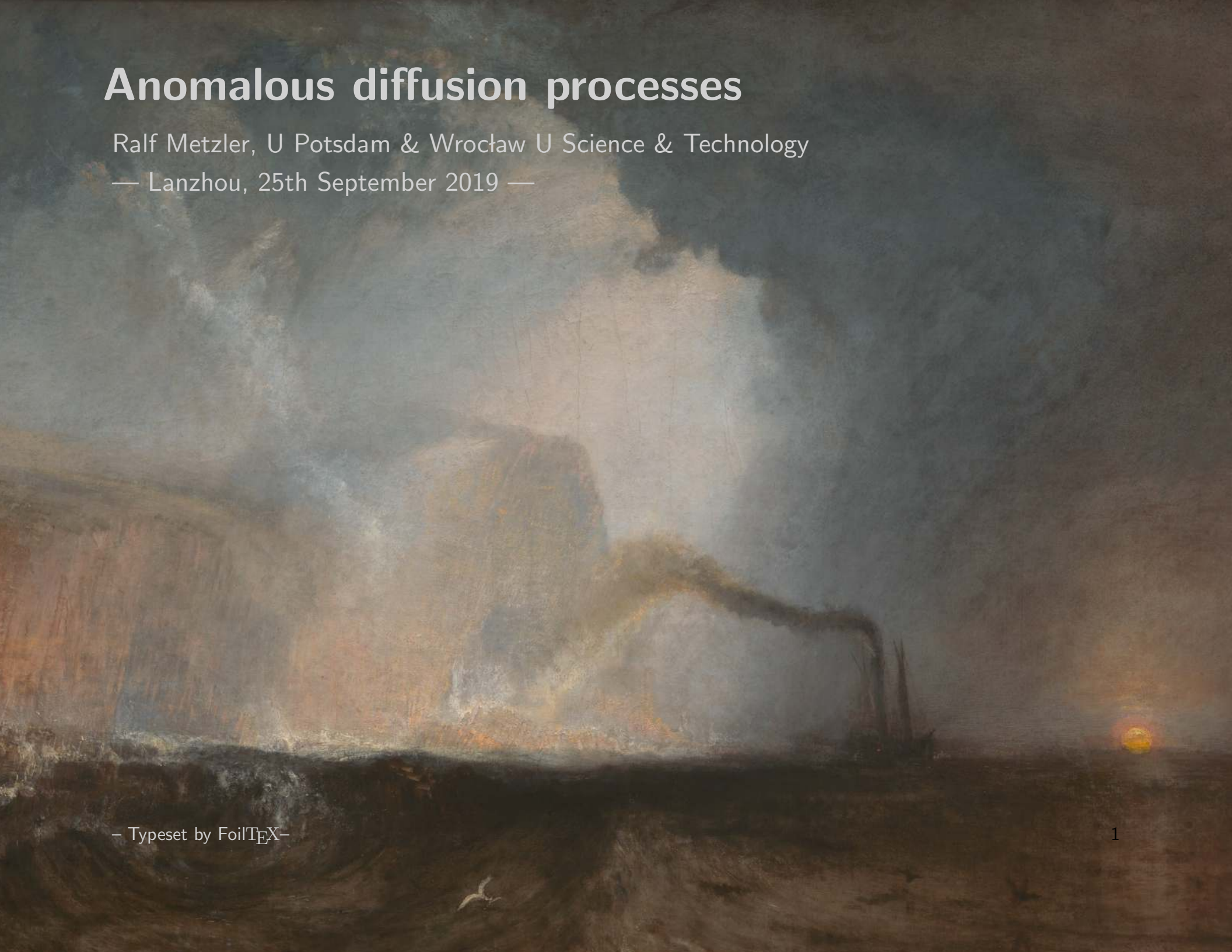


Anomalous diffusion processes

Ralf Metzler, U Potsdam & Wrocław U Science & Technology

— Lanzhou, 25th September 2019 —



Relaxation processes

Deborah number:

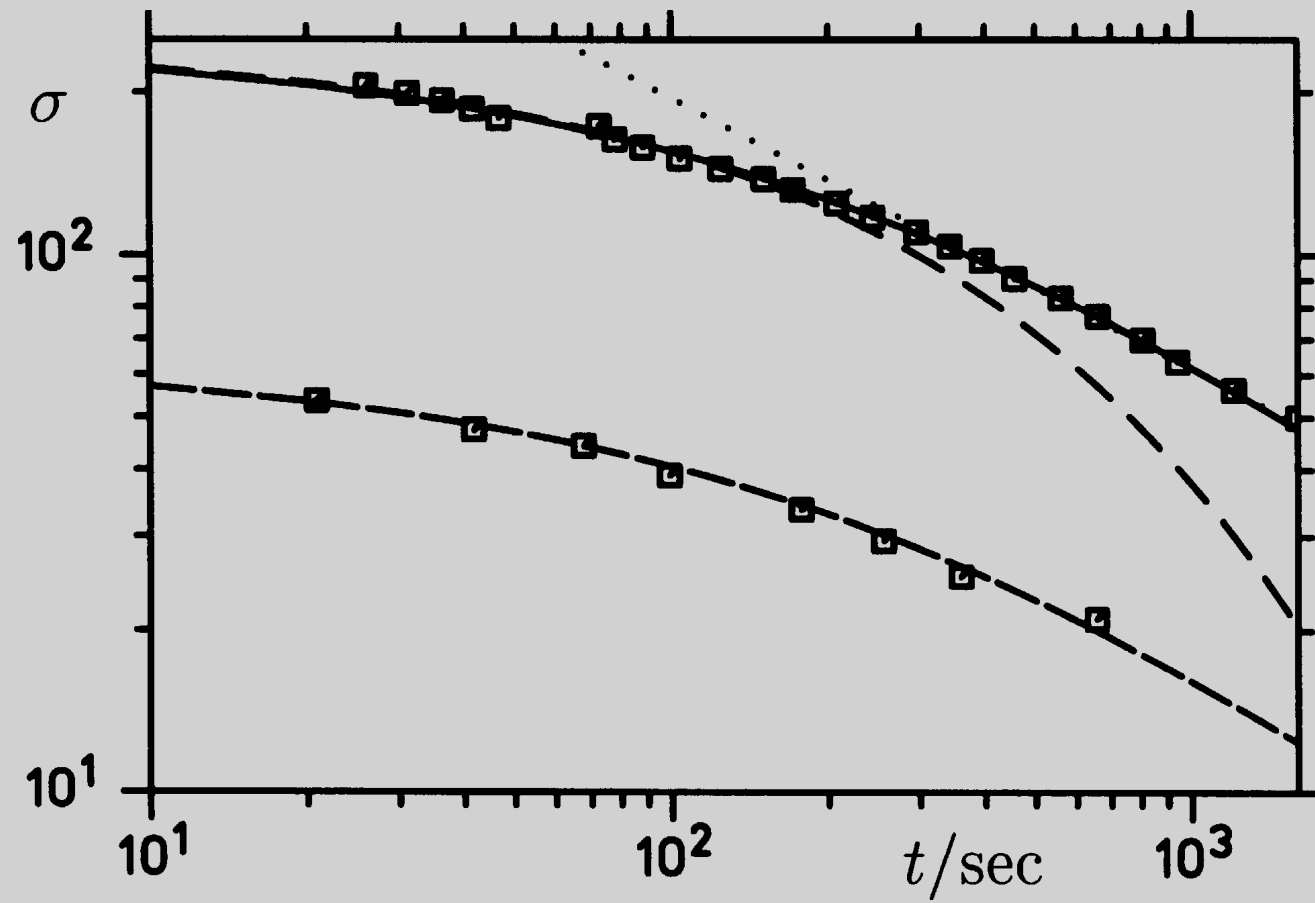
$$\tau = \frac{\text{Relaxation time scale}}{\text{Experimental time}}$$

The tar pitch drop experiment (U Queensland):

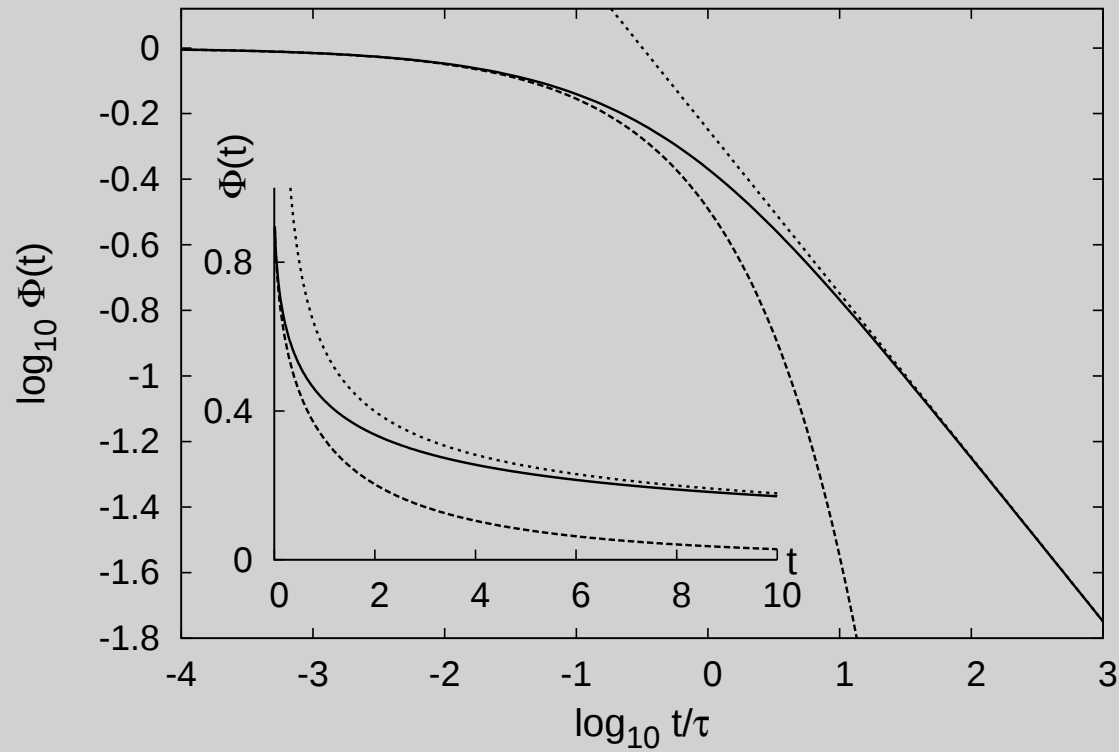
Date	Event
1927	Setup experiment
1930	Stem cut
12/1938	1st drop fell
02/1947	2nd drop fell
04/1954	3rd drop fell
05/1962	4th drop fell
08/1970	5th drop fell
04/1979	6th drop fell
07/1988	7th drop fell
11/2000	8th drop fell



Mittag-Leffler relaxation in rheology



Mittag-Leffler relaxation



$$E_{\alpha} \left(-(t/\tau)^{\alpha} \right) = \sum_{n=0}^{\infty} \frac{(-1)^n}{\Gamma(1 + \alpha n)} \left(\frac{t}{\tau} \right)^{\alpha n} \sim - \sum_{n=1}^{\infty} \frac{(-1)^n}{\Gamma(1 - \alpha n)} \left(\frac{\tau}{t} \right)^{\alpha n}$$

$$\frac{d\phi(t)}{dt} = {}_0D_t^{1-\alpha} \phi(t) \equiv \frac{d}{dt} \int_0^t \frac{(t-t')^{\alpha-1}}{\Gamma(\alpha)} dt'$$

Agricultural Rheology



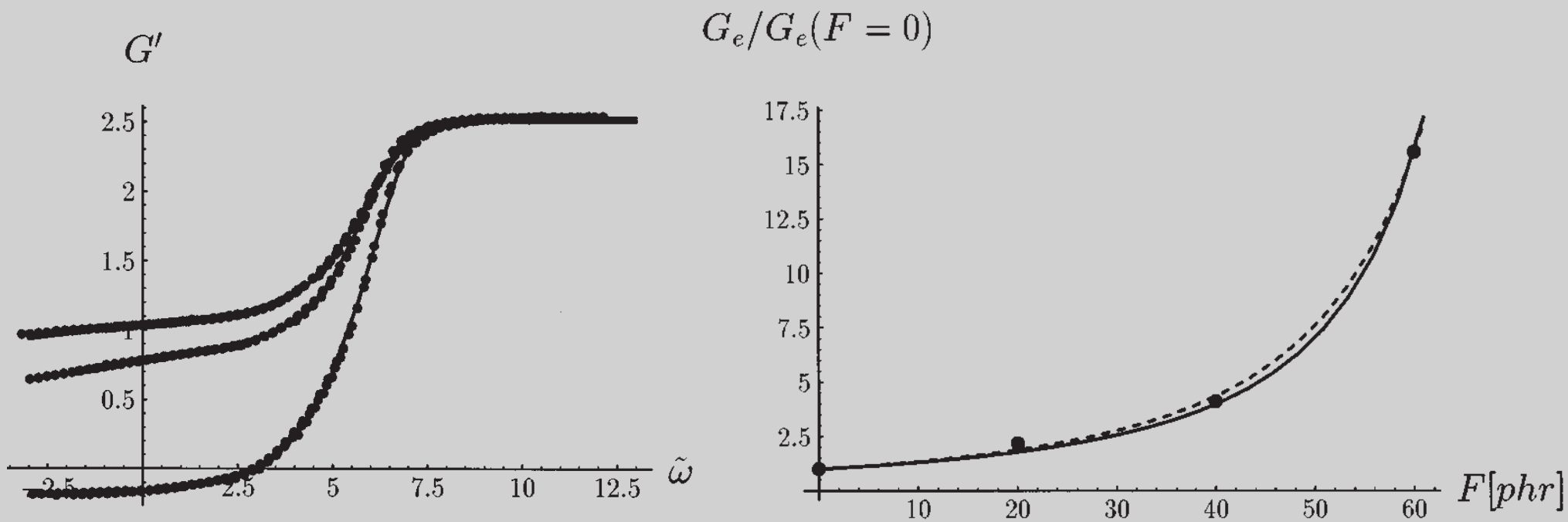
G. W. SCOTT BLAIR
AND MARKUS REINER



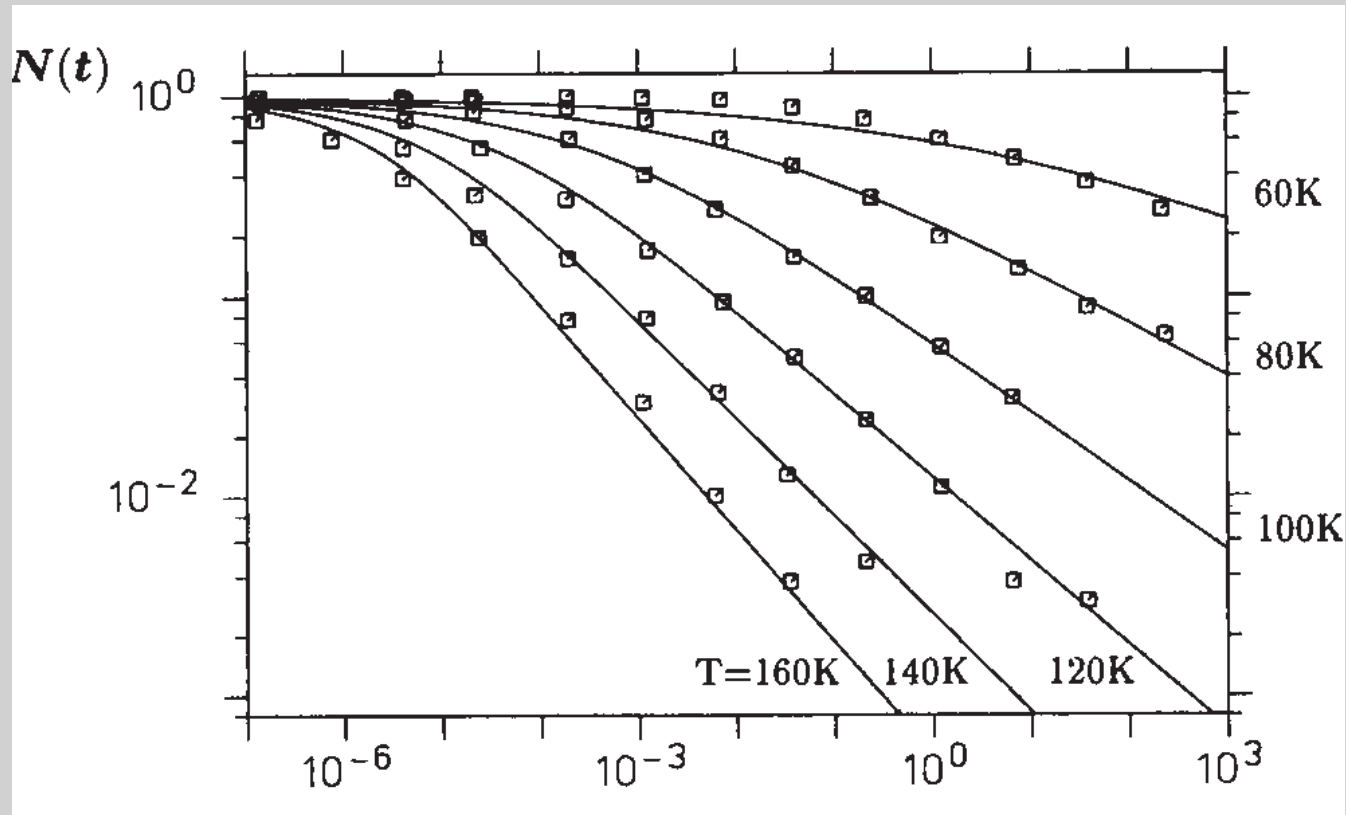
Moduli of carbon-black filled rubbers

Fractional Zener model of rheology:

$$\sigma(t) + \tau_0^q {}_0D_t^q (\sigma(t) - \sigma_0) = \left(G_m + G_e + G_\infty \right) {}_0D_t^q \tau_0^q (\epsilon(t) - \epsilon_0) + G_e \tau_0^{q-\mu} {}_0D_t^{q-\mu} \epsilon(t) + G_\infty \epsilon(t)$$



Rebinding of CO to myoglobin

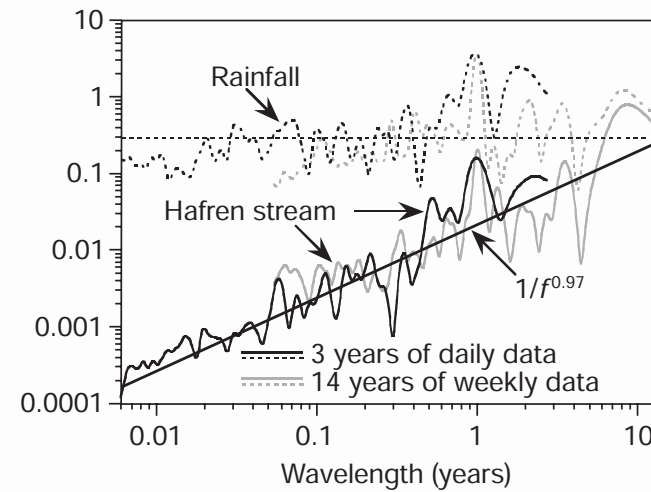
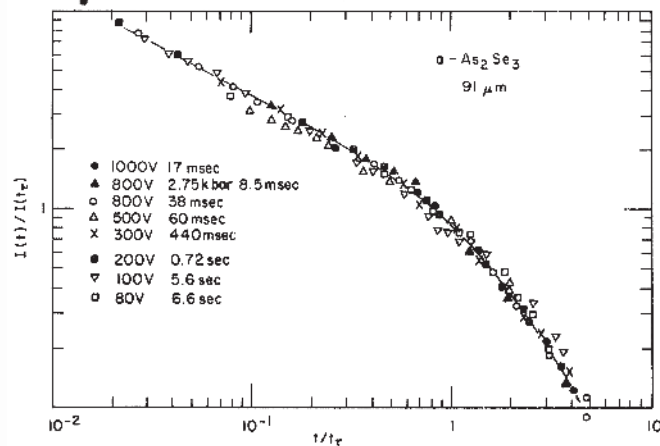


Anomalous diffusion occurs from micro to macro scales

Glasses/disordered systems

Tracer dispersion in groundwater

Scher & Montroll, PRB (1975)



