

9.3 Symmetric stable laws: PDFs and moments (continued)

Remark 5. Important!

For any CF/PDF small k -behavior of $\hat{p}(k) \leftrightarrow$

$\leftrightarrow$ large- x behavior of $p(x)$

$$\hat{l}_\alpha(k) = e^{-\sigma^\alpha |k|^\alpha}, \quad 0 < \alpha < 2; \quad \alpha = 2 \text{ Gaussian CF}$$

$$k \rightarrow 0 \quad \hat{l}_\alpha(k) \sim 1 - \sigma^\alpha |k|^\alpha + \dots$$

$$|x| \rightarrow \infty \quad l_\alpha(x) \sim C_1(\alpha) \frac{\sigma^\alpha}{|x|^{1+\alpha}}, \quad C_1(\alpha) = \frac{1}{\pi} \sin\left(\frac{\pi\alpha}{2}\right) \Gamma(1+\alpha)$$

$$\hat{l}_\alpha(k) = \int_{-\infty}^{\infty} e^{ikx} l_\alpha(x) dx = 2 \int_0^{\infty} \cos(kx) l_\alpha(x) dx = 1 - 2 \int_0^{\infty} (1 - \cos(kx)) l_\alpha(x) dx \sim$$

$$\int_{-\infty}^{\infty} l_\alpha(x) dx = 2 \int_0^{\infty} l_\alpha(x) dx = 1$$

$$\sim 1 - 2 C_1(\alpha) \sigma^\alpha \int_0^{\infty} \frac{1 - \cos kx}{x^{1+\alpha}} dx \Rightarrow$$

$$\int_0^{\infty} \frac{1 - \cos(kx)}{x^{1+\alpha}} dx = \{y = |k|x\} = |k|^\alpha \int_0^{\infty} \frac{1 - \cos y}{y^{1+\alpha}} dy = \frac{|k|^\alpha}{\alpha} \int_0^{\infty} \frac{\sin y}{y^\alpha} dy =$$

$$= \frac{|k|^\alpha}{2 C_1(\alpha)}$$

$$\Rightarrow \hat{l}_\alpha(k) \sim 1 - \sigma^\alpha |k|^\alpha$$

Opposite direction: $1 - \sigma^\alpha |k|^\alpha \rightarrow C_1(\alpha) \frac{\sigma^\alpha}{|x|^{1+\alpha}}$

Requires the theory of improper integrals

9.4 Generalized Central Limit Theorem for symmetric probability laws

$X_1, X_2, \dots, X_n$ iid $p_X(x)$, $p_X(x) = p_X(-x)$, $\langle X^2 \rangle = \infty$

$$\text{Let } p(x) \sim \frac{A}{|x|^{1+\alpha}} \text{ as } |x| \rightarrow \pm\infty, \quad 0 < \alpha < 2 \Rightarrow \langle X^2 \rangle = 2 \int_0^{\infty} x^2 p(x) dx = \infty$$

$$A = \text{const} \quad \frac{1}{x^{\alpha-1}} \text{ divergence at } x = \infty$$

Properly normalized sum (Reminder: CLT $Y = \frac{1}{\sqrt{n}} \sum_{j=1}^n X_j$)

$$Y_n = \frac{1}{n^{1/\alpha}} \sum_{j=1}^n X_j$$

$$\text{Theorem } p_{Y_n}(y) \xrightarrow{n \rightarrow \infty} l_\alpha(y) \sim \frac{A}{|y|^{1+\alpha}}, \quad y \rightarrow \infty, \quad 0 < \alpha < 2$$

Two important statements

1. Only asymptotic behavior of $p_X(x)$ at $x \rightarrow \infty$ is important for convergence to stable law.
2. The limiting stable PDF $l_\alpha(y)$ has exactly the same

for convergence to stable
 2. The limiting stable PDF $l_y(y)$ has exactly the same asymptotic behavior as $p_X(x)$

Remark $l_x(y) \div \hat{l}_x(k) = e^{-\sigma_y^\alpha |k|^\alpha}$

$$A = \frac{\sigma_y^\alpha}{\pi} \sin\left(\frac{\pi\alpha}{2}\right) \Gamma(1+\alpha)$$

$$\sigma_y = \left[\frac{\pi A}{\sin\left(\frac{\pi\alpha}{2}\right) \Gamma(1+\alpha)} \right]^{1/\alpha}$$

