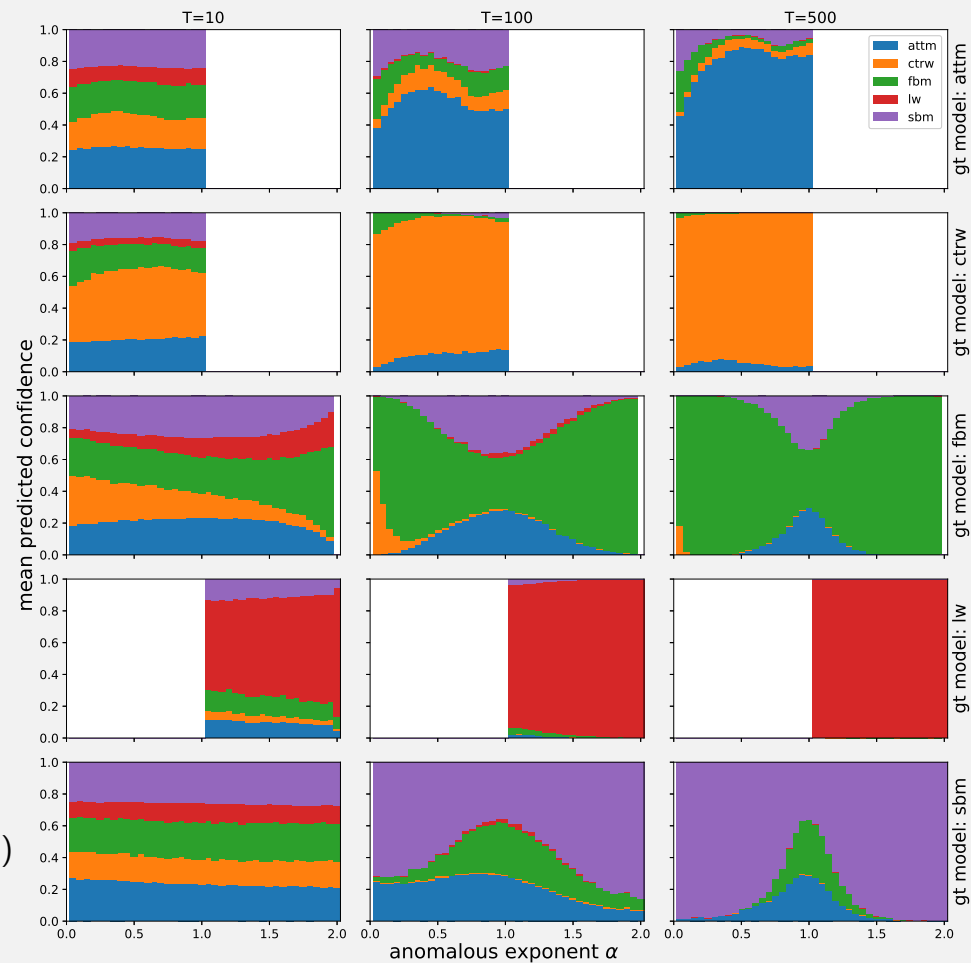
Objective comparison of methods to decode anomalous diffusion

Gorka Muñoz-Gil ¹, Giovanni Volpe ²✉, Miguel Angel Garcia-March ³, Erez Aghion ⁴, Aykut Argun ², Chang Beom Hong ⁵, Tom Bland ⁶, Stefano Bo ⁴, J. Alberto Conejero ³, Nicolás Firbas ³, Òscar Garibo i Orts ³, Alessia Gentili ⁷, Zihan Huang ⁸, Jae-Hyung Jeon ⁵, Hélène Kabbech ⁹, Yeongjin Kim ⁵, Patrycja Kowalek ¹⁰, Diego Krapf ¹¹, Hanna Loch-Olszewska ¹⁰, Michael A. Lomholt ¹², Jean-Baptiste Masson ¹³, Philipp G. Meyer ⁴, Seongyu Park ⁵, Borja Requena ¹, Ihor Smal ⁹, Taegeun Song ^{5,14,15}, Janusz Szwabiński ¹⁰, Samudrajit Thapa ^{16,17,18}, Hippolyte Verdier ¹³, Giorgio Volpe ⁷, Artur Widera ¹⁹, Maciej Lewenstein ^{1,20}, Ralf Metzler ¹⁶ & Carlo Manzo ^{1,21}✉

Bayesian-weighted deep learning model selection



“Beyond the Black Box problem”



H Seckler & RM, Nat Comm (2022)