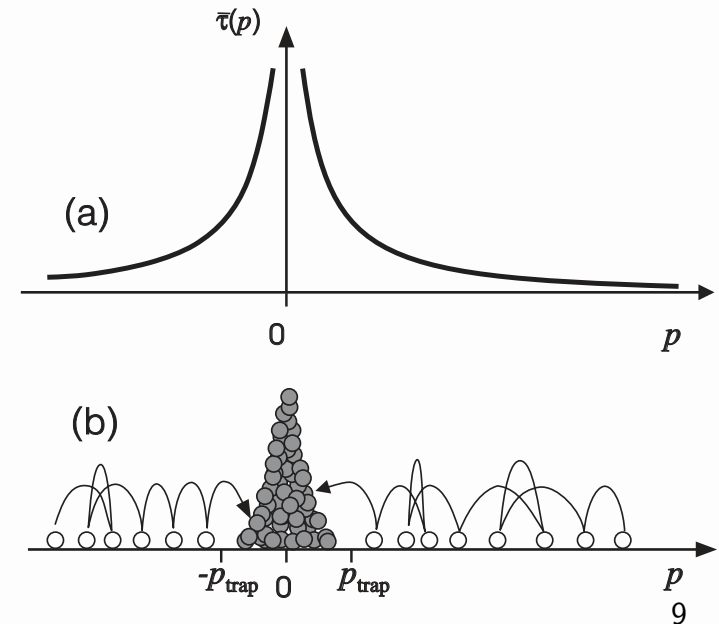
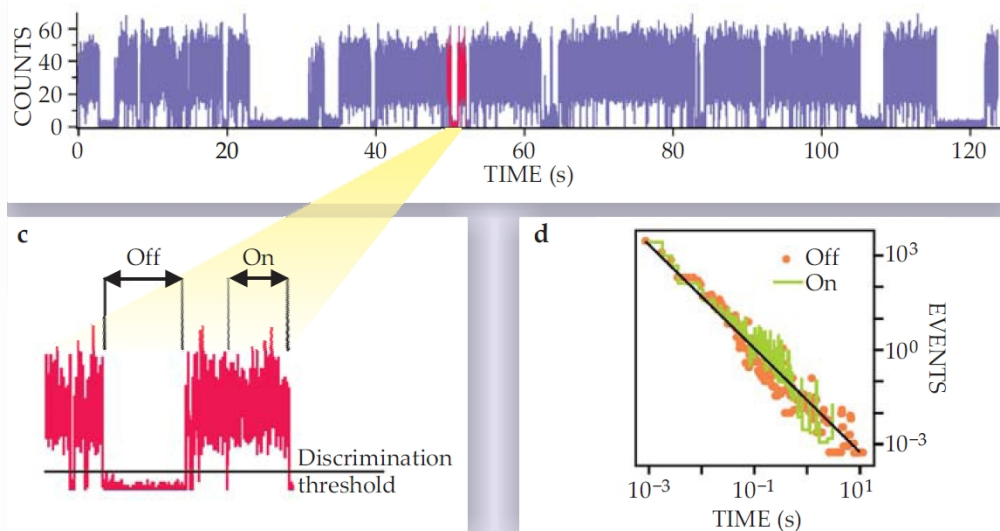
Scher et al, GRL (2002)

Blinking quantum dots

$$\langle \mathbf{r}^2(t) \rangle \simeq t^\alpha$$

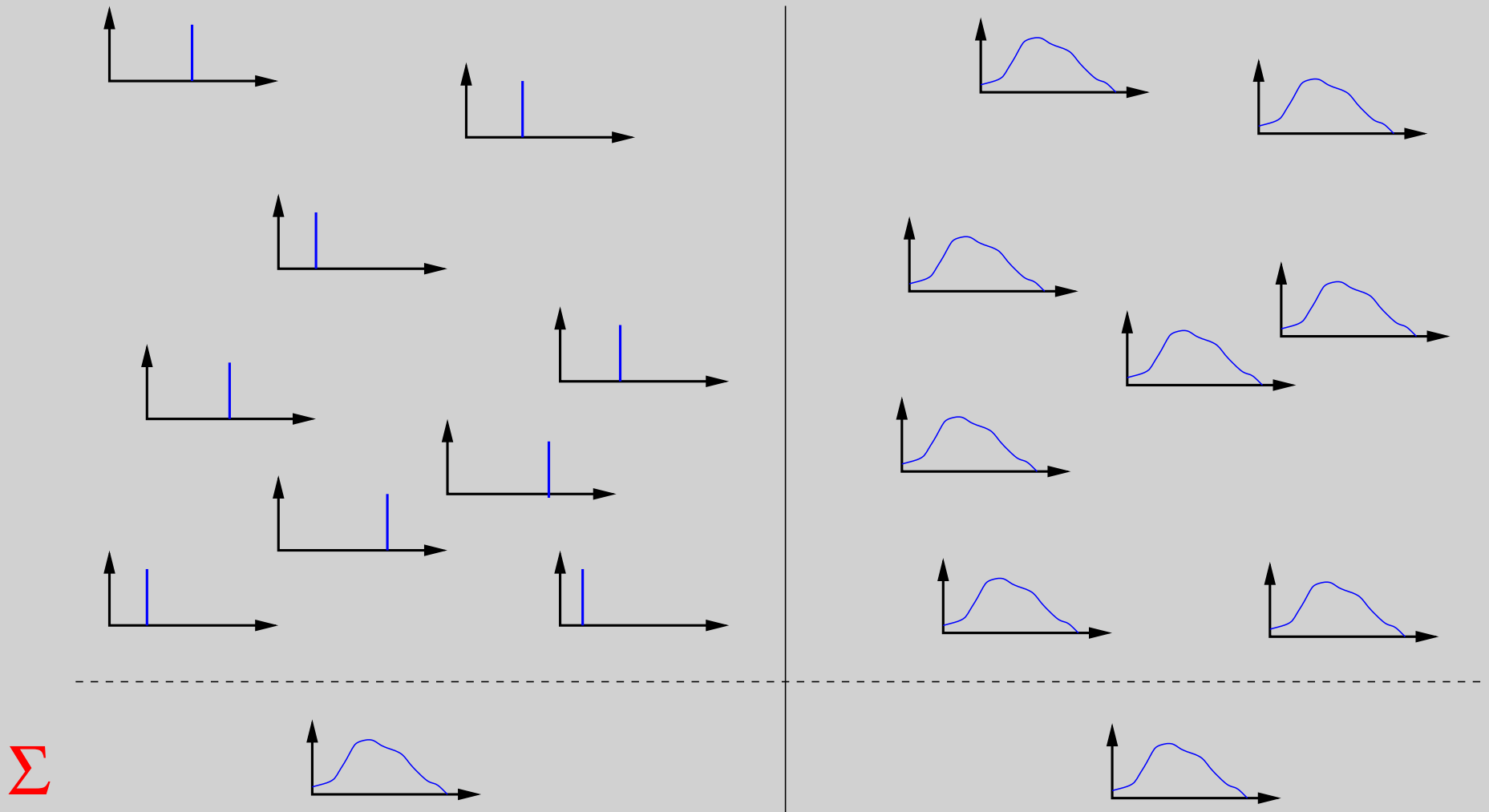
Subrecoil laser cooling

Stefani et al, Phys Today (2009)



Bardou, Chaotic Dynamics etc (2005)

Single molecule experiments vs ensemble average



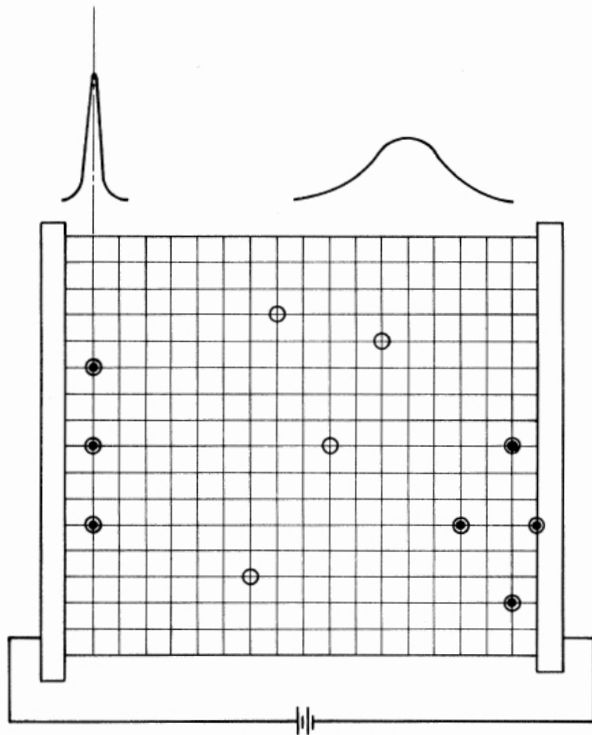
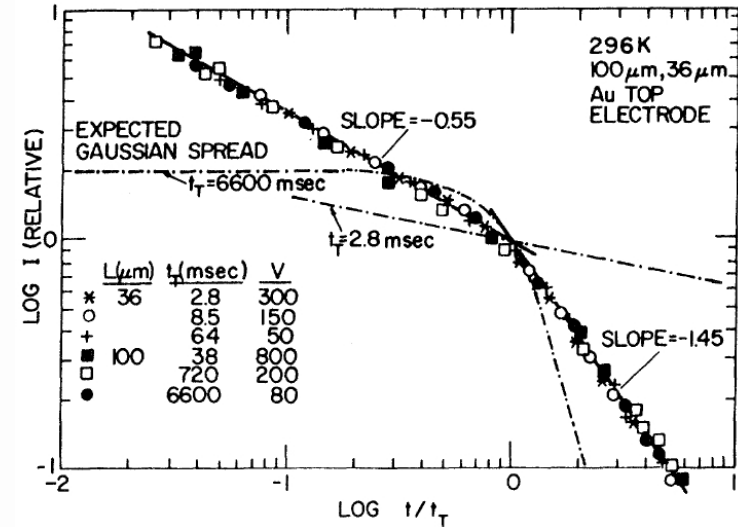
Charge carrier diffusion in amorphous semiconductors

Waiting time distribution

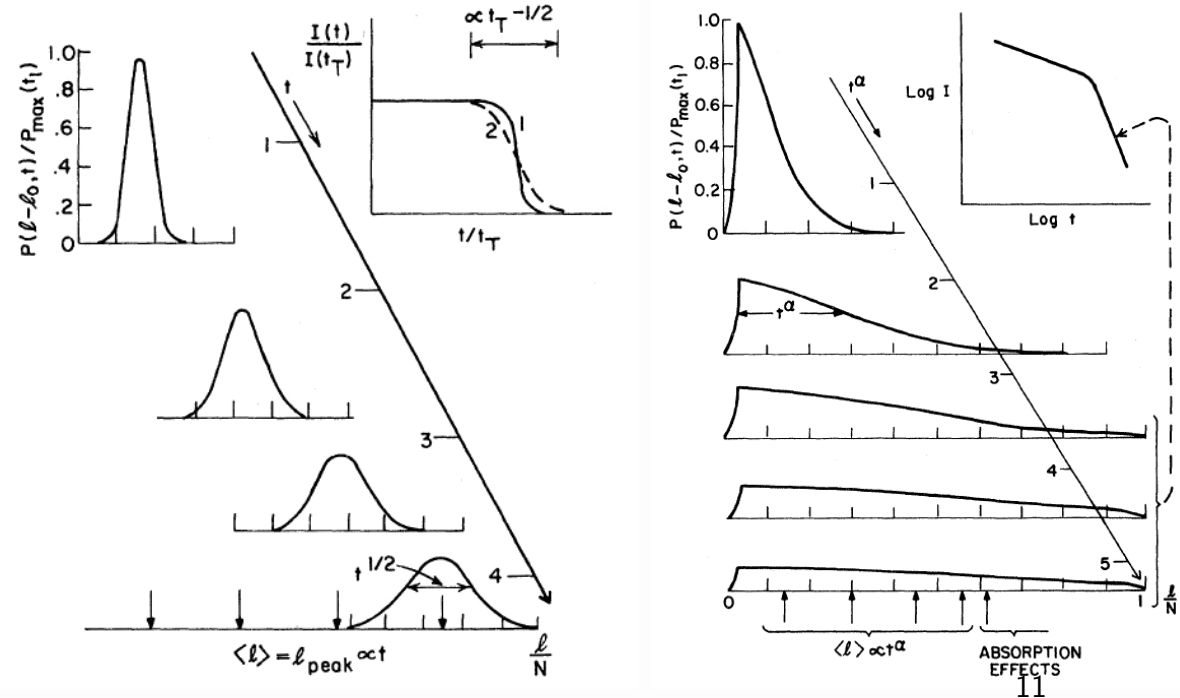
$$\psi(t) \simeq \frac{\tau^\alpha}{t^{1+\alpha}} \Leftrightarrow \psi(u) = \exp(-[u\tau]^\alpha)$$

In the bias of the electrical field:

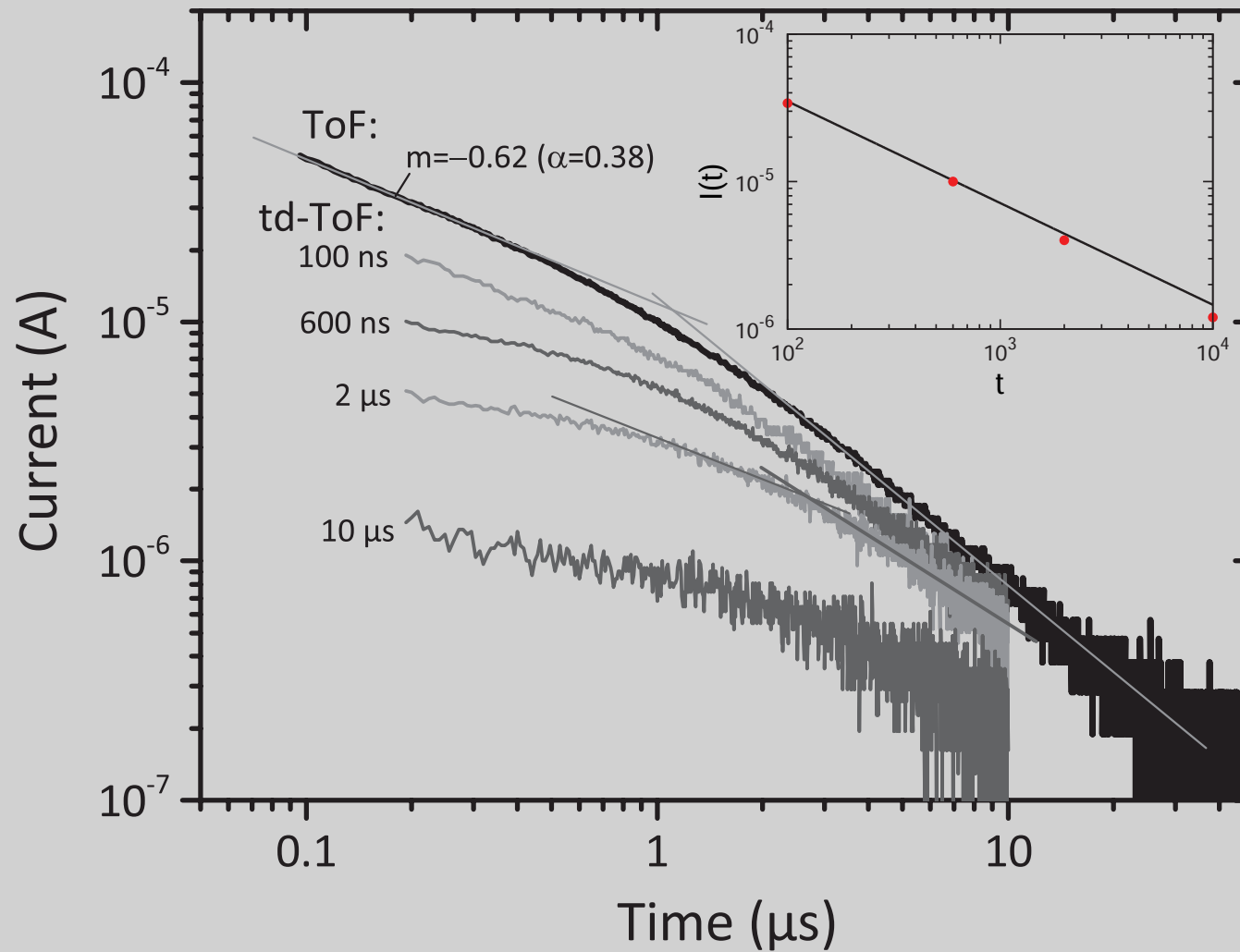
$$\frac{\sqrt{\langle (\Delta x)^2 \rangle}}{\langle x(t) \rangle} \simeq t^0 \quad \longleftrightarrow \quad \underbrace{\frac{\sqrt{\langle (\Delta x)^2 \rangle}}{\langle x(t) \rangle}}_{\alpha=1} \simeq t^{-1/2}$$



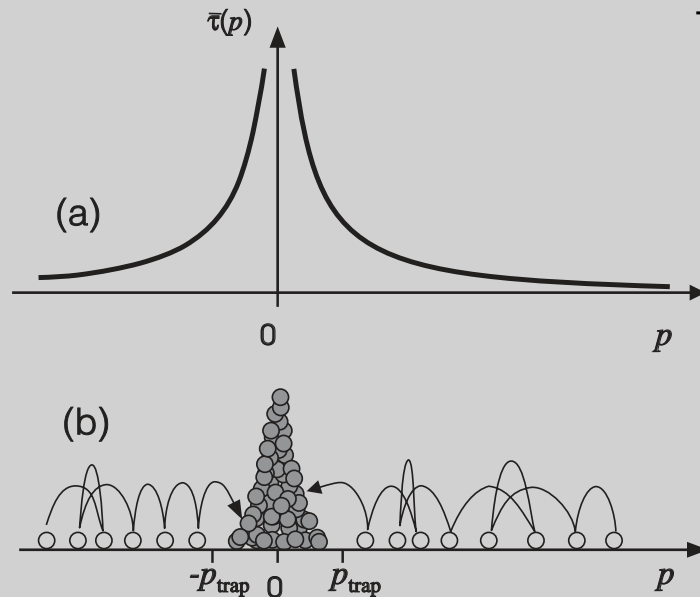
H Scher & E Montroll, PRB (1975)



Ageing charge carrier motion in polymeric semiconductors



Power-law waiting time dynamics in subrecoil laser cooling



Trapping time PDF of particles near zero momentum:

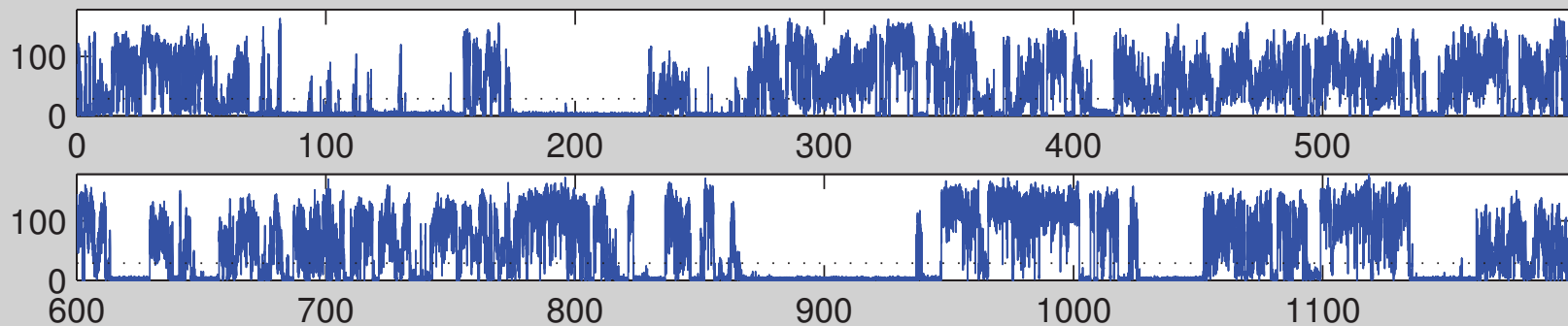
$$\psi(\tau) \simeq \tau^{-3/2}$$

F Bardou, JP Bouchaud, O Emile, A Aspect

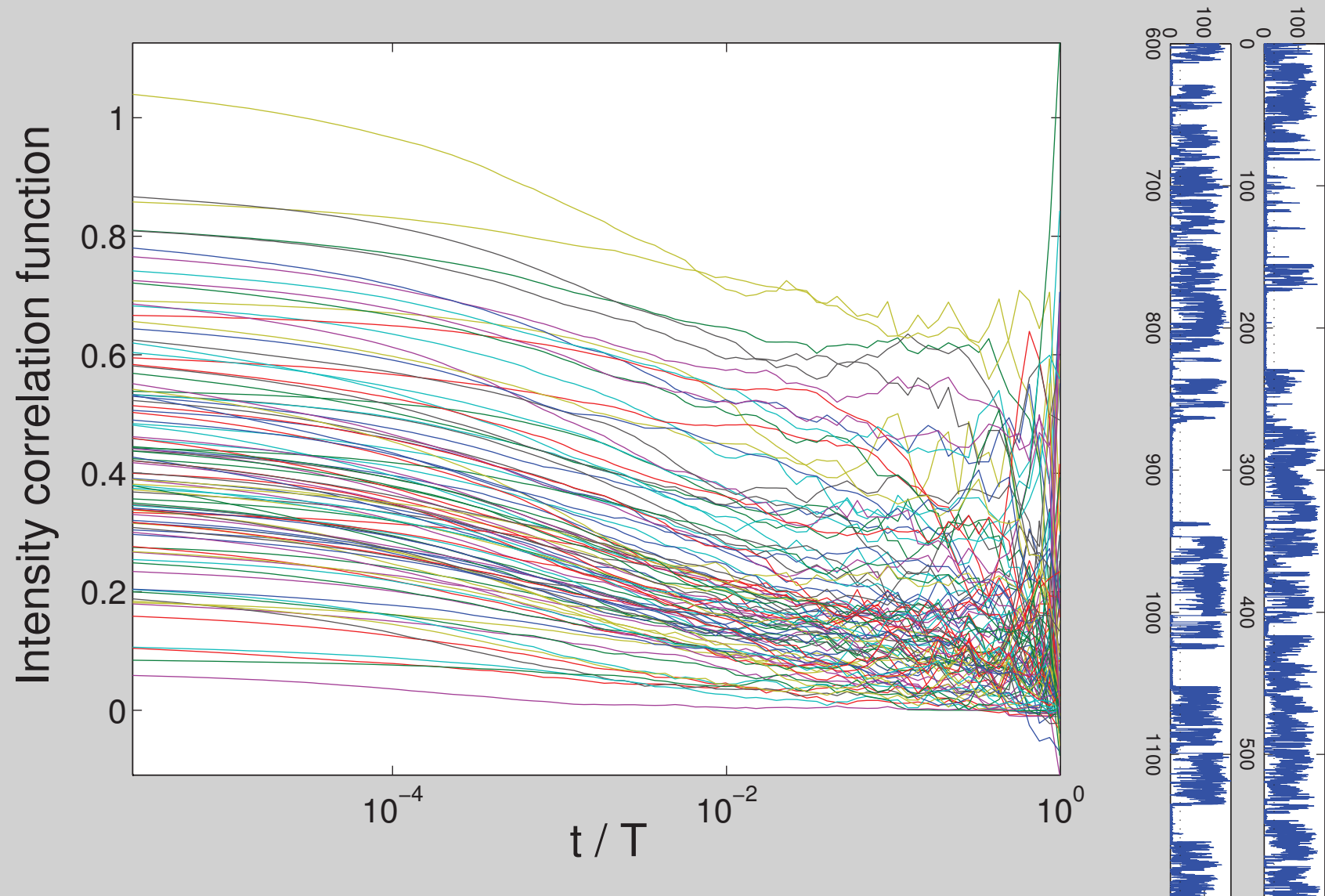
& C Cohen-Tannoudji, PRL (1994)