Summary of GCLT

$X_1, \dots, X_n$ iid $p_X(x)$

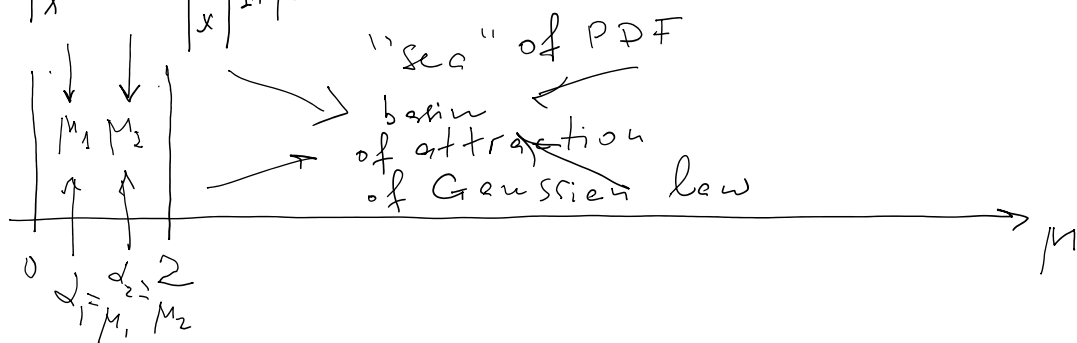
CLT $\rightarrow$ basin of attraction of the Gaussian law
 $p_X(x)$ with finite variance ($\alpha=2$)

GCLT $\rightarrow$ basin of attraction of α -stable laws
 $p_X(x)$ with infinite variance and the same
 asymptotic behavior as α -stable law, $0 < \alpha < 2$

Remark $p(x) \sim \frac{1}{x^3}$? $Y_n = \frac{1}{\sqrt{n \ln n}} \sum_{j=1}^n X_j$

$p_{Y_n}(y) \xrightarrow{n \rightarrow \infty}$ Gaussian law

$p_X(x) \sim \frac{A}{|x|^{1+\mu}}$, $0 < \mu < 2$, $x \rightarrow \infty \Rightarrow p_Y(y) \xrightarrow{n \rightarrow \infty}$ stable law, $\alpha = \mu$



(9.5) Linear transformation $-|k|^\alpha$ stable analogue of the standard normal law ($\sigma=1$)
 r.v. $l_x(x) \div \hat{l}_x(k) = e$

$Y = \sigma X + \mu$
 $\hat{p}_Y(k) = \langle e^{iky} \rangle = e^{i\mu k} \hat{l}_x(\sigma k) \div l_x(y) = l_x(y = \sigma x + \mu) \equiv$

$$\hat{p}_Y(k) = \langle e^{iky} \rangle = e^{i\mu k} \hat{l}_\alpha(\sigma k) \div l_\alpha(y) = l_\alpha(y = \sigma x + \mu) \equiv l_\alpha(y; \sigma, \mu)$$

Terminology σ scale parameter $\leftrightarrow$ standard deviation at $\alpha = 2$
 μ shift parameter

$$\langle Y \rangle = \mu \leftrightarrow \mu = \text{mean value if } 1 < \alpha \leq 2$$

$$\mu = 0 \Rightarrow p_Y(y) = p_Y(-y) \text{ symmetric PDF}$$

Important remark μ and σ can be eliminated by a simple linear transformation

$$\boxed{l_\alpha(y; \sigma, \mu) = \frac{1}{\sigma} l_\alpha\left(\frac{y - \mu}{\sigma}; 1, 0\right)} \quad (A)$$

Take (A), $\times e^{iky}$ and $\int_{-\infty}^{\infty} dy$:

$$\text{lhs of (A)} = \int_{-\infty}^{\infty} e^{iky} l_\alpha(y; \sigma, \mu) dy = e^{i\mu k - \sigma^\alpha/k|^\alpha}$$

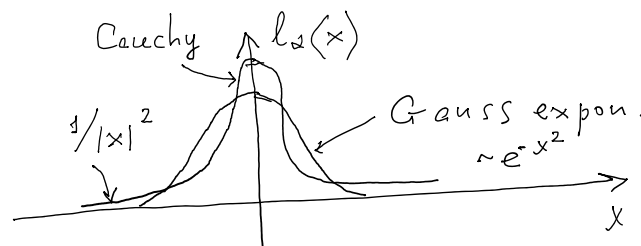
$$\begin{aligned} \text{rhs of (A)} &= \frac{1}{\sigma} \int_{-\infty}^{\infty} e^{iky} l_\alpha\left(\frac{y - \mu}{\sigma}; 1, 0\right) dy = \left\{ x = \frac{y - \mu}{\sigma} \right\} = \\ &= \int_{-\infty}^{\infty} e^{ik(\sigma x + \mu)} l_\alpha(x; 1, 0) dx = e^{i\mu k} \int_{-\infty}^{\infty} e^{ik\sigma x} l_\alpha(x; 1, 0) dx = \\ &= e^{i\mu k} \hat{l}_\alpha(\sigma k) \end{aligned}$$