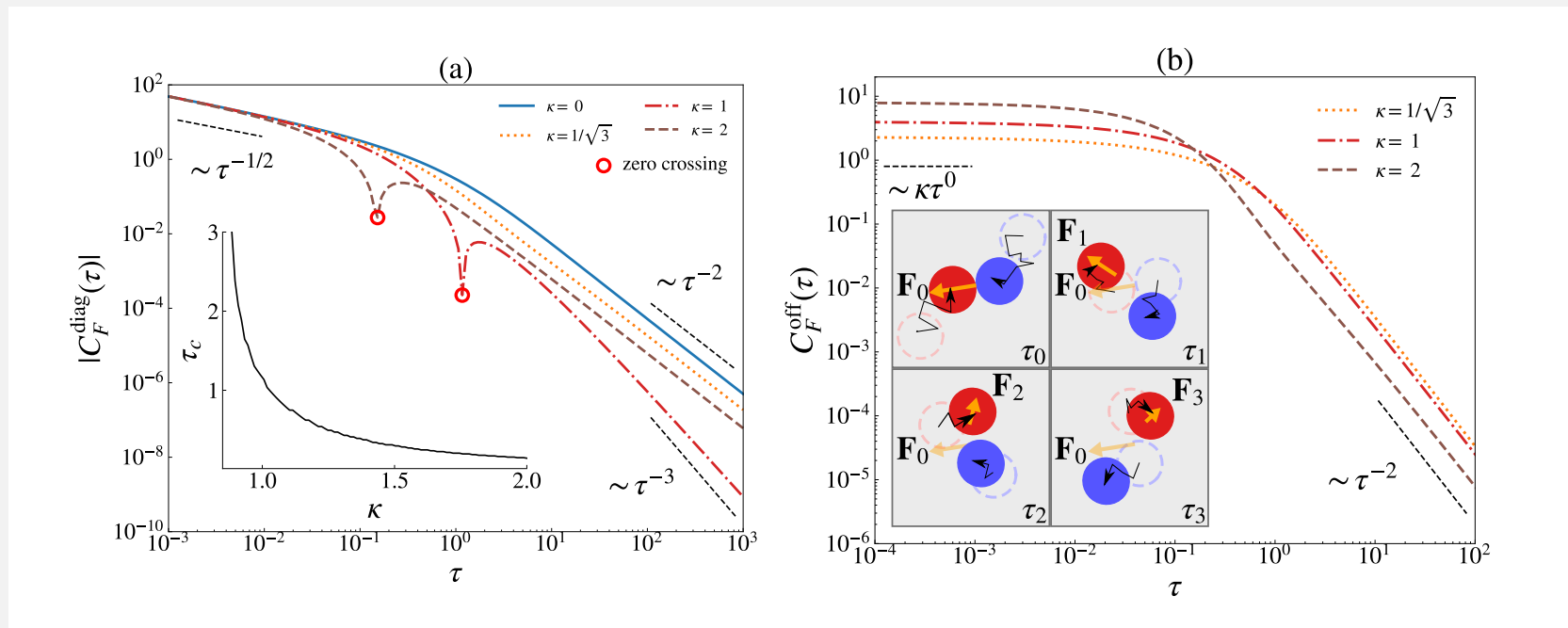
Odd-diffusive systems at equilibrium

Odd-diffusive system has diffusion tensor $\mathbf{D} = D_0(\mathbf{1} + \kappa\boldsymbol{\varepsilon})$: Lorentz force, skyrmions

$\boldsymbol{\varepsilon}$ Levi-Civita: $\varepsilon_{xy} = -\varepsilon_{yx} = 1$; $\varepsilon_{xx} = \varepsilon_{yy} = 0$

Force autocorrelation shows oscillations at equilibrium, contradicting claims that it should decay monotonically [M Caraglio & T Franosch, PRL (2022)]

Conjecture for tail of the correlation function $\simeq \tau^{-3/2}$ also violated



Σ ummary

- I Single particle tracking & supercomputing provide access to unprecedented insight into microscopic dynamics of complex systems
- II Long-range correlated motion is revealed across scales, from drug molecules in silica slabs over intracellular motion to flying birds
- III Inhomogeneity of tracer ensemble and/or environment cause pronouncedly non-Gaussian displacement distributions
- IIII Long-range correlated processes defy mathematical analysis (e.g., first-passage) & show strange interactions with boundaries
- IIII Objective data analysis methods aid in identification of models & parameters: Bayesian-maximum likelihood, deep learning, & their combinations. Measurement error! AnDi challenge special issue of JPA & AnDi 2!

For slides or questions: write to rmetzler@uni-potsdam.de