S Schaufler, WP Schleich, VP Yakovlev, PRL (1999)

Power-law switching times in quantum dots

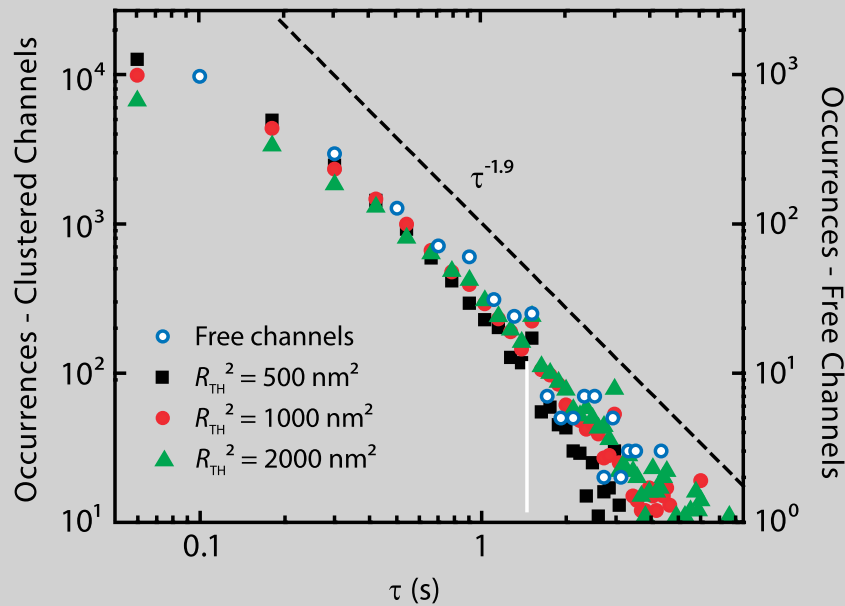
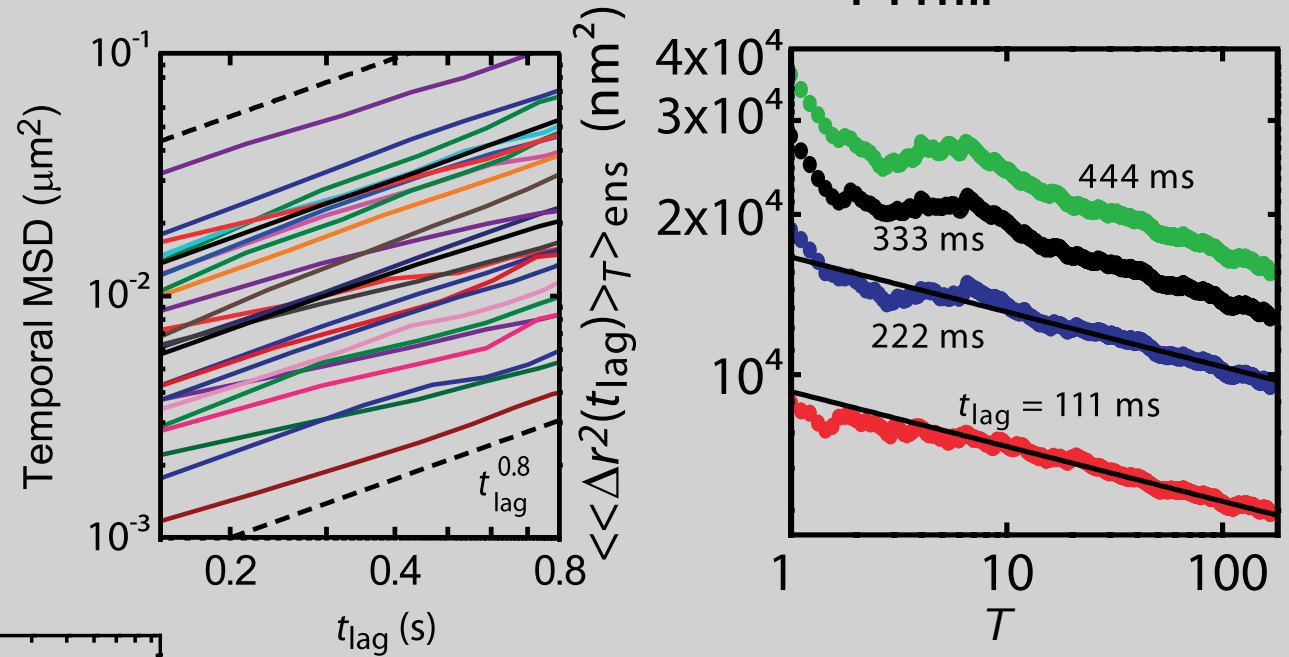
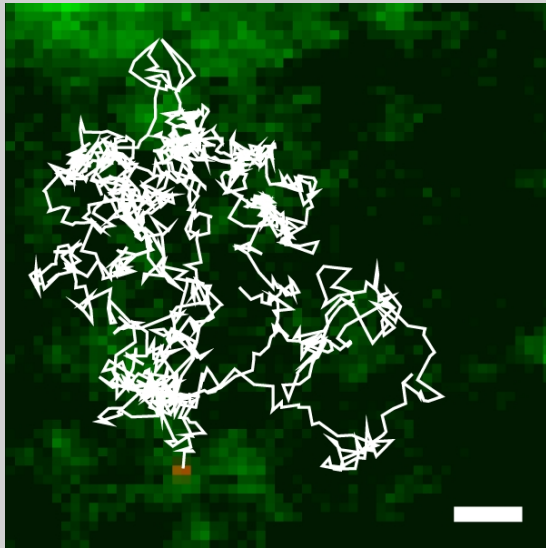


Blinking dynamics in quantum dots: amplitude scatter



On/off time statistics indeed follow power-laws with $\alpha \approx 0.5$

CTRW-like motion of Ka channels in plasma membrane



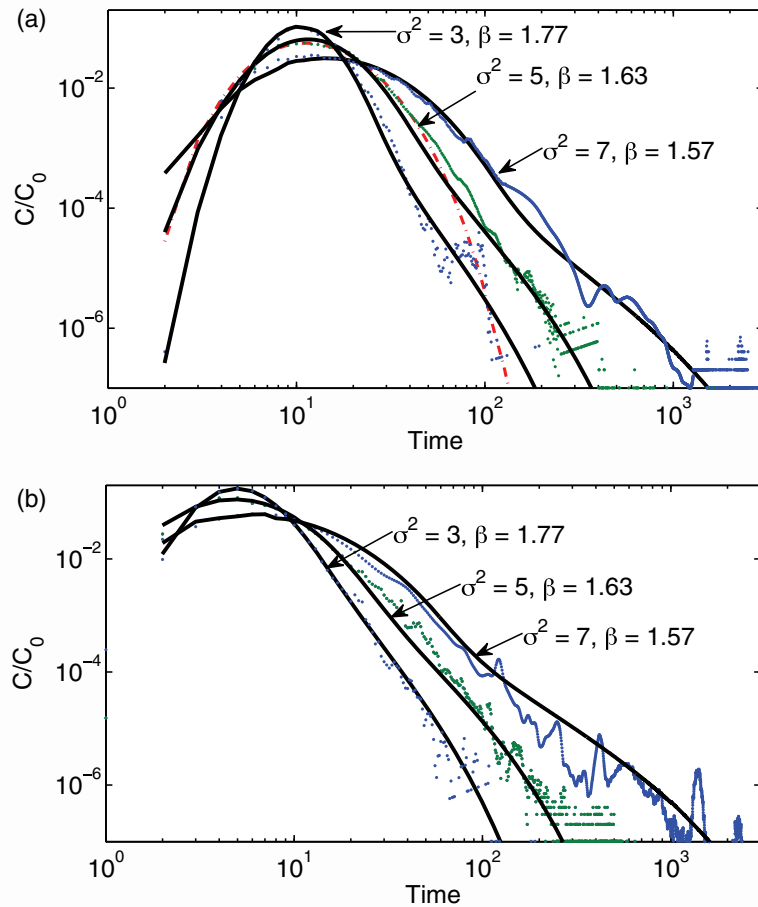
$$\psi(\tau) \simeq \tau^{-1-\alpha} \text{ scale free}$$

$$\overline{\delta^2(\Delta)} \text{ apparently random}$$

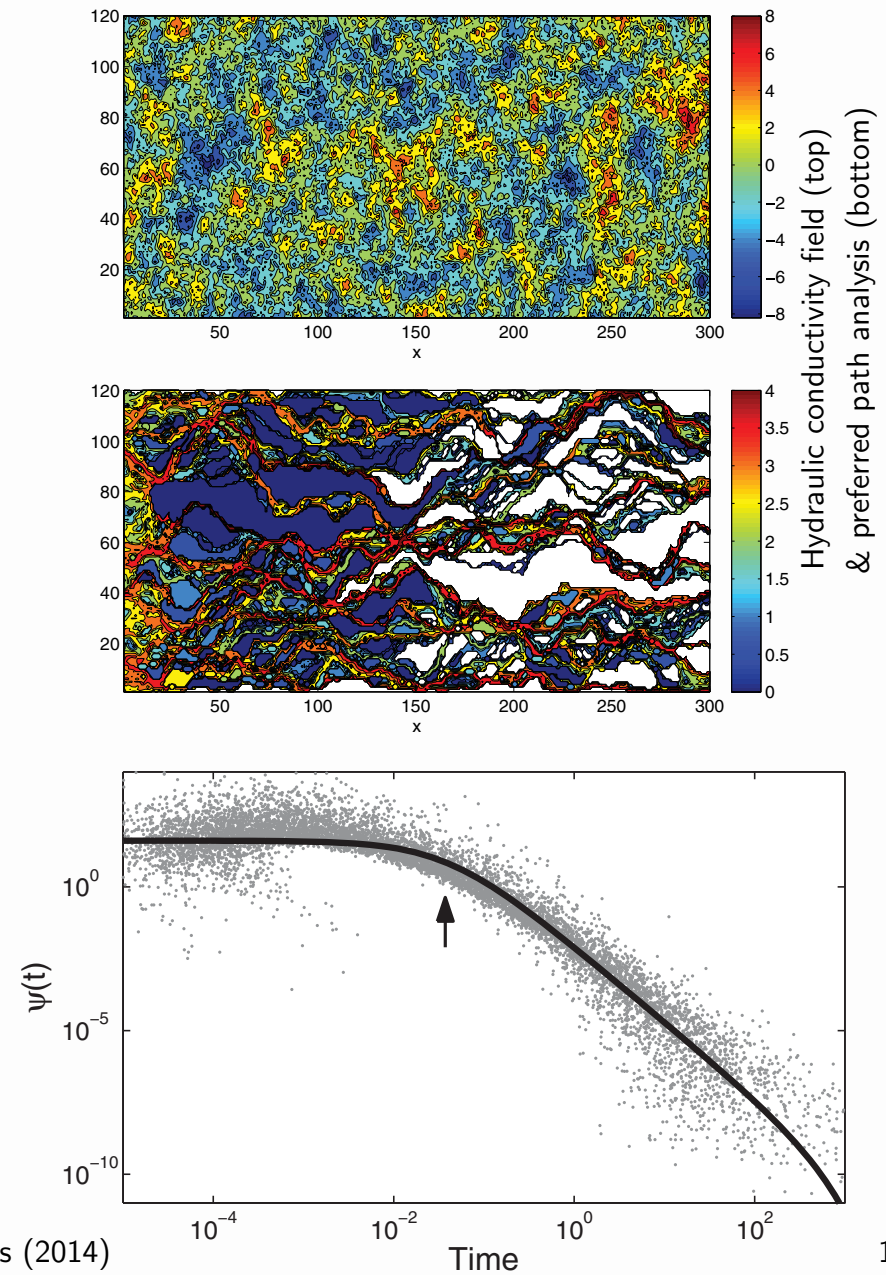
$$\Delta/T^{1-\alpha} \simeq \overline{\delta^2(\Delta)} \neq \langle \mathbf{r}^2(\Delta) \rangle \simeq \Delta^\alpha$$

$$P(\mathbf{r}, t) \simeq \exp \left(-\beta r^{1/[1-\alpha/2]} \right)$$

Reconstruction of waiting time PDF from hydraulic field



Reconstructed $\psi(t)$ fitted by parameters from the corresponding breakthrough curves





Waiting time distribution $\psi(t)$; characteristic waiting time $\langle t \rangle = \int_0^\infty t\psi(t)dt \rightarrow \infty$

[illegible]

Scher-Montroll equation

$$\eta(x, t) = \delta(t)\delta(x) + \int_0^t dt' \int_{-\infty}^{\infty} dx' \eta(x', t') \lambda(x - x') \psi(t - t')$$

$$\eta(k, u) = 1 + \eta(k, u) \lambda(k) \psi(u)$$

$$P(x, t) = \int_0^t dt' \eta(x, t') \Psi(t - t') \quad \therefore \quad \Psi(t) = 1 - \int_0^t dt' \psi(t')$$

$$P(k, u) = \eta(k, u) \Psi(u) = \eta(k, u) \frac{1 - \psi(u)}{u}$$

$$P(k, u) = \frac{1 - \psi(u)}{u} \frac{1}{1 - \lambda(k) \psi(u)} \quad \therefore \quad \psi(u) \sim 1 - (u\tau)^\alpha, \quad \lambda(k) \sim 1 - (\sigma|k|)^\mu$$

A primer on the fractional Fokker-Planck equation

Waiting time and jump length densities:

$$\psi(t) \simeq \frac{\tau^\alpha}{t^{1+\alpha}}, \quad 0 < \alpha < 1; \quad \Lambda(x, x') = \lambda(x-x') \left(A(x')\Theta(x-x') + B(x')\Theta(x'-x) \right)$$

Continuum limit: fractional Fokker-Planck equation:

$$\frac{\partial}{\partial t} P(x, t) = {}_0D_t^{1-\alpha} \left(\frac{\partial}{\partial x} \frac{V'(x)}{m\eta_\alpha} + K_\alpha \frac{\partial^2}{\partial x^2} \right) P(x, t) \quad \therefore \quad K_\alpha = \frac{k_B T}{m\eta_\alpha}$$

Riemann-Liouville operator:

Subordination:

$${}_0D_t^{1-\alpha} f(t) = \frac{1}{\Gamma(\alpha)} \frac{\partial}{\partial t} \int_0^t \frac{f(t')}{(t-t')^{1-\alpha}} dt' \quad \therefore \quad P(x, t) = \int_0^\infty \mathcal{E}_\alpha(s, t) P_M(x, s) ds$$

Generalised Laplace kernel ($s^* \equiv \eta_\alpha s / \eta_1$):

$$\mathcal{E}_\alpha(s, t) = \mathcal{L}^{-1} \left\{ \frac{\eta_\alpha}{\eta_1 u^{1-\alpha}} \exp(-u^\alpha s^*) \right\} = \frac{1}{\alpha s} H_{1,1}^{1,0} \left[\frac{(s^*)^{1/\alpha}}{t} \middle| \begin{matrix} (1, 1) \\ (1, 1/\alpha) \end{matrix} \right]$$

Mittag-Leffler form of the mode relaxation:

$$T_n(t) = T_0 E_\alpha \left(-\lambda_n t^\alpha \right) = \sum \frac{(-\lambda_n t^\alpha)^n}{\Gamma(1 + \alpha n)} \sim \begin{cases} 1 - \lambda_n t^\alpha / \Gamma(1 + \alpha) \\ 1 / (\Gamma(1 - \alpha) \lambda_n t^\alpha) \end{cases}$$

Fractional diffusion equation

Time-fractional diffusion equation:

$$\frac{\partial}{\partial t} P(x, t) = {}_0D_t^{1-\alpha} K_\alpha \frac{\partial^2}{\partial x^2} P(x, t)$$

Equivalent version:

$${}_0D_t^\alpha P(x, t) - \frac{1}{\Gamma(1-\alpha)} t^{-\alpha} = \frac{\partial^2}{\partial x^2} P(x, t)$$

Initial value also needed for consistency: after integration $\int \cdot dx$,

$$\text{lhs: } {}_0D_t^\alpha 1 - \frac{1}{\Gamma(1-\alpha)} t^{-\alpha} = 0$$

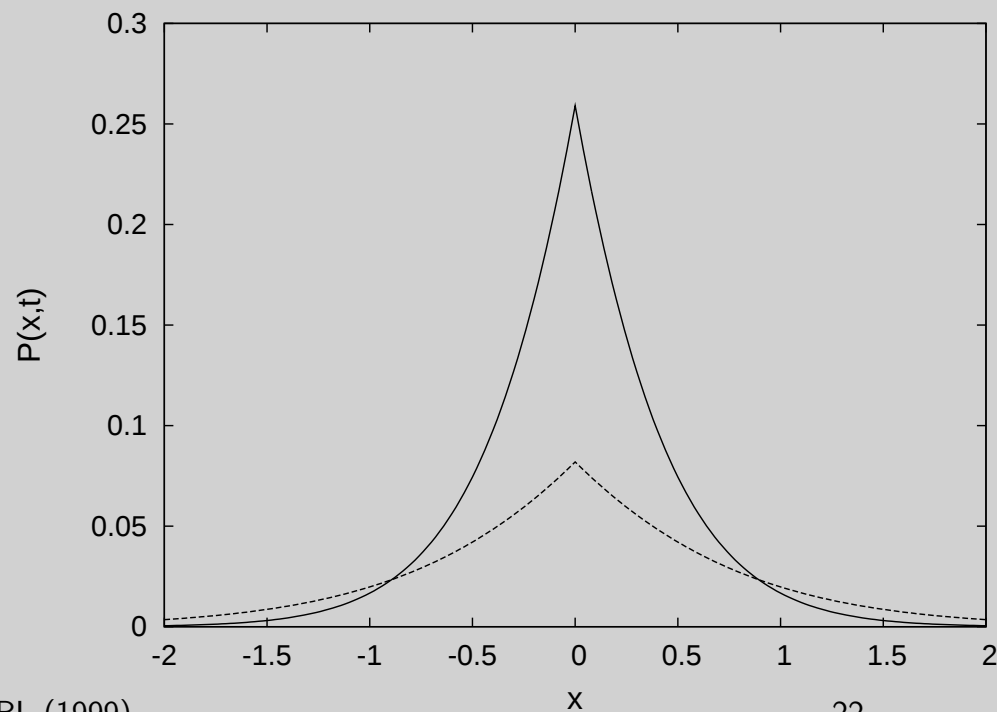
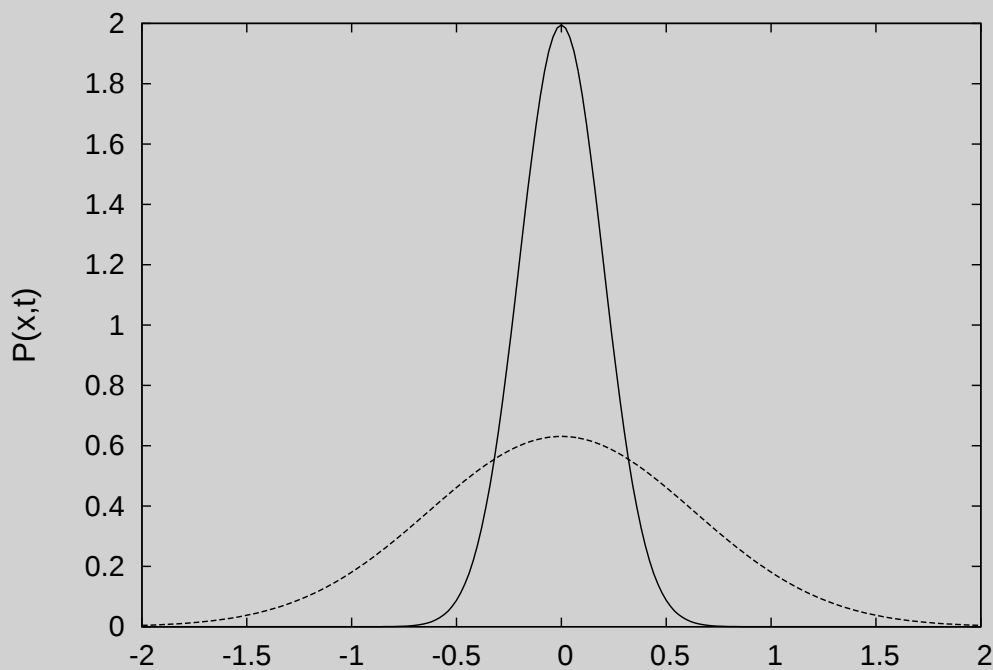
$$\text{rhs: } K_\alpha [\partial P(x, t) / \partial x]_{-\infty}^{\infty} = 0$$