(9.6) PDF plots

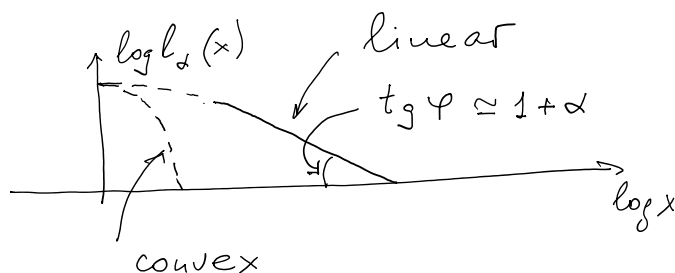
$$\begin{aligned} l_\alpha(x=0) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-ikx} \hat{l}_\alpha(k) dk \Big|_{x=0} = \frac{2}{2\pi} \int_0^{\infty} e^{-k^\alpha} dk = \\ &= \left\{ y = k^\alpha, \quad k = y^{1/\alpha}, \quad dk = \frac{1}{\alpha} y^{\frac{1}{\alpha}-1} dy \right\} = \frac{1}{\pi\alpha} \int_0^{\infty} e^{-y} y^{\frac{1}{\alpha}-1} dy = \\ &= \frac{\Gamma(1/\alpha)}{\pi\alpha} \end{aligned}$$

$$\alpha = 2 \quad l_2(0) = \frac{\Gamma(1/2)}{2\pi} \frac{1}{2\sqrt{\pi}}$$

$$\alpha = 1 \quad l_1(0) = \frac{1}{\pi}$$



$$l_\alpha(x) \sim \frac{1}{x^{1+\alpha}}, \quad x \rightarrow \infty$$



9.7 General form of CF for α -stable probability laws

$$\hat{l}_{\alpha, \beta}(k; \sigma, \mu) = \exp \left[i \mu k - \sigma^\alpha |k|^\alpha \left(1 - i \beta \operatorname{sgn}(k) w(k, \alpha) \right) \right]$$

$0 < \alpha \leq 2$ index of stability (Lévy index)

$-1 \leq \beta \leq 1$ parameter of asymmetry

μ shift parameter, real

σ scale parameter, positive

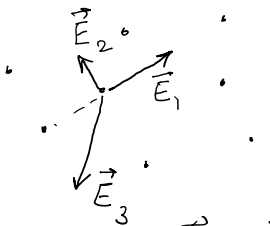
• μ and σ are not so important as α and β :
can be removed by linear transformation, see Eq. (A)

• $\beta = 0$ for symmetric laws $l_\alpha(x) = l_\alpha(-x)$

$$w(k, \alpha) = \begin{cases} \tan\left(\frac{\pi\alpha}{2}\right), & \alpha \neq 1, \quad 0 < \alpha < 2 \\ -\frac{2}{\pi} \ln|k|, & \alpha = 1 \end{cases}$$

$\alpha = 2 \Rightarrow \beta$ irrelevant

(9.8) Holtzmark distribution (1919)
system of identical charged particles distributed randomly in space

$$\vec{E}_i = \frac{e \vec{r}_i}{r_i^3} \quad n = \frac{N}{V} = \text{const}$$


$$\vec{E} = \vec{E}_1 + \dots + \vec{E}_n, \quad p(\vec{E}) = ?$$

$$p(\vec{E}) = \int_{-\infty}^{\infty} e^{-i \vec{k} \cdot \vec{E}} \hat{p}_E(\vec{k}) \frac{d\vec{k}}{2\pi} \sim a/|\vec{k}|^{3/2}$$

Rigorous calculation $\hat{p}_E(\vec{k}) = e$

$$\hat{p}(\vec{E}) \sim \frac{1}{|\vec{E}|^{5/2}}$$

3-dimensional symmetric
 α -stable law with $\alpha = 3/2$

Similar examples: systems with long-range interactions

- gravity field of stars $\alpha = 3/2$
- electric field of dipoles $\alpha = 1$
- velocity field of point vortices in fully developed turbulence $\alpha = 3/2$

Discussion Assignment 1

#(1.2) Skewness & Kurtosis $K_3 = \frac{\overline{M}_3}{\sigma^3}, \quad K_4 = \frac{\overline{M}_4}{\sigma^4}$

$$\overline{M}_n = \langle (x - \mu_1)^n \rangle \quad \mu_1 = \langle x \rangle$$

At the lecture $\mu = \langle x \rangle = 0$

$$\langle x \rangle_G = \int_{-\infty}^{\infty} x \frac{1}{\sqrt{2\pi}\sigma^2} \exp\left[-\frac{(x-\mu)^2}{2\sigma^2}\right] dx =$$

$$= \int_{-\infty}^{\infty} (x - \mu + \mu) \dots dx = \int_{-\infty}^{\infty} (x - \mu) \dots dx = 0$$

$$\overline{M}_2 = \int_{-\infty}^{\infty} (x - \mu)^2 \frac{1}{\sqrt{2\pi}\sigma^2} \exp\left[-\frac{(x-\mu)^2}{2\sigma^2}\right] dx = \sigma^2$$

$$= \int_{-\infty}^{\infty} y \frac{1}{\sqrt{2\pi}\sigma^2} \exp\left[-\frac{y^2}{2\sigma^2}\right] dy + \mu \int_{-\infty}^{\infty} \frac{e^{-y/2\sigma^2}}{\sqrt{2\pi}\sigma^2} dy = \mu$$

$$\text{Var}\{x\} = \frac{1}{\sqrt{2\pi}\sigma^2} \int_{-\infty}^{\infty} (x-\mu)^2 \exp\left[-\frac{(x-\mu)^2}{2\sigma^2}\right] dx = \sigma^2$$

$$\text{Check. } K_4 = \frac{\overline{\mu_4}}{\sigma^4} = \frac{1}{\sigma^4} \int_{-\infty}^{\infty} (x-\mu)^4 \frac{1}{\sqrt{2\pi}\sigma^2} \exp\left[-\frac{(x-\mu)^2}{2\sigma^2}\right] dx = 3$$

$$\text{Excessive kurtosis } K_{ex4} = K_4 - 3$$

$$\text{Laplace } p_L(x) = \frac{1}{2b} \exp\left(-\frac{|x-\mu|}{b}\right)$$

$$\text{Norm } \int_{-\infty}^{\infty} p_L(x) dx = \int_{-\infty}^{\infty} \delta y = x - \mu = \frac{1}{2b} \int_{-\infty}^{\infty} e^{-|x|/b} dx = \frac{2}{2b} \int_0^{\infty} e^{-x/b} dx =$$

$$= \int_0^{\infty} e^{-y} dy = 1$$

$$\text{1st moment } \int_{-\infty}^{\infty} x p_L(x) dx = \int_{-\infty}^{\infty} (x-\mu + \mu) \frac{1}{2b} e^{-|x|/b} dx = \mu \underbrace{\int_{-\infty}^{\infty} \frac{1}{2b} e^{-|x|/b} dx}_{=1} = \mu$$

$$\text{Var}\{x\} = \int_{-\infty}^{\infty} (x-\mu)^2 \frac{1}{2b} e^{-|x-\mu|/b} dx = \frac{1}{2b} \int_{-\infty}^{\infty} x^2 e^{-|x|/b} dx = \frac{1}{b} \int_0^{\infty} x^2 e^{-x/b} dx = \frac{1}{b} \left[\frac{x^2}{2} \int_0^{\infty} e^{-\alpha x} dx \right]_{\alpha=1/b} = \frac{1}{b} \left[\frac{x^2}{2\alpha^2} \frac{1}{\alpha} \right]_{\alpha=1/b} =$$

$$= \frac{1}{b} \left[\frac{\partial}{\partial \alpha} \left(-\frac{1}{\alpha^2} \right) \right]_{\alpha=1/b} = \frac{2}{b} \frac{1}{\alpha^3} \Big|_{\alpha=1/b} = 2b^2 = \sigma^2$$

$$\overline{\mu_4} = \frac{1}{2b} \int_{-\infty}^{\infty} x^4 e^{-|x|/b} dx = \frac{2}{2b} \int_0^{\infty} x^4 e^{-x/b} dx = \frac{1}{b} \left[\frac{x^4}{2\alpha^4} \int_0^{\infty} e^{-\alpha x} dx \right]_{\alpha=1/b} =$$