Alternatively, use Caputo operator

Solution of the fractional diffusion equation

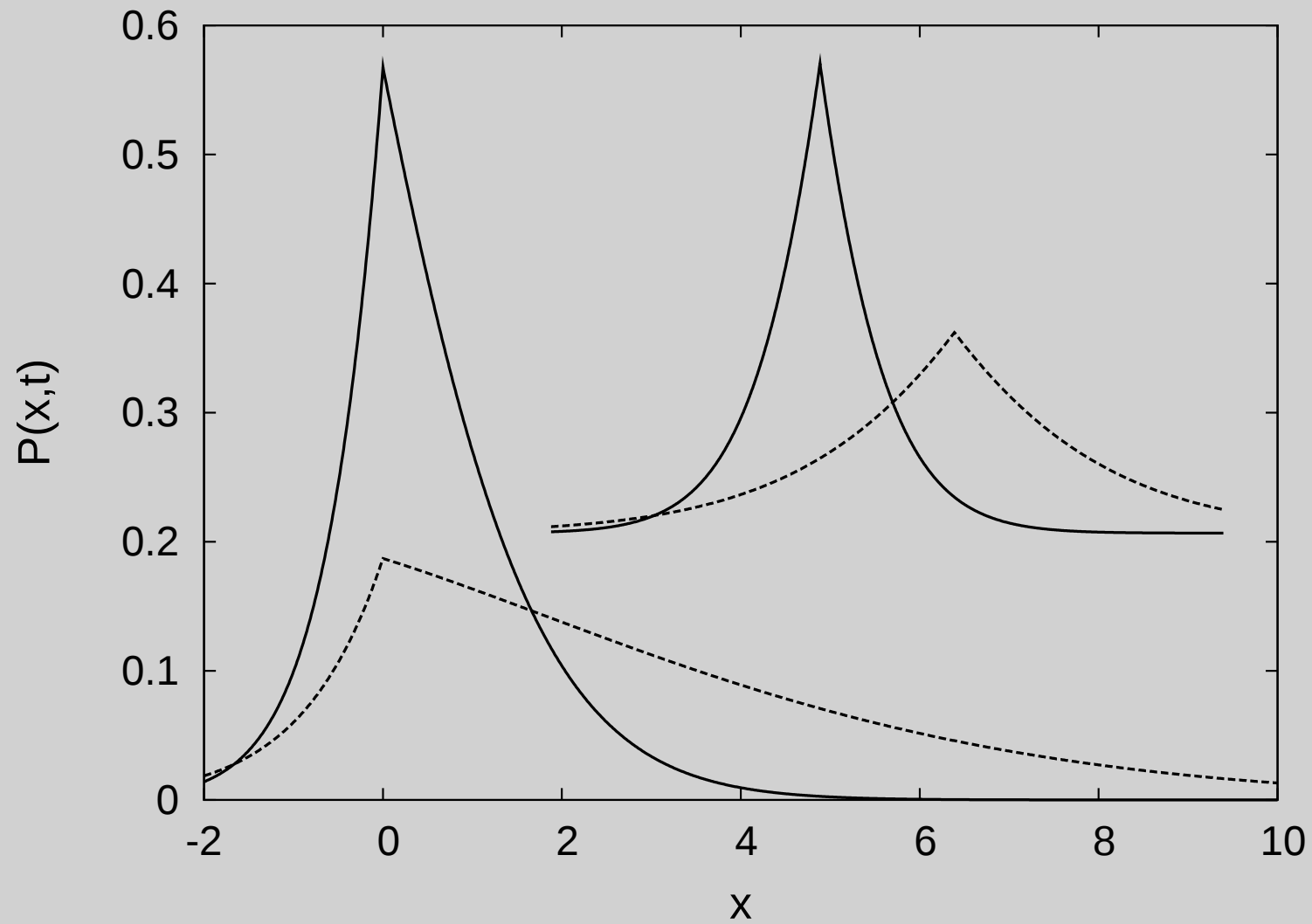
$$P(x, t) = \frac{1}{\sqrt{4\pi K_\alpha t^\alpha}} H_{1,2}^{2,0} \left[\frac{x^2}{4K_\alpha t^\alpha} \middle| \begin{matrix} (1 - \frac{\alpha}{2}, \alpha) \\ (0, 1), (\frac{1}{2}, 1) \end{matrix} \right]$$

$$P(x, t) \sim \frac{1}{\sqrt{4\pi K_\alpha t^\alpha}} \sqrt{\frac{1}{2-\alpha}} \left(\frac{2}{\alpha}\right)^{(1-\alpha)/(2-\alpha)} \left(\frac{|x|}{\sqrt{K_\alpha t^\alpha}}\right)^{-(1-\alpha)/(2-\alpha)} \\ \times \exp\left(-\frac{2-\alpha}{2} \left(\frac{\alpha}{2}\right)^{\alpha/(2-\alpha)} \left[\frac{|x|}{\sqrt{K_\alpha t^\alpha}}\right]^{1/(1-\alpha/2)}\right)$$



RM & J Klafter, Phys Rep (2000); RM, E Barkai & J Klafter, PRL (1999)

Fractional diffusion-advection



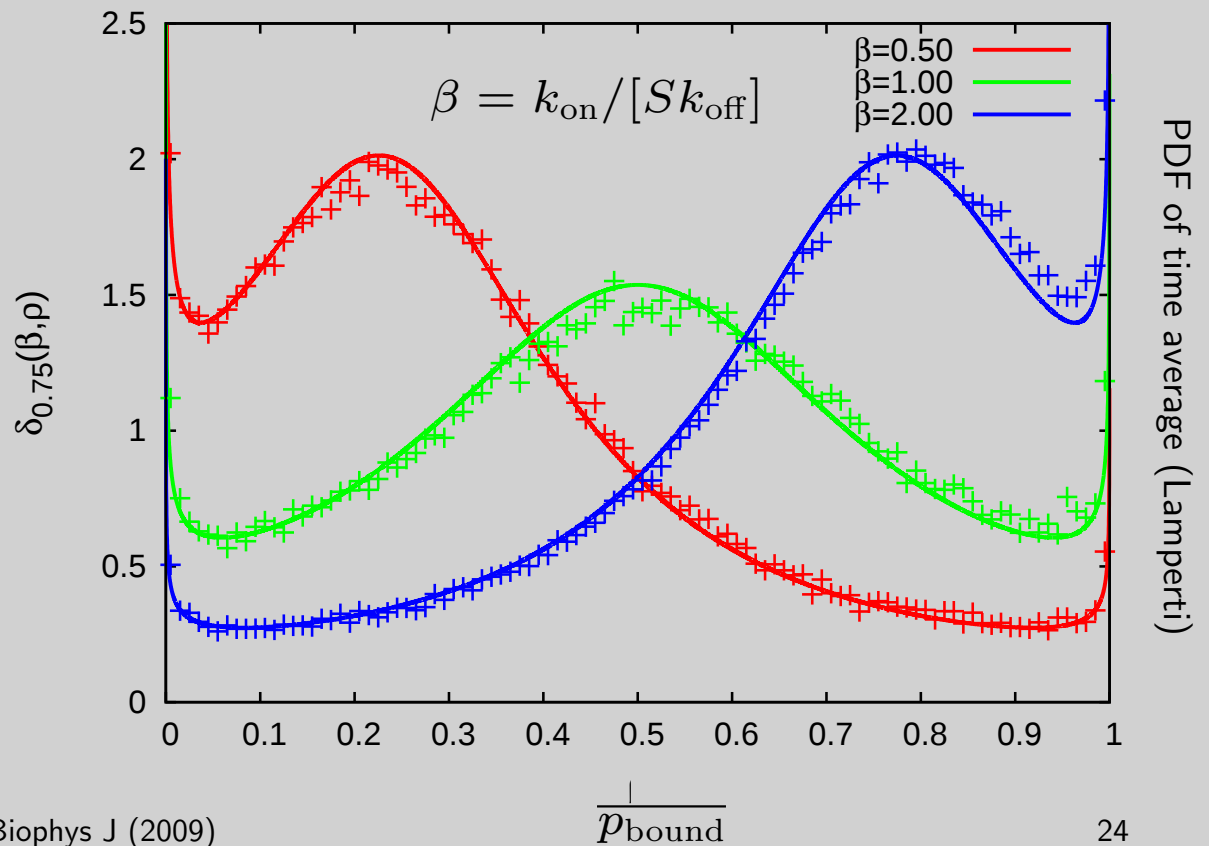
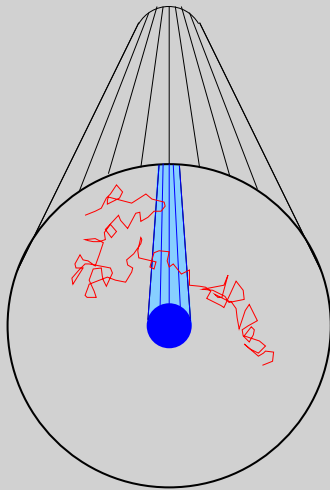
Lack of time scale causes weak ergodicity breaking (WEB)

Unbinding time PDF to bulk, rebinding time PDF to cylinder for $\psi(t) \simeq t^{-1-\alpha}$:

$$\wp_{\text{unb}}(t) \sim \frac{\kappa^{1-\alpha}}{\kappa_{\text{off}} t^{1+\alpha}}; \quad \wp_{\text{reb}}(t) \sim \frac{S}{k_{\text{on}} t^{1+\alpha}}, \quad S = \pi(R_2^2 - R_1^2)$$

In this system time averages are random & different from ensemble averages: WEB

$$\overline{p_{\text{bound}}} \equiv \lim_{t \rightarrow \infty} \frac{t_{\text{bound}}}{t} \neq \langle p_{\text{bound}} \rangle \equiv \frac{k_{\text{on}}}{k_{\text{on}} + S k_{\text{off}} / k_{\text{on}}}$$



Extracting information from single Brownian trajectories



Ensemble averaged MSD for normal diffusion (on average # jumps $\sim$ elapsed time t):

$$\langle \mathbf{r}^2(t) \rangle = \int \mathbf{r}^2 P(\mathbf{r}, t) d\mathbf{r} = 2dK_1 t \quad \left(= \langle \delta \mathbf{r}^2 \rangle \frac{t}{\tau}, \quad K_1 = \frac{\langle \delta \mathbf{r}^2 \rangle}{2d\tau} \right)$$

Single particle trajectory $\mathbf{r}(t)$, $t \in [0, T]$:

$$\overline{\delta^2(t)} = \frac{1}{T-t} \int_0^{T-t} [\mathbf{r}(t' + t) - \mathbf{r}(t')]^2 dt' = \frac{1}{T-t} \int_0^{T-t} \langle \delta \mathbf{r}^2 \rangle \frac{t}{\tau} dt'$$

Single trajectory information equals ensemble information (*Boltzmann-Khinchin*):

$$\lim_{T \rightarrow \infty} \overline{\delta^2(t)} = 2dK_1 t = \langle \mathbf{r}^2(t) \rangle$$

Anomalous diffusion is not always ergodic (weak or strong violation):

$$\langle \mathbf{r}^2(t) \rangle \simeq K_\alpha t^\alpha \neq \overline{\delta^2(t)} \simeq t/T^{1-\alpha}$$

WEB: signature of non-stationarity of process. SEB: discontinuity of phase space

Time averaged MSD & weak ergodicity breaking (WEB)

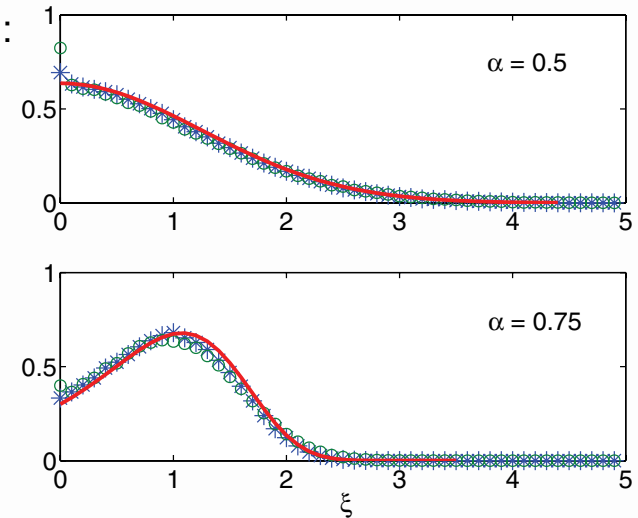
Time averaged MSD $\simeq \Delta$ is pseudo-Brownian and ageing ($\langle x^2(t) \rangle \simeq K_\alpha t^\alpha$):

$$\langle \overline{\delta^2(\Delta)} \rangle \sim \frac{1}{N} \sum_i^N \overline{\delta_i^2(\Delta)} \sim \frac{2dK_\alpha}{\Gamma(1+\alpha)} \frac{\Delta}{T^{1-\alpha}} \quad \therefore \quad K_\alpha \equiv \frac{\langle \delta \mathbf{r}^2 \rangle}{2\tau^\alpha}$$

Amplitude distribution $\overline{\delta^2}$ of trajectories ($\xi \equiv \overline{\delta^2} / \langle \overline{\delta^2} \rangle$):

$$\phi_\alpha(\xi) \sim \frac{\Gamma^{1/\alpha}(1+\alpha)}{\alpha \xi^{1+1/\alpha}} L_\alpha^+ \left(\frac{\Gamma^{1/\alpha}(1+\alpha)}{\xi^{1/\alpha}} \right)$$

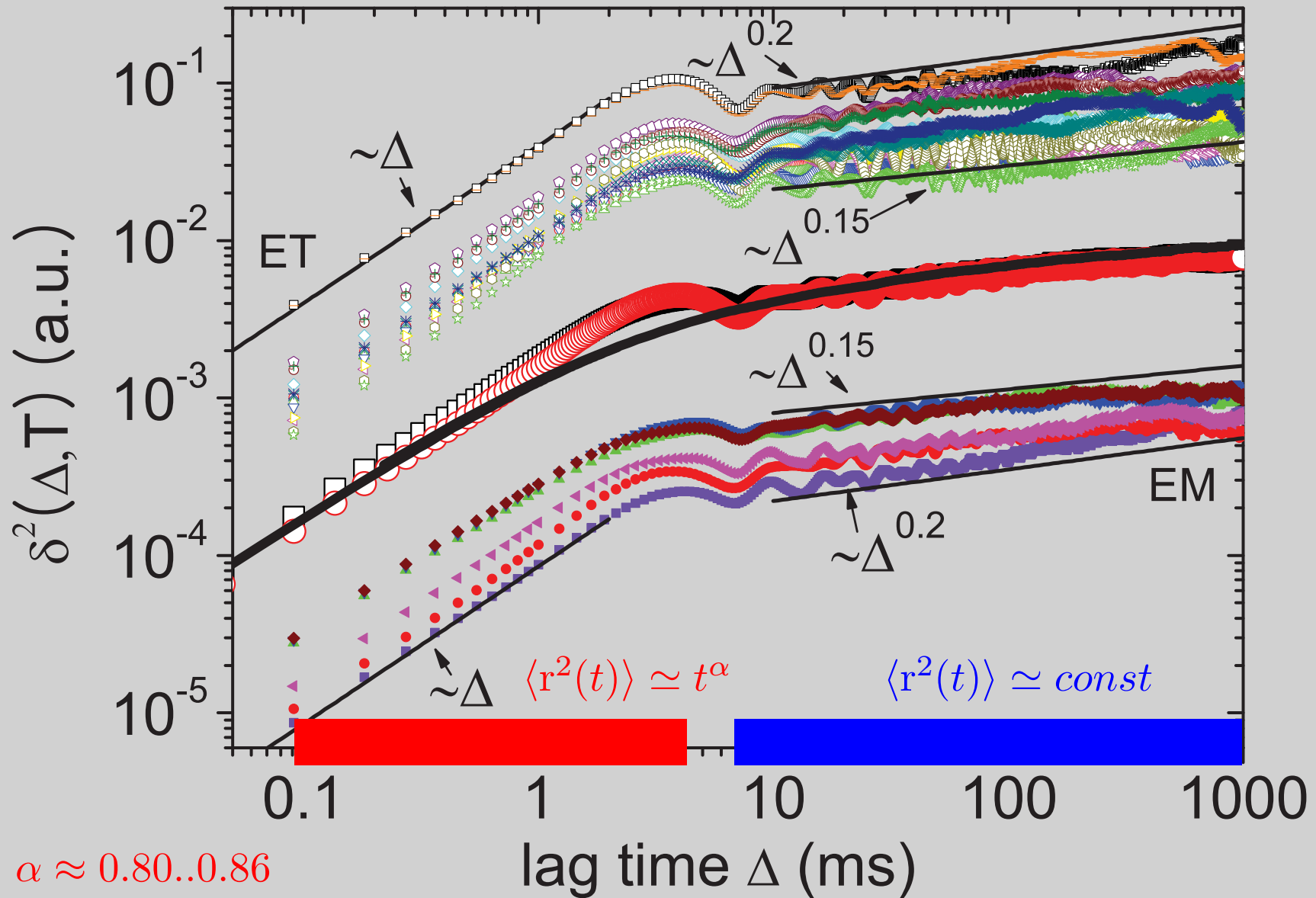
$$\phi_{1/2}(\xi) = \frac{2}{\pi} \exp \left(-\frac{\xi^2}{\pi} \right); \quad \phi_1(\xi) = \delta(\xi - 1)$$



Confinement does not effect a plateau ($\langle x^2(t) \rangle \simeq \text{const}(T)$):

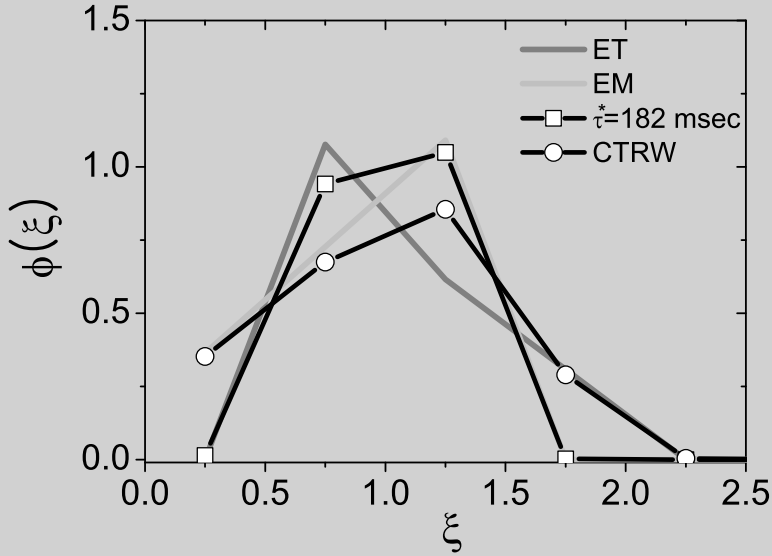
$$\langle \overline{\delta^2(\Delta)} \rangle \sim \left(\langle x^2 \rangle_B - \langle x \rangle_B^2 \right) \frac{2 \sin(\pi\alpha)}{(1-\alpha)\alpha\pi} \left(\frac{\Delta}{T} \right)^{1-\alpha}; \quad \frac{1}{(K_\alpha \lambda_1)^{1/\alpha}} \ll \Delta \ll T$$

Granule subdiffusion in harmonic optical tweezer potential

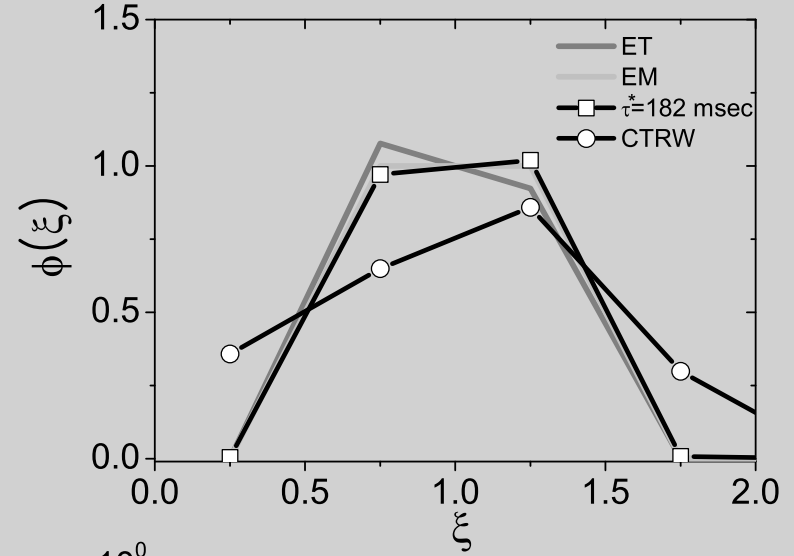


Further analysis of the lipid granule data

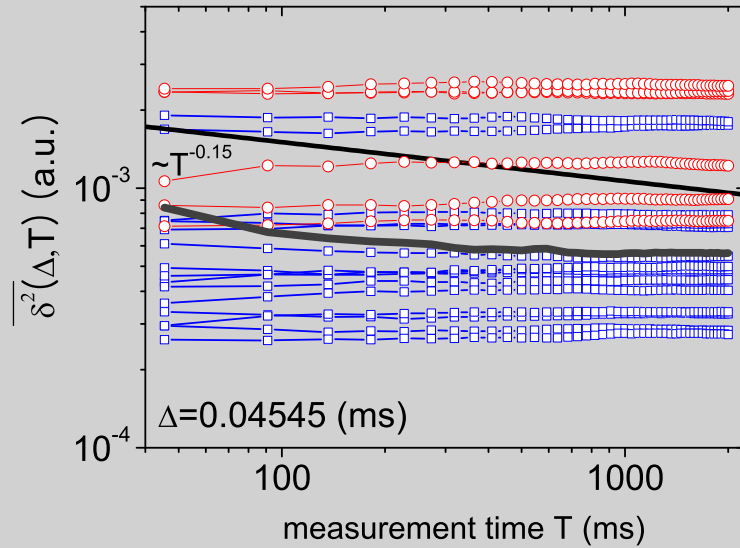
$\Delta = 0.04545 \text{ msec}$



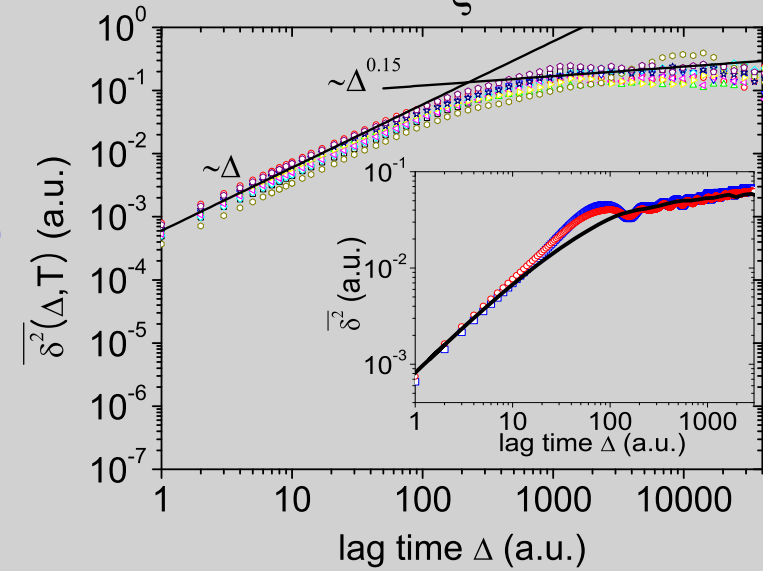
$\Delta = 4.545 \text{ msec}$



Ageing plot:

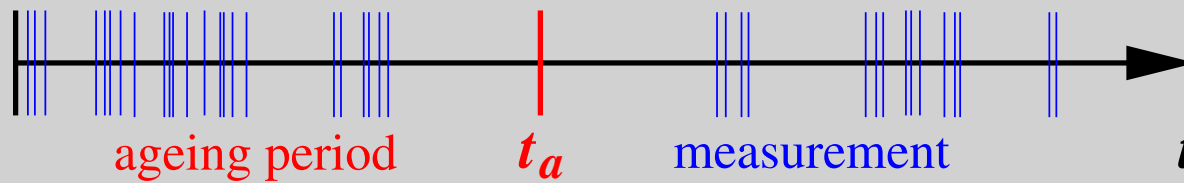


Time averaged MSD:



$$\psi(t) = \frac{d}{dt} \left[1 - \tau^\alpha (\tau + t)^{-\alpha} \exp \left(-\frac{t}{\tau^*} \right) \right] \therefore \tau^* \approx 182 \text{ msec}$$

Ageing effects in single trajectory time averages

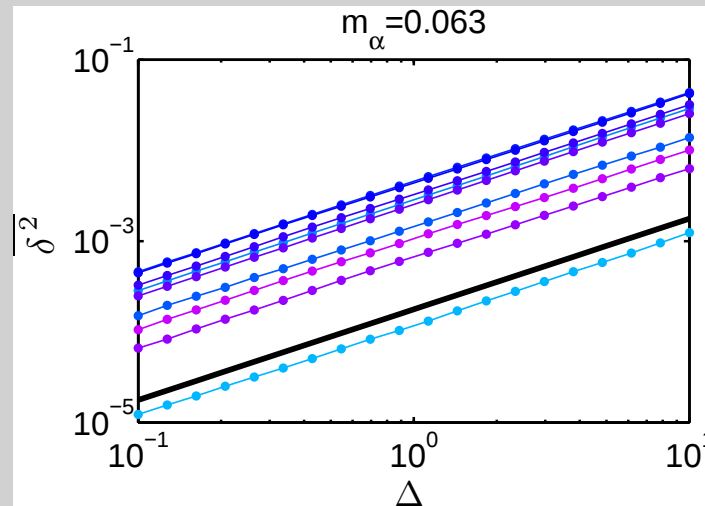
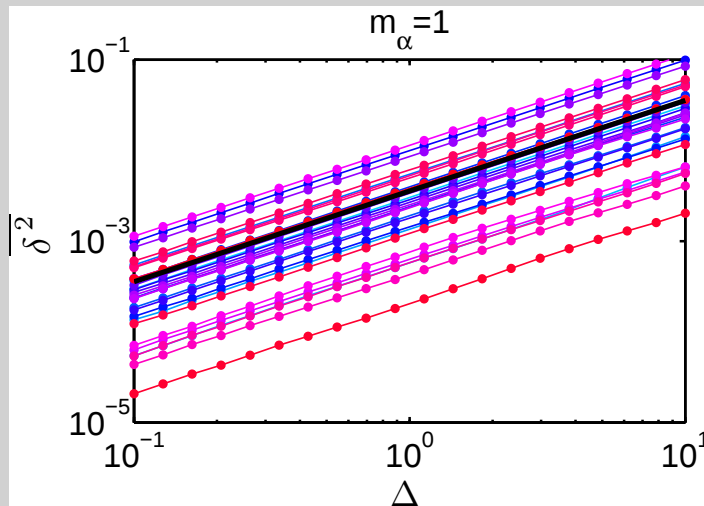


Ageing mean squared displacement ($\Lambda(z) = (1 + z)^\alpha - z^\alpha$)

$$\left\langle \overline{\delta^2(\Delta)} \right\rangle_a = \frac{\Lambda_\alpha(t_a/T) g(\Delta)}{\Gamma(1 + \alpha) T^{1-\alpha}} \Leftrightarrow \langle x^2(t) \rangle_a \simeq \begin{cases} t^\alpha, & t_a \ll t \\ t_a^{\alpha-1} t, & t_a \gg t \end{cases}$$

Probability to make at least one step during $[t_a, t_a + T]$: *population splitting*

$$m_\alpha(T/t_a) \simeq (T/t_a)^{1-\alpha}, \quad T \ll t_a$$



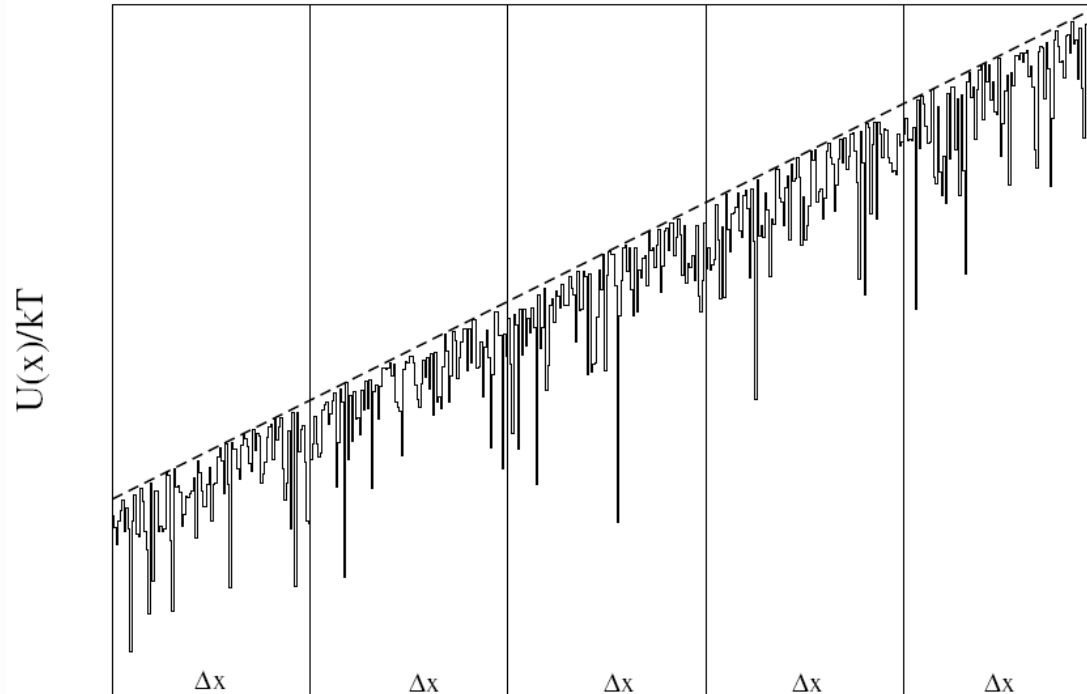
Quenched trap model for scale-free process

At each lattice site E_x is the depth of the trap. $\{E_x\}$ are iid random variables with PDF

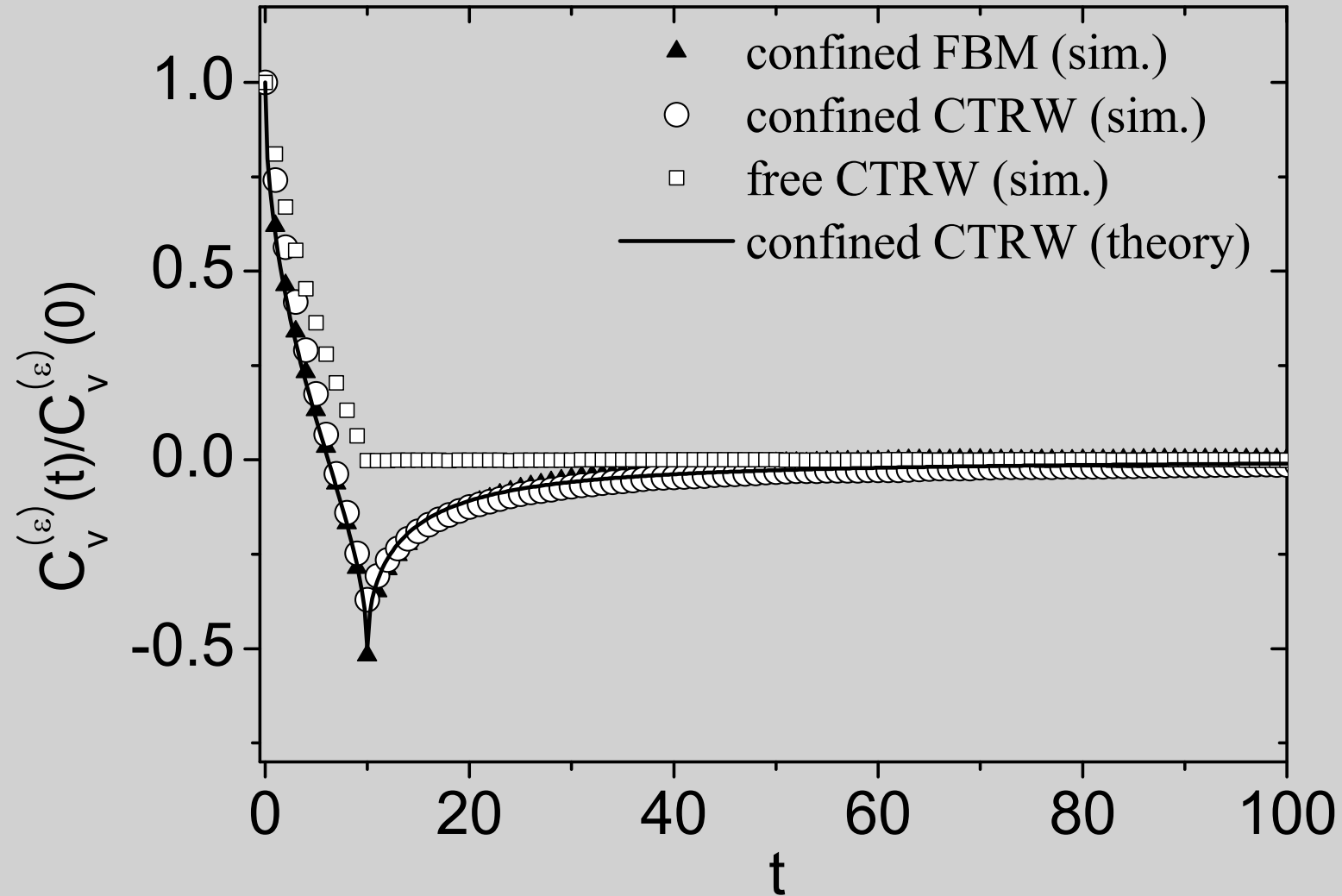
$$\varrho(E) = \frac{1}{k_B T_g} \exp\left(-\frac{E}{k_B T_g}\right)$$

By Arrhenius' law for escape, $\tau_x = \exp(E_x/k_B T)$, we obtain

$$\psi(t) \simeq \frac{T}{T_g} t^{-(1+T/T_g)}$$



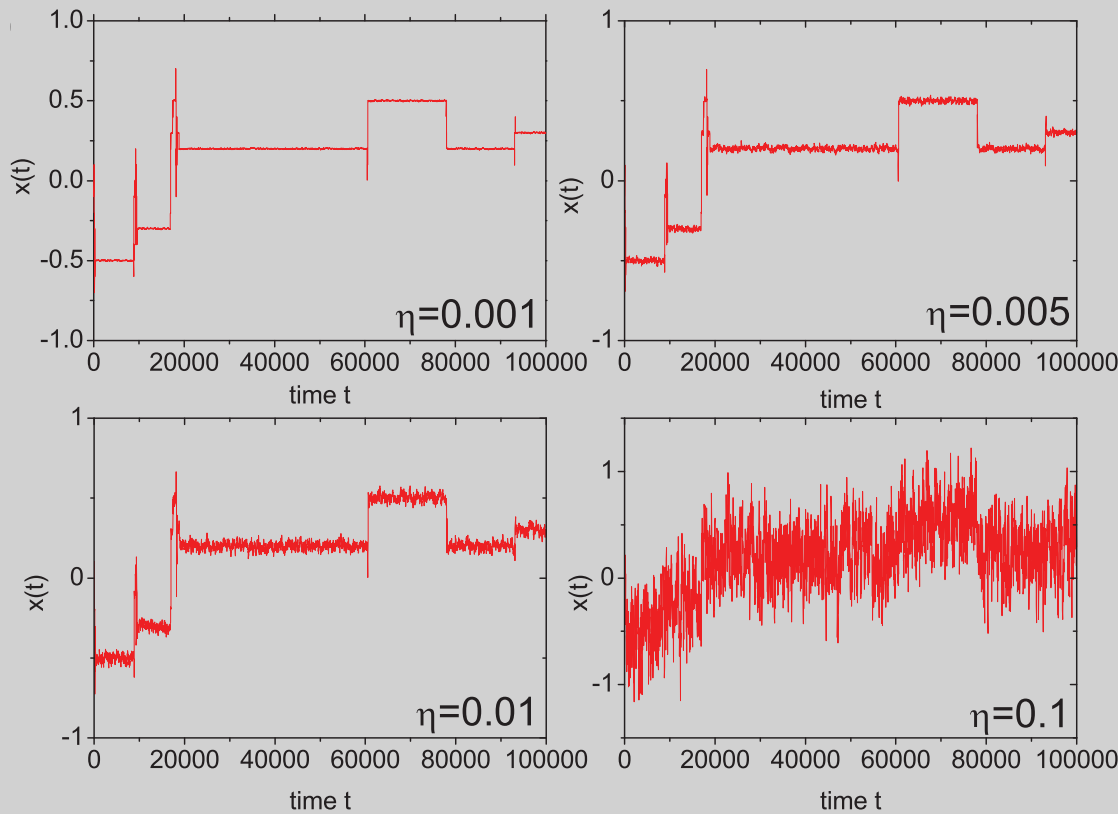
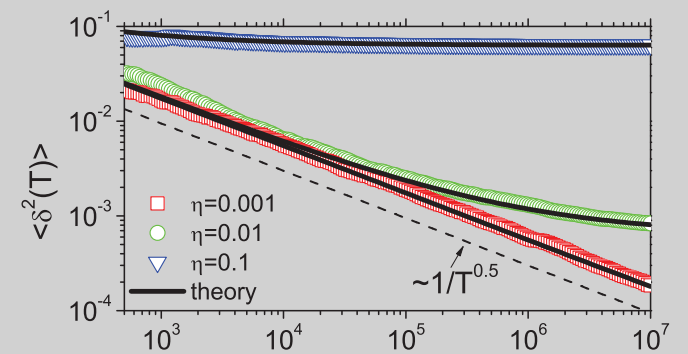
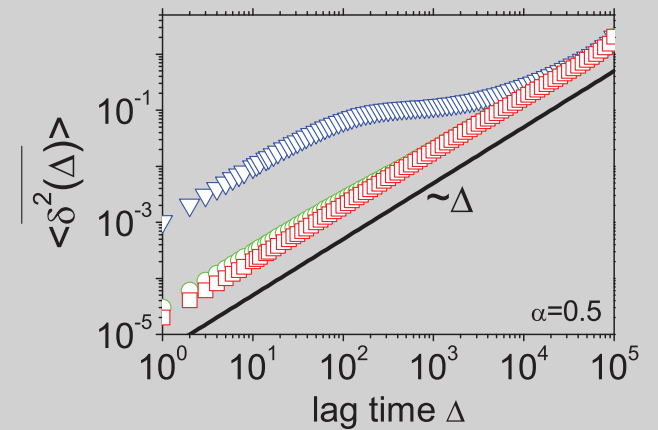
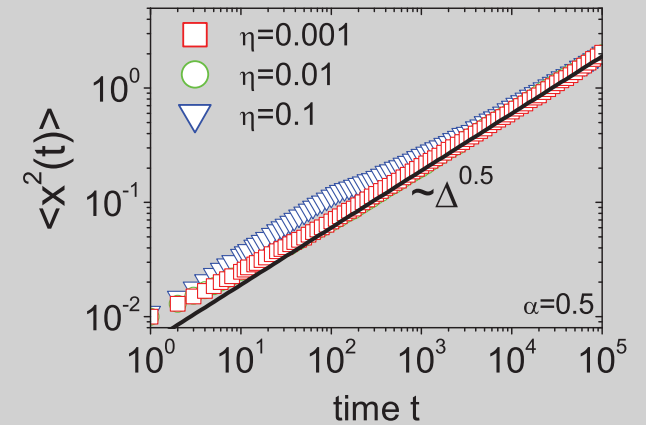
Velocity autocorrelation function: CTRW vs FBM



Noisy CTRW processes /w OU-noise

$$\langle x^2(t) \rangle = \frac{2K_\alpha}{\Gamma(1+\alpha)} t^\alpha + \frac{\eta^2 D}{k} (1 - e^{-2kt})$$

$$\langle \overline{\delta^2(\Delta)} \rangle \sim \frac{2K_\alpha}{\Gamma(1+\alpha)} \frac{\Delta}{T^{1-\alpha}} + \frac{2\eta^2 D}{k} (1 - e^{-k\Delta})$$



Correlated Continuous Time Random Walks

Waiting time ψ_i is conditioned by waiting time ψ_{i-1} :

$$P(\psi_i, \psi_{i-1}) = P(\psi_i | \psi_{i-1}) P(\psi_{i-1})$$

By iteration this means that ψ_i depends on all previous waiting time:

$$P(\psi_i, \psi_{i-1}, \dots, \psi_0) = \prod_{n=1}^i P(\psi_n | \psi_{n-1}) P(\psi_0)$$

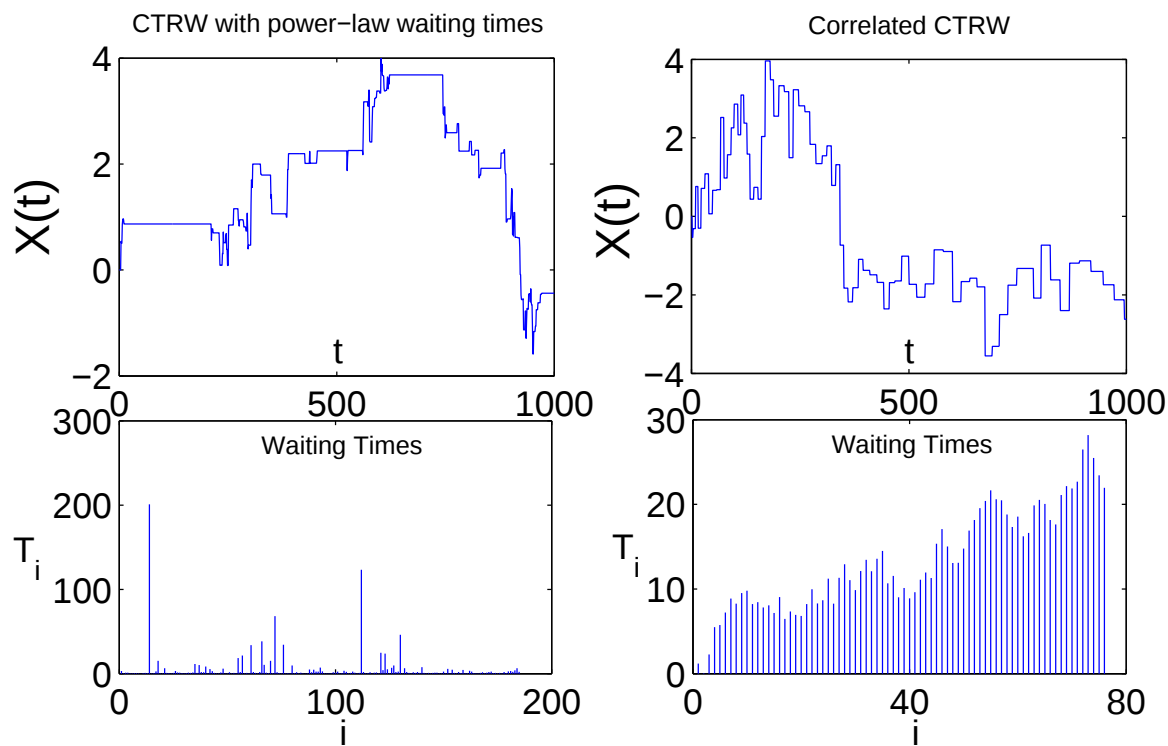
Problem: difficult to evaluate

Question: can we find a more direct way of introducing correlations?

Correlated CTRW: coupling of waiting times

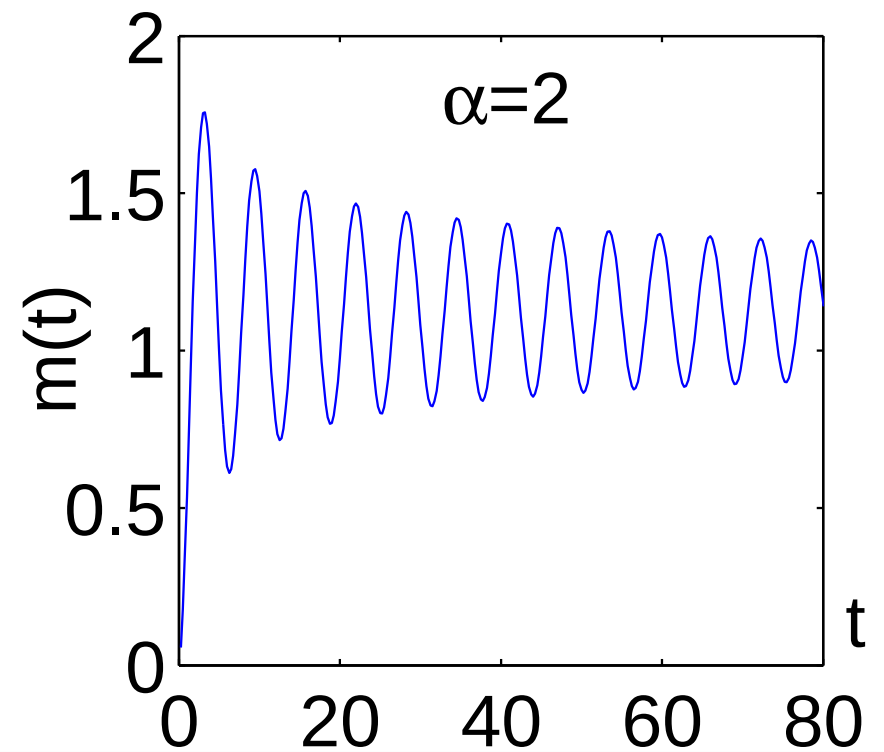
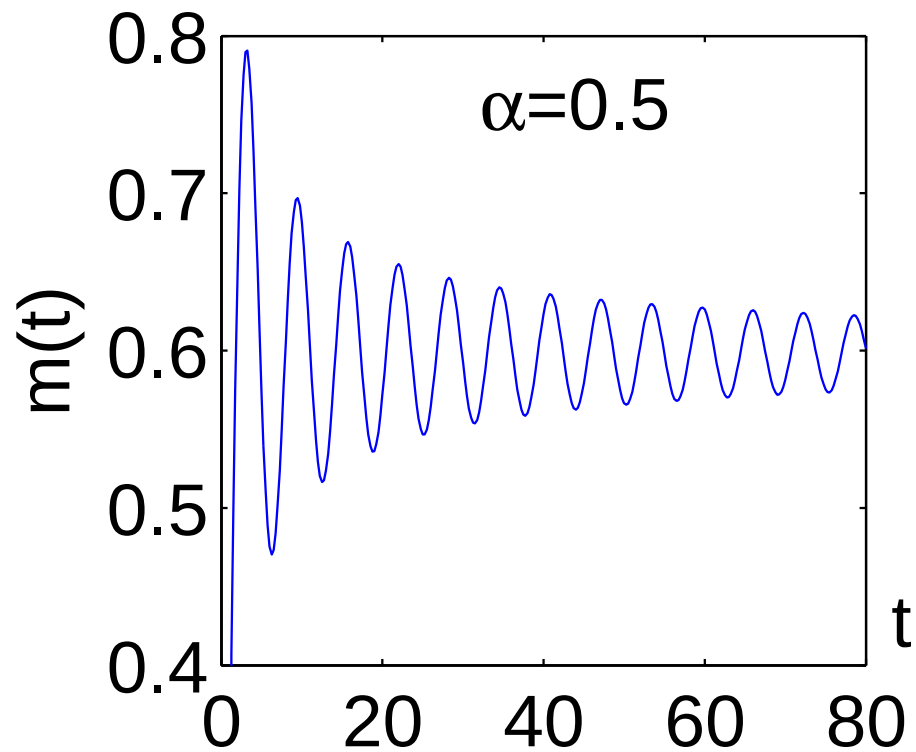
Correlated waiting times: $\tau_i = |\tau_{i-1} + \xi_i|$, such that $\langle \tau \rangle \simeq t^{1/\alpha}$

$$\langle \overline{\delta^2(\Delta, T)} \rangle \simeq \frac{2K_\gamma \Delta}{T^{1-\alpha/(1+\alpha)}}, \quad \langle x^2(t) \rangle = 2K_\gamma t^\gamma, \quad \gamma = \frac{\alpha}{\alpha + 1} \in (0, 2/3]$$



CCTRW: Ageing behaviour

Ageing response to sinusoidal force:

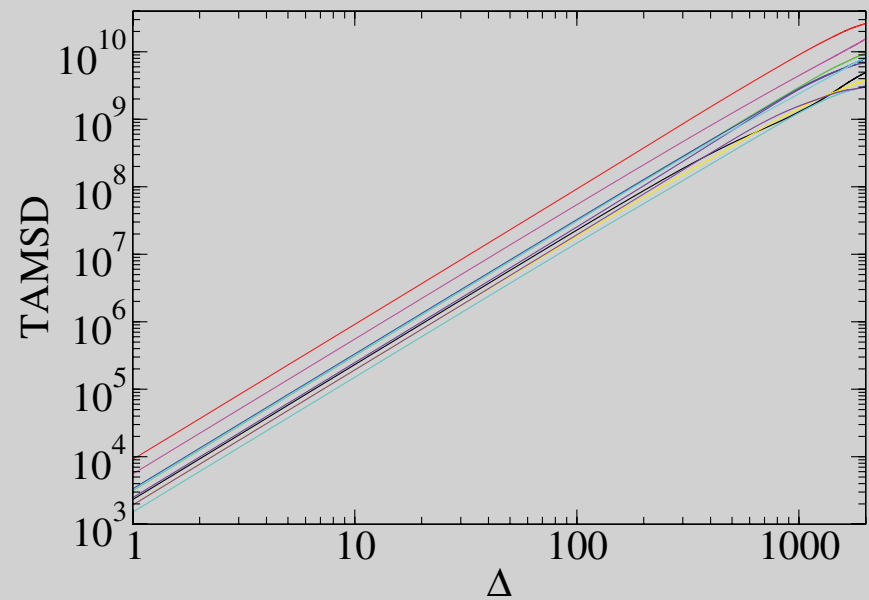
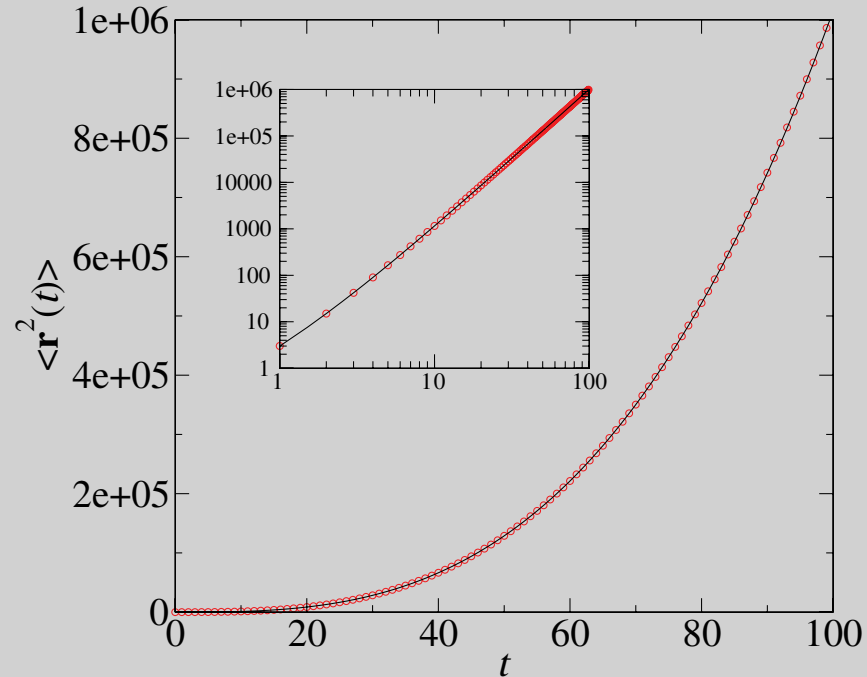


Coupling of jump lengths

Correlated jump lengths: $\lambda_i = \lambda_{i-1} \pm \xi_i$

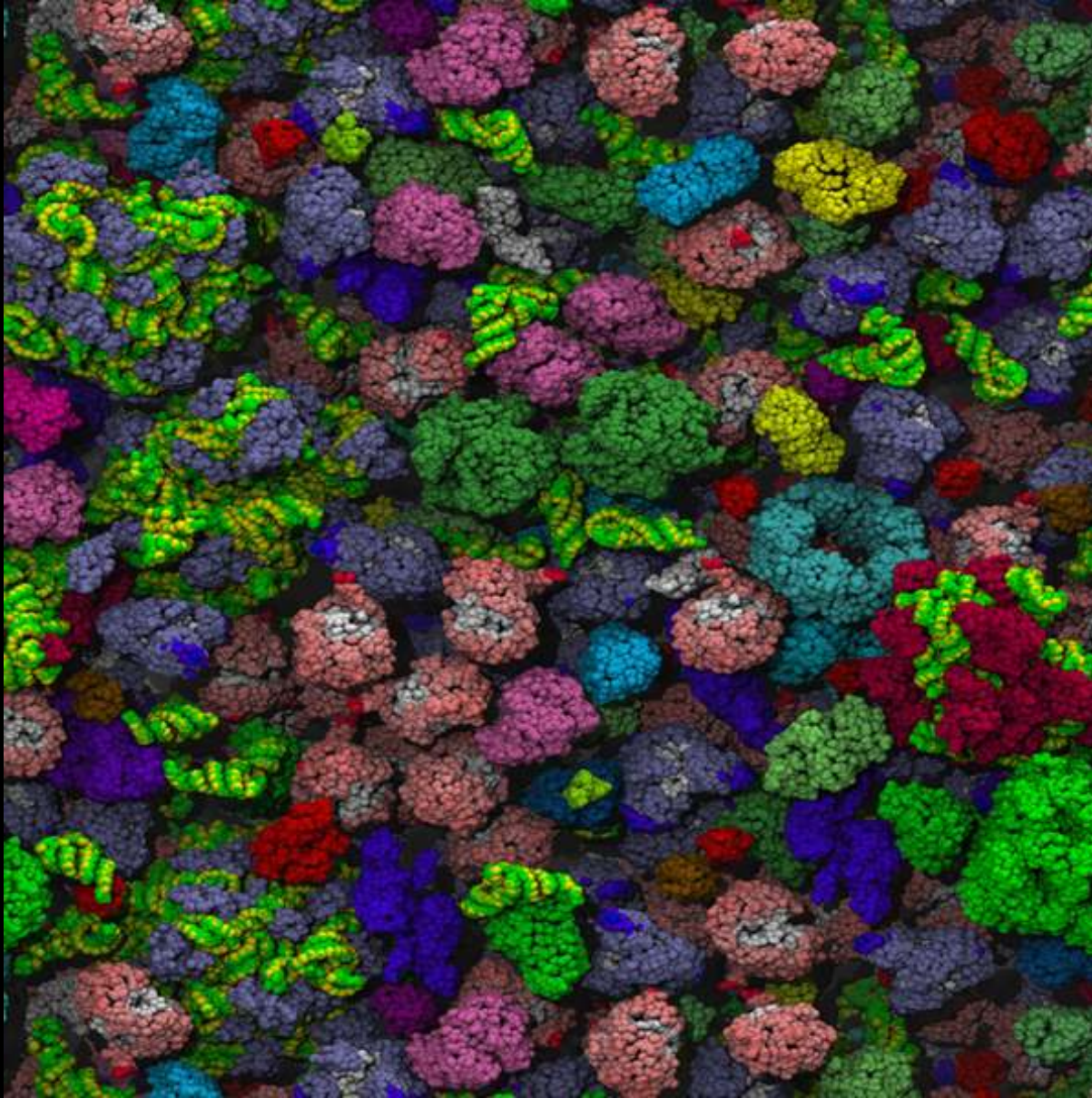
Gaussian distribution of jump lengths:

$$\langle x^2(t) \rangle = \frac{t(t+1)(2t+1)\sigma^2}{4} \Leftrightarrow \langle \overline{\delta^2(\Delta, T)} \rangle \sim \frac{3\sigma^2}{4} \Delta^2 T, \quad \Delta \ll T$$





Macromolecular crowding in living cells



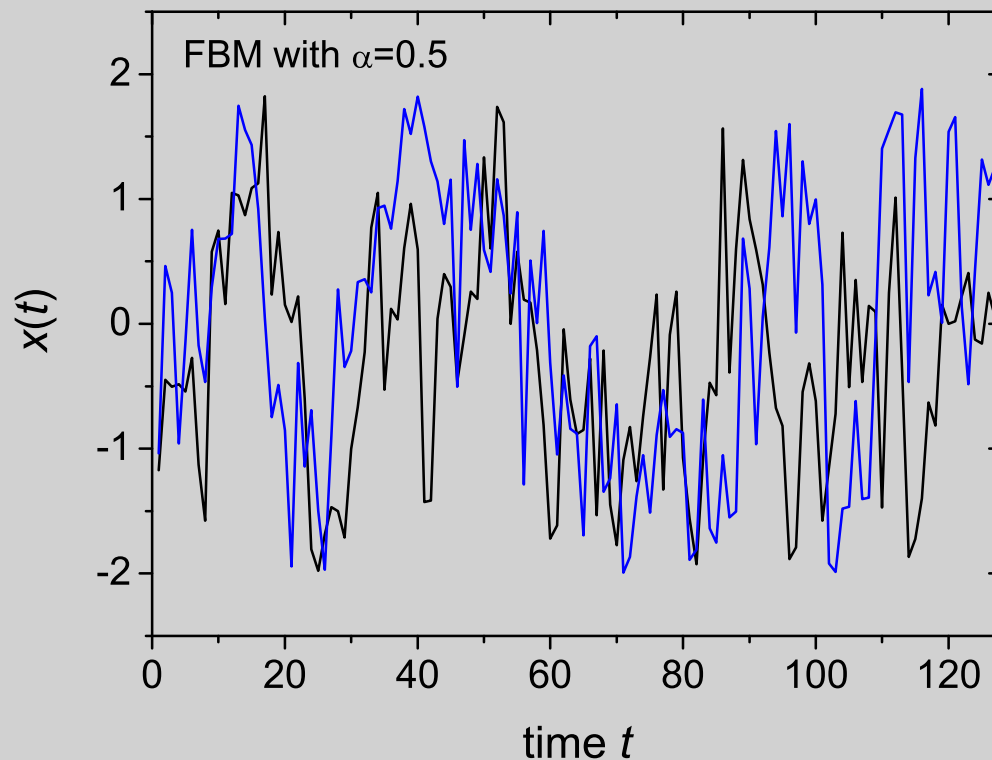
SR McGuffee & AH Elcock, PLoS Comp Biol (2010)

Fractional Brownian motion (viscoelastic environment)

Fractional Gaussian noise: $\langle \xi(t_1)\xi(t_2) \rangle \sim \alpha K_\alpha (\alpha - 1) |t_1 - t_2|^{\alpha-2}$

Position autocorrelation: $\langle x(t_1)x(t_2) \rangle = K_\alpha \left(t_1^\alpha + t_2^\alpha - |t_1 - t_2|^\alpha \right); \quad t_1, t_2 > 0$

Ergodicity is satisfied: $\langle \overline{\delta^2(\Delta, T)} \rangle \sim 2K_\alpha \Delta^\alpha = \langle x^2(\Delta) \rangle$



Viscoelastic diffusion $\langle \mathbf{r}^2(t) \rangle \simeq K_\alpha t^\alpha$ is asymptotically ergodic

Fractional Langevin equation:

$$\ddot{\mathbf{r}}(t) + \int_0^t \eta(t-t') \dot{\mathbf{r}}(t') dt' = \boldsymbol{\zeta}(t)$$

with $\eta(t) = \sum_{i=1}^N a_i(k) e^{-\nu_i t} \rightarrow t^{-\alpha}$ &
 $\langle \zeta_i(t) \zeta_j(t') \rangle = \delta_{ij} k_B T \eta(|t-t'|)$, Gauss

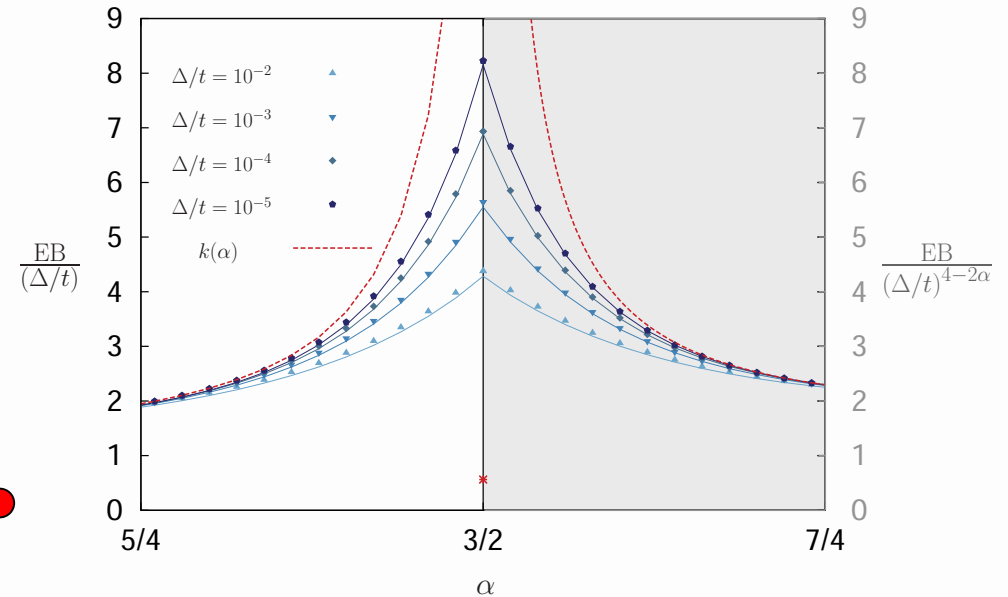
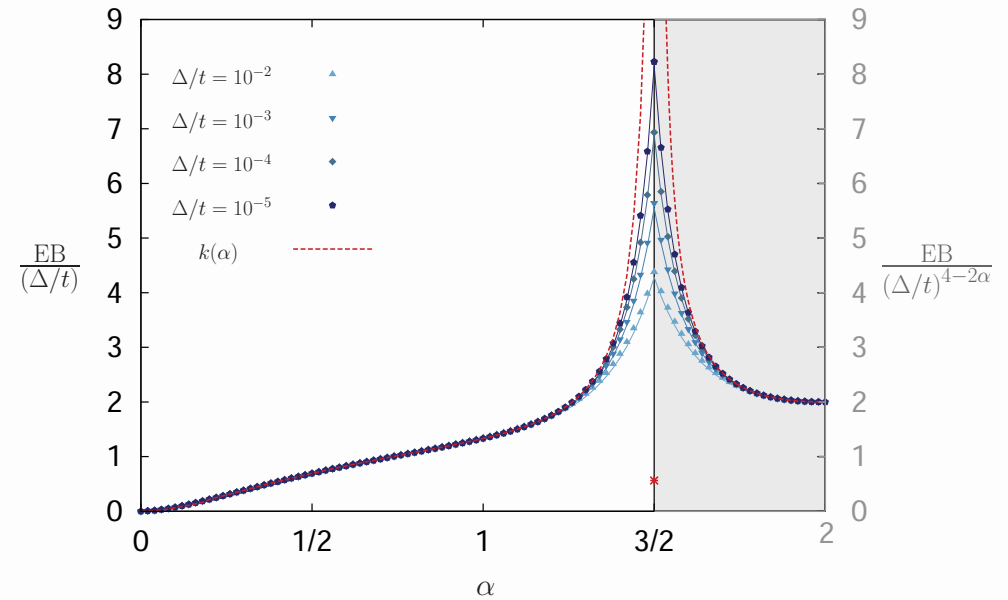
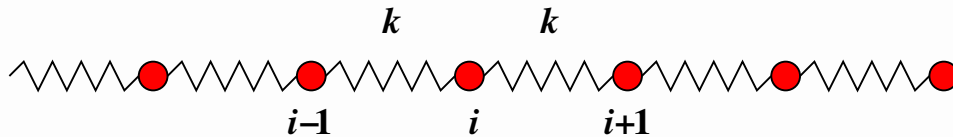
Fractional Brownian motion:

$$\dot{\mathbf{r}}(t) = \boldsymbol{\zeta}(t)$$

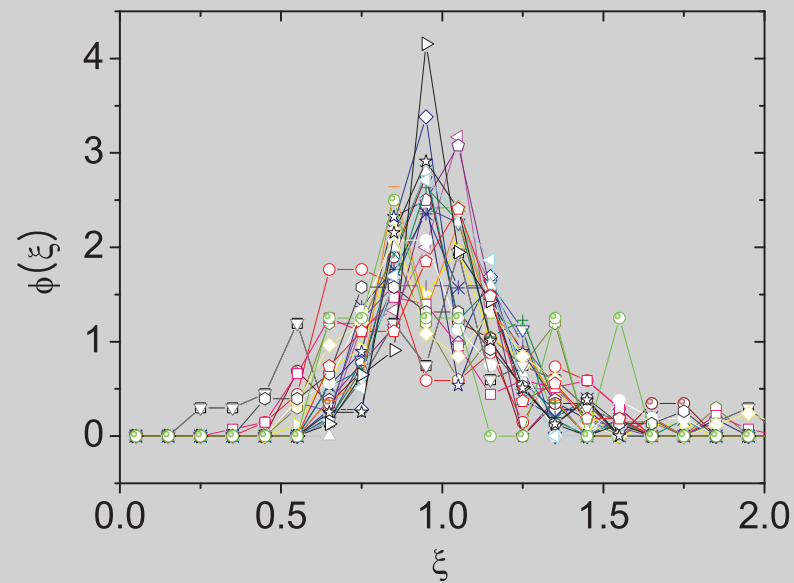
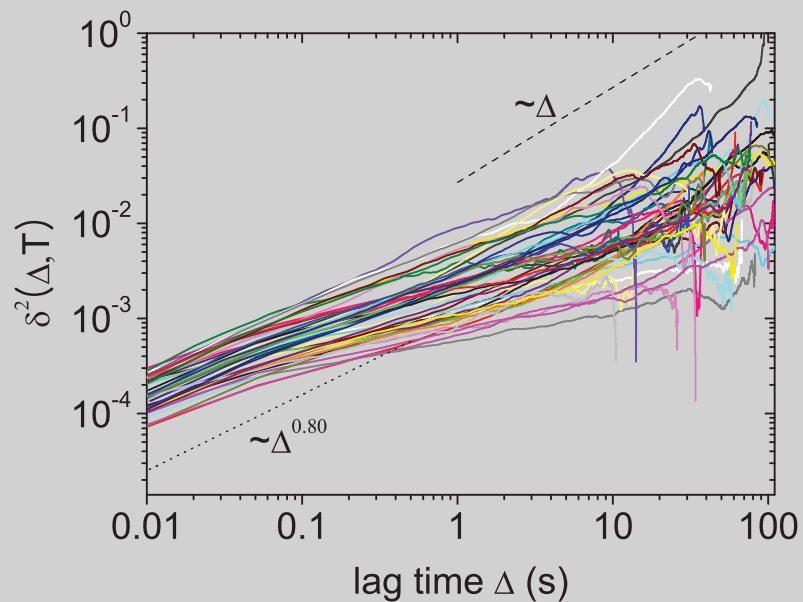
In both cases: $\lim_{T \rightarrow \infty} \overline{\delta^2(\Delta)} = \langle \mathbf{r}^2(\Delta) \rangle$

Ergodicity breaking parameter:

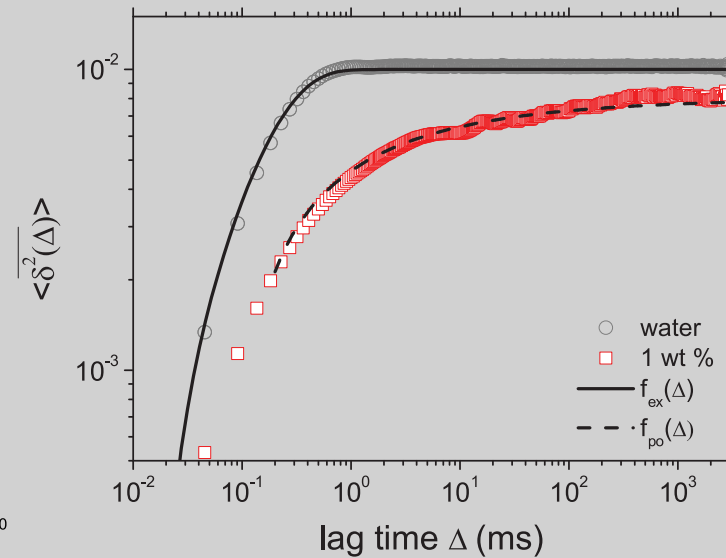
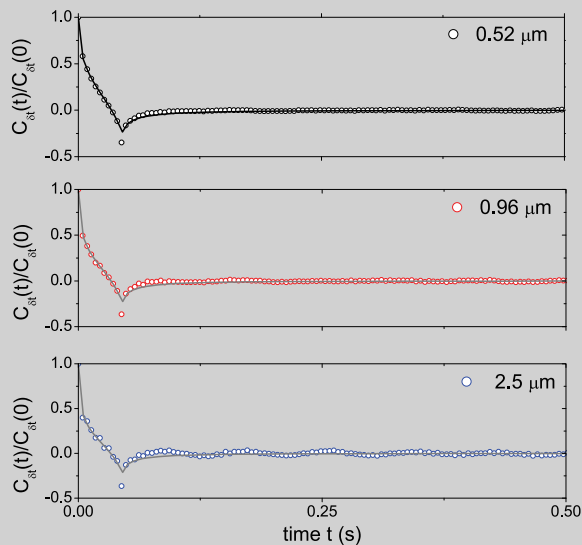
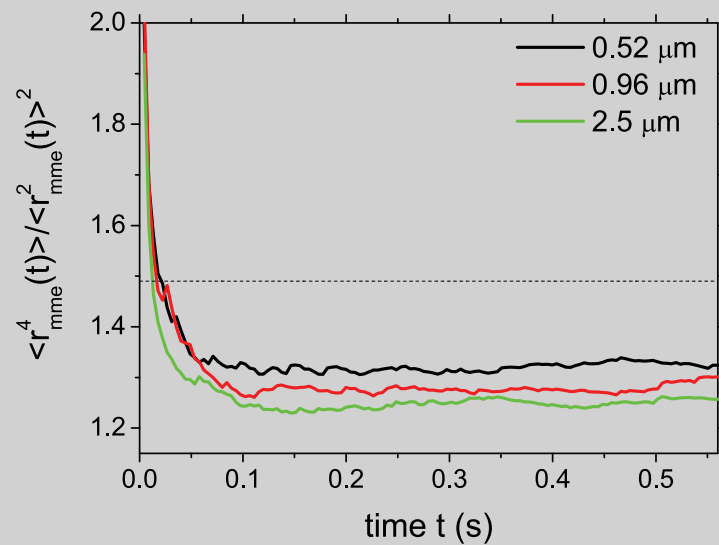
$$\text{EB} = \langle \xi^2 \rangle - 1, \xi = \overline{\delta^2(\Delta)} / \langle \overline{\delta^2(\Delta)} \rangle$$



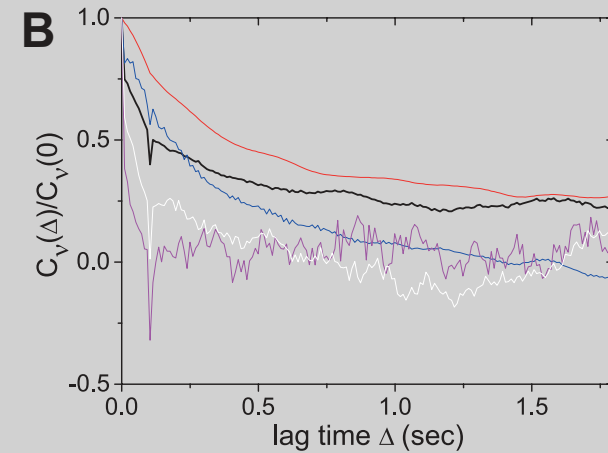
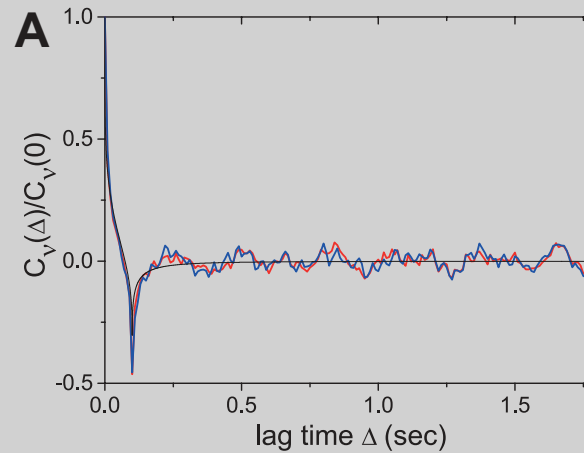
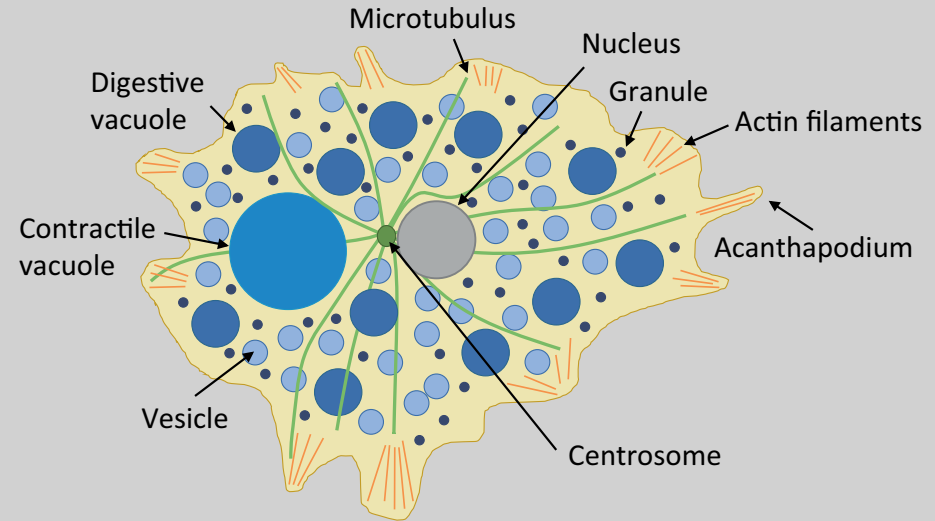
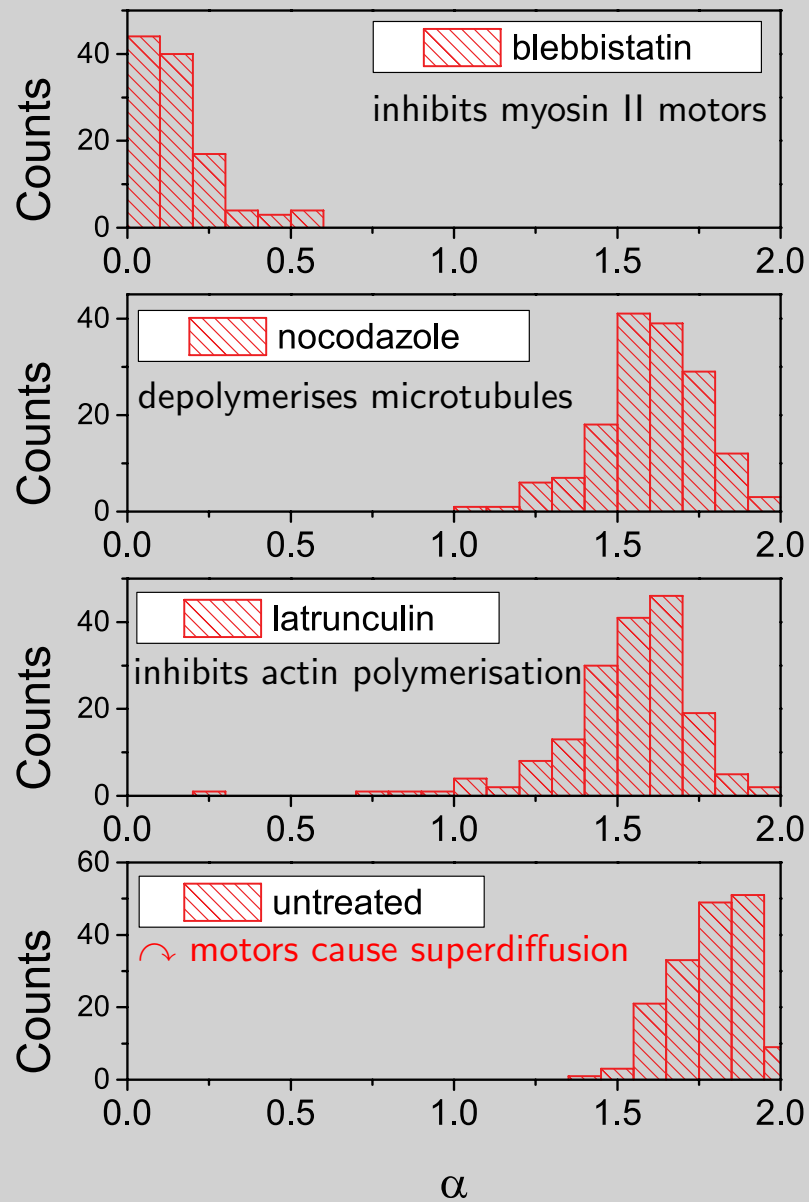
Passive motion of submicron tracers in cells is viscoelastic



Tracer beads in wormlike micellar solution $\downarrow$
Lipid granules in living yeast cells $\downarrow$



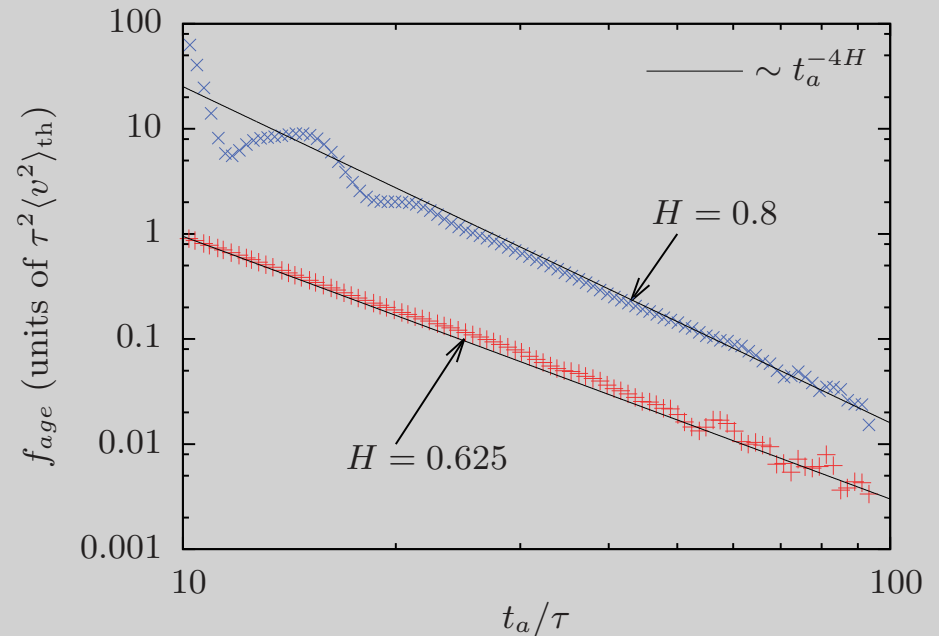
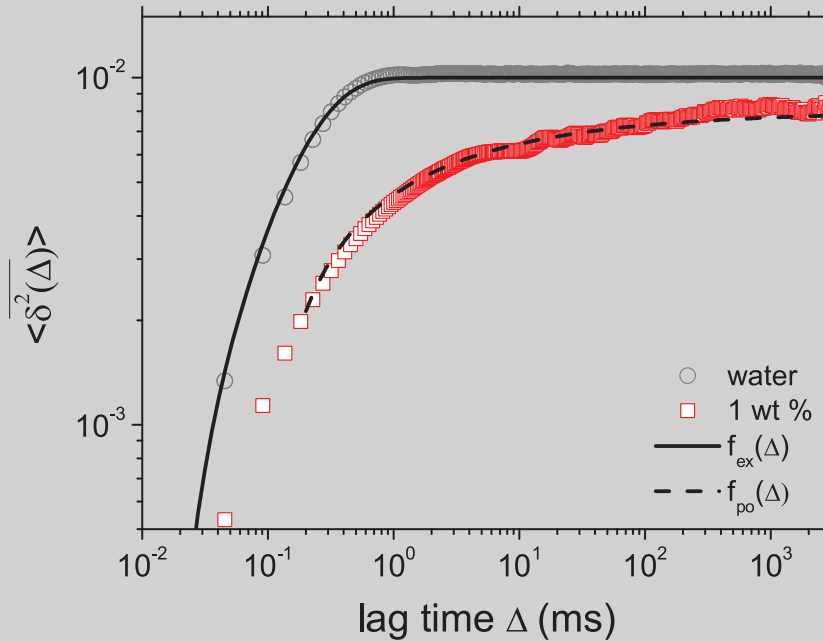
Superdiffusion in supercrowded *Acanthamoeba castellani*



Viscoelastic anomalous diffusion: transient WEB & ageing

Unbounded motion is ergodic: $\overline{\delta^2(\Delta)} = \langle r^2(\Delta) \rangle$

The motion does not age



Transient WEB in confiment:

$$\overline{\delta^2(\Delta)} \sim 2 \langle x^2 \rangle_{th} \left(1 - \frac{\gamma}{k \Delta^{2-\alpha}} \right)$$

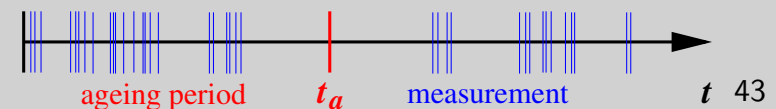
Transient ageing:

$$\overline{\delta^2(\Delta)} = f_{st}(\Delta) + f_{age}(\Delta, t_a, T)$$

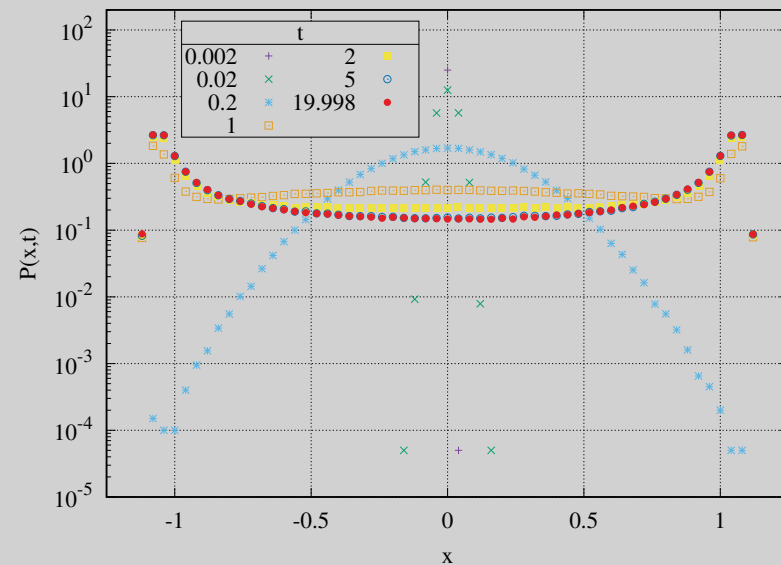
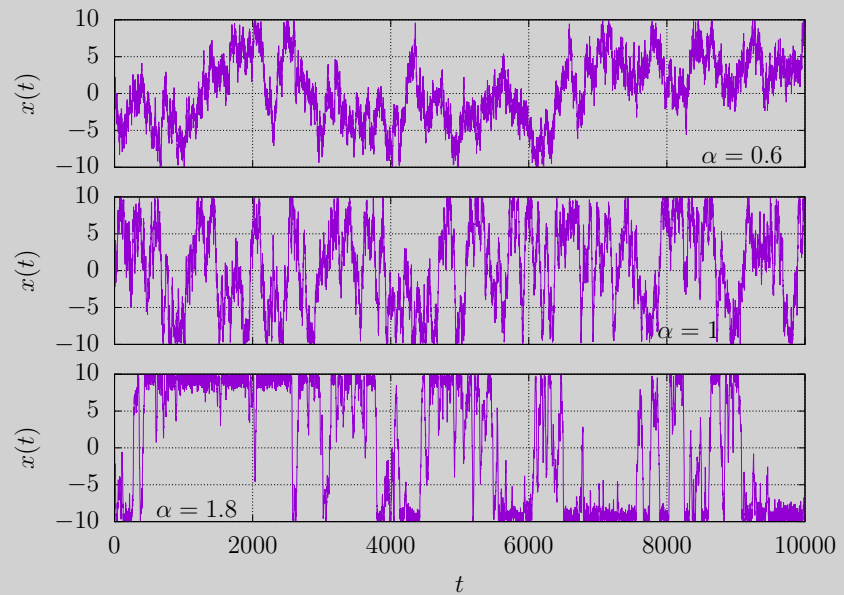
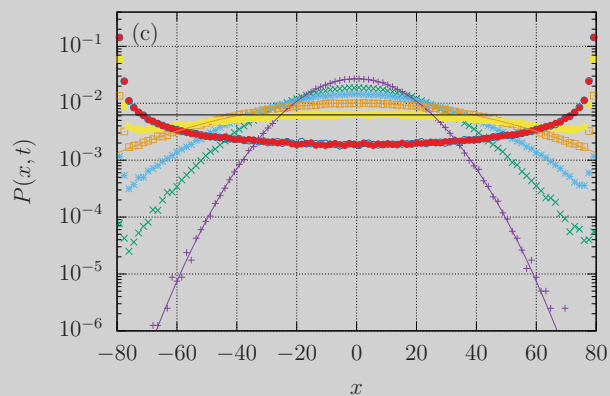
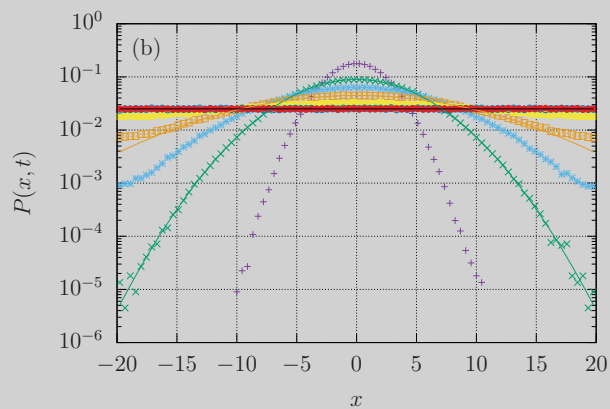
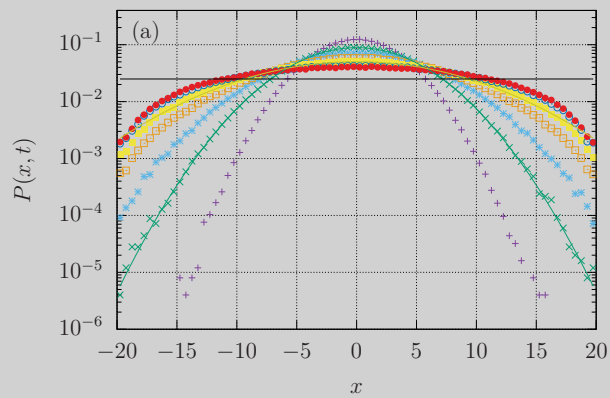
[J-H Jeon, N Leijnse, L Oddershede & RM, New J Phys (2013),
J-H Jeon & RM, PRE (2012)]

[J Kursawe, J Schulz & RM, PRE (2013)]

W Deng & E Barkai, PRE (2009)

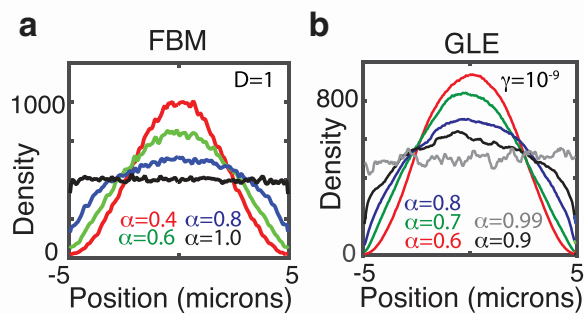
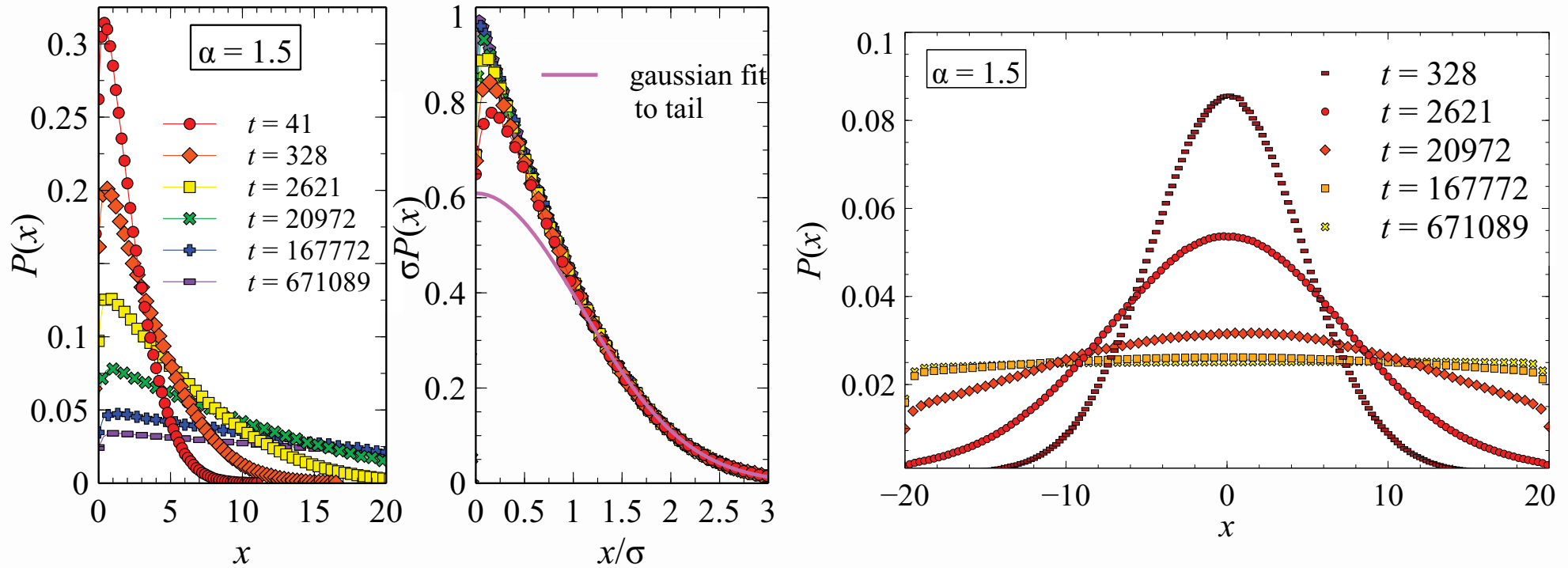


FBM: accretion & depletion effects near boundaries



T Guggenberger, G Pagnini, T Vojta & RM, NJP (2019); see also AHO Wada & T Vojta, PRE (2018)

FLE shows equilibrated behaviour

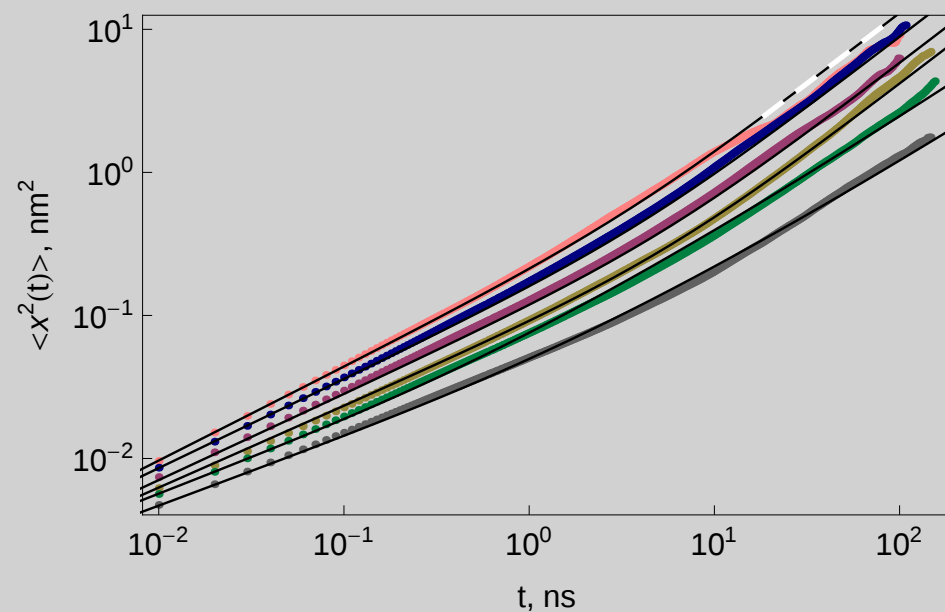
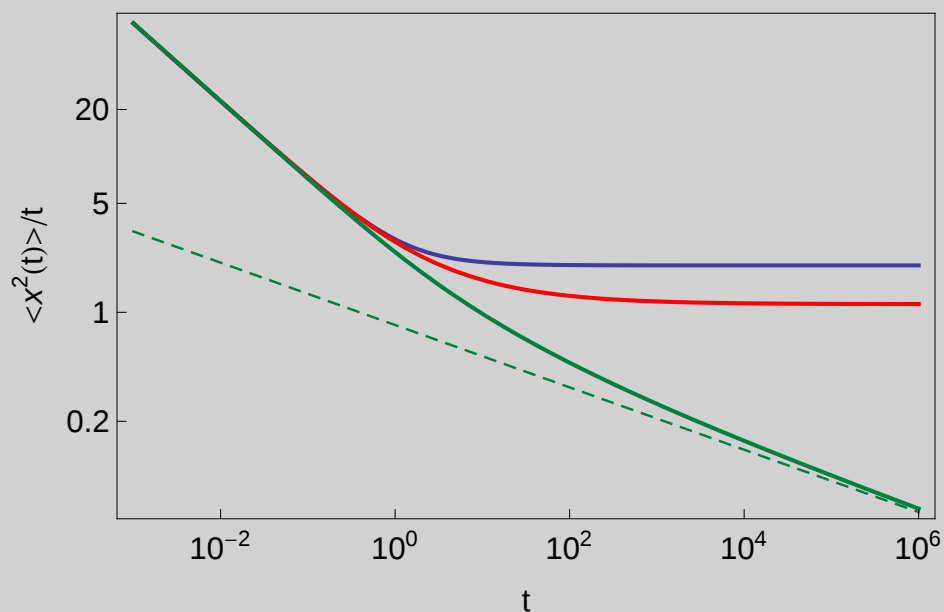
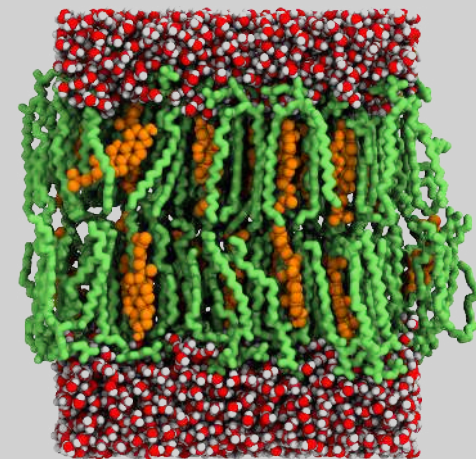


Holmes, BPJ (2019) using hard BC

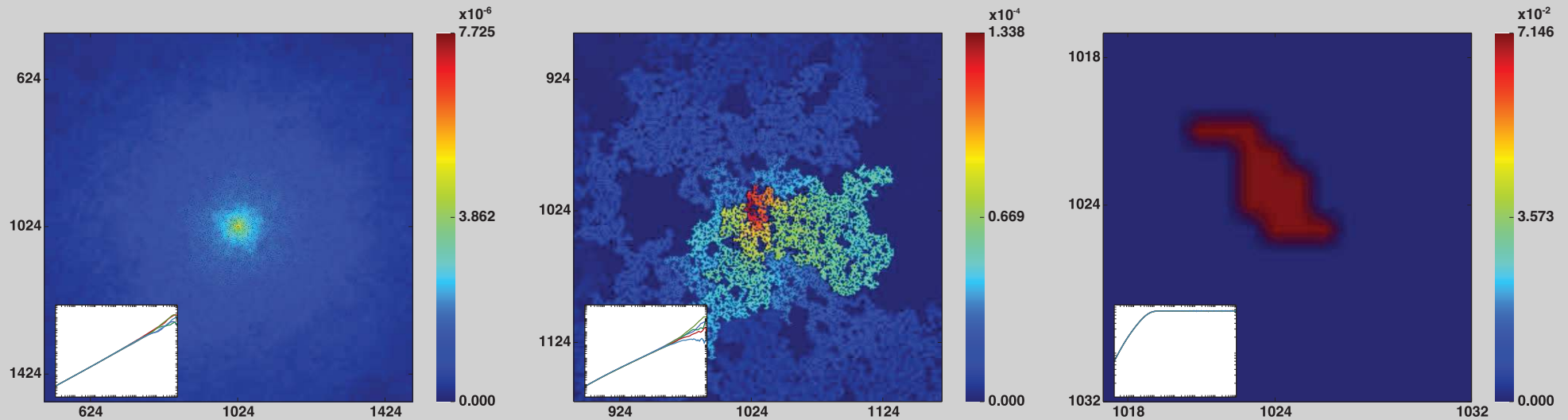
Tempered FLE motion: crossover to faster diffusion

Tempered fractional Gaussian noise:

$$\langle \xi(t) \xi(t + \tau) \rangle = \begin{cases} \frac{C}{\Gamma(2H - 1)} \tau^{2H-2} e^{-\tau/\tau_\star} \\ \frac{C}{\Gamma(2H - 1)} \tau^{2H-2} \left(1 + \frac{\tau}{\tau_\star} \right)^{-\mu} \end{cases}$$

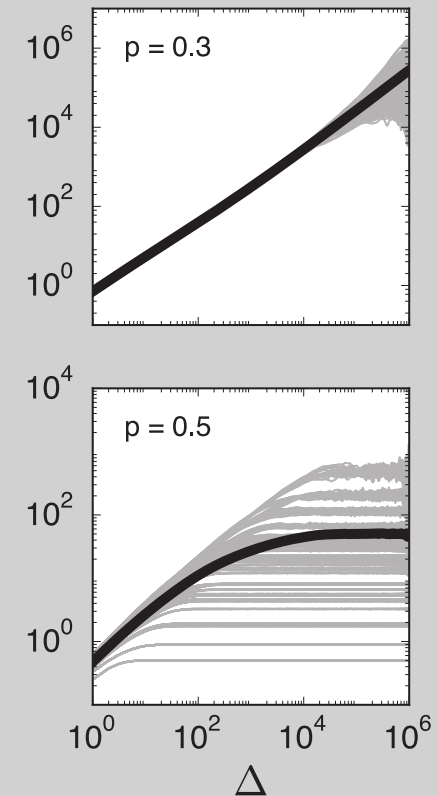
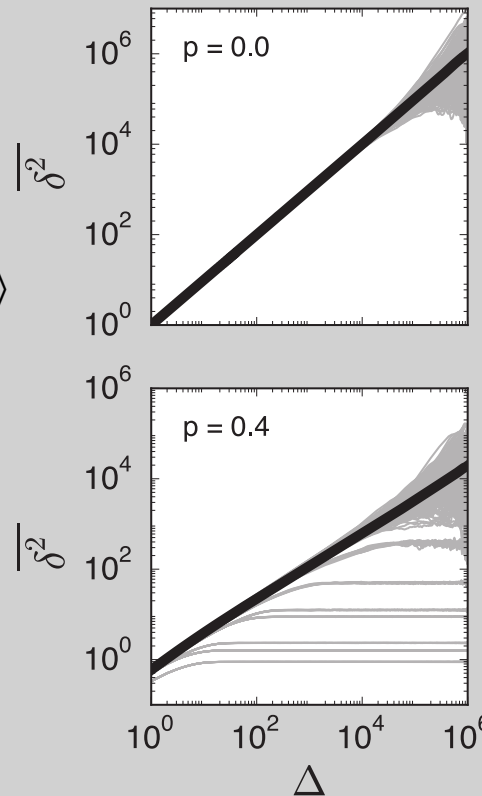


Geometry induced strong non-ergodicity in random fractals



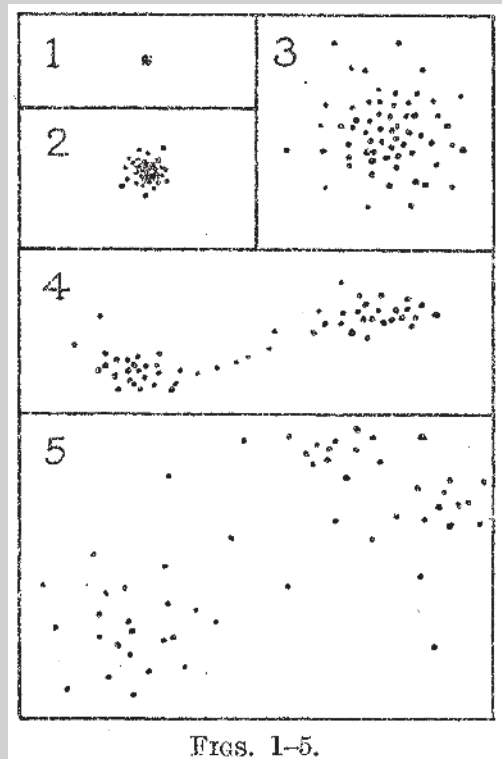
$$\langle \widetilde{\overline{\delta^2(\Delta)}} \rangle = \sum_{i,\zeta} \frac{\overline{\delta_{i,\zeta}^2(\Delta)}}{N \times N_\zeta} = \langle \widetilde{\mathbf{r}^2(t)} \rangle$$

Non-ergodic: $\lim_{T \rightarrow \infty} \overline{\delta^2(\Delta)}$



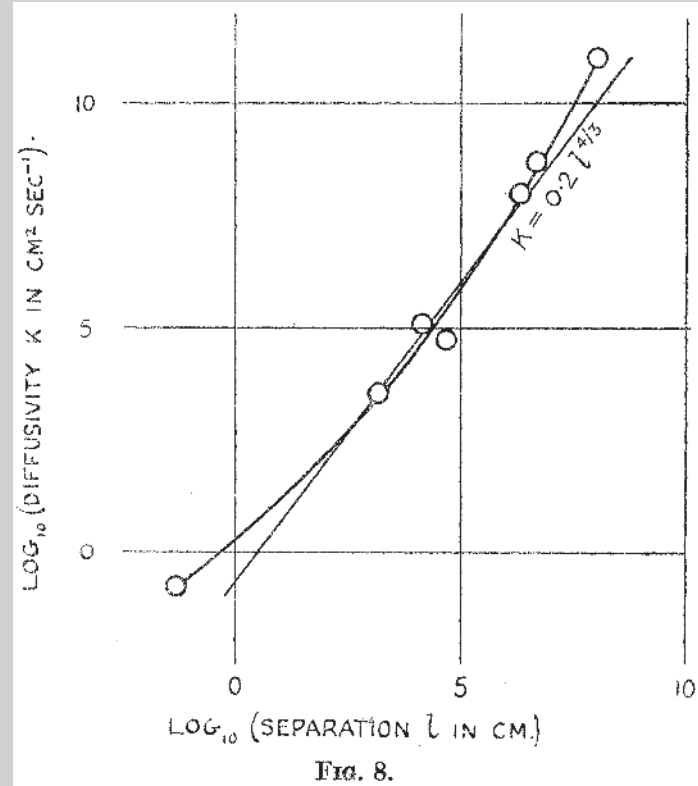
Y Mardoukhi, J-H Jeon & RM