$$= \frac{1}{b} \left[\frac{\partial^4}{\partial \alpha^4} \frac{1}{\alpha} \right]_{\alpha=1/b} = 24b^4$$

$$K_4 = \frac{24b^4}{4b^4} = 6$$

Problem 4
$$\rho = \frac{\text{cov}(X, Y)}{\sqrt{\text{Var}[X] \text{Var}[Y]}}$$

$$\langle X \rangle = \langle Y \rangle = 0, \quad \text{Var}[X] = \langle X^2 \rangle, \quad \text{Var}[Y] = \langle Y^2 \rangle$$

a) start from
$$\left\langle \left(\frac{X}{\sigma_x} \pm \frac{Y}{\sigma_y} \right)^2 \right\rangle \geq 0$$

$$\underbrace{\frac{\langle X^2 \rangle}{\sigma_x^2}}_{=1} \pm \underbrace{\frac{2\langle XY \rangle}{\sigma_x \sigma_y}}_{=2\rho} + \underbrace{\frac{\langle Y^2 \rangle}{\sigma_y^2}}_{=1} \geq 0$$

$$2 \pm 2\rho \geq 0 \Rightarrow -1 \leq \rho \leq 1$$

b) $Y = aX$

$$\rho = \frac{\langle XY \rangle}{\sqrt{\langle X^2 \rangle \langle Y^2 \rangle}} = \frac{a \langle X^2 \rangle}{|a| \langle X^2 \rangle} = \text{sign } a = \pm 1$$

Problem 6
$$p_z(z) = \iint_{-\infty}^{\infty} p(x, y) \delta(z - f(x, y)) dx dy$$

$$z = \frac{x}{y} \quad p_z(z) = \iint_{-\infty}^{\infty} \underbrace{p(x, y)}_{p(x)p(y)} \delta\left(z - \frac{x}{y}\right) dx dy = \left\{ \delta(ax) = \frac{1}{|a|} \delta(x) \right\}$$

$$= \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dx p_X(x) p_Y(y) \delta\left[\frac{1}{|y|}(x - zy)\right] = \int_{-\infty}^{\infty} dy \sqrt{|y|} p_Y(y) \int_{-\infty}^{\infty} dx p_X(x) \delta(x - zy) \Rightarrow$$

$$\Rightarrow \boxed{p_z(z) = \int_{-\infty}^{\infty} p_Y(y) p_X(zy) |y| dy}$$

$$p_X(x) = \frac{1}{\sqrt{2\pi\sigma_x^2}} \exp\left(-\frac{x^2}{2\sigma_x^2}\right), \quad p_Y(y) = \frac{1}{\sqrt{2\pi\sigma_y^2}} e^{-x^2/(2\sigma_y^2)}$$

$$p_z(z) = \frac{1}{2\pi\sigma_x\sigma_y} \int_{-\infty}^{\infty} e^{-y^2/2\sigma_y^2 - z^2 y^2/2\sigma_x^2} |y| dy = \left\{ A = \frac{z^2}{2\sigma_x^2} + \frac{1}{2\sigma_y^2} \right\} =$$

$$= \frac{1}{\pi\sigma_x\sigma_y} \int_0^{\infty} y e^{-Ay^2} dy = \left\{ \eta = y^2, d\eta = 2y dy \right\} =$$

$$\frac{1}{2\pi\sigma_x\sigma_y} \int_0^{\infty} e^{-A\eta} d\eta = \frac{1}{2\pi\sigma_x\sigma_y} \left[-\frac{1}{A} e^{-A\eta} \right]_0^{\infty} = \frac{1}{2\pi\sigma_x\sigma_y} \frac{1}{A} = \frac{1}{2\pi\sigma_x\sigma_y} \frac{2\sigma_x^2\sigma_y^2}{z^2\sigma_y^2 + \sigma_x^2} = \frac{2\sigma_x\sigma_y}{z^2\sigma_y^2 + \sigma_x^2}$$

$$\begin{aligned}
 & \pi \sigma_x \sigma_y \text{ or } 1 \\
 & = \frac{1}{2\pi \sigma_x \sigma_y} \int_0^\infty e^{-A\eta} d\eta = \frac{1}{2\pi \sigma_x \sigma_y A(z)} = \{ \sigma_x^2 = \sigma_y^2 = 1 \} = \\
 & = \frac{1}{2\pi \left(\frac{z^2}{2} + \frac{1}{2} \right)} = \frac{1}{\pi (1+z^2)} = \text{Cauchy}
 \end{aligned}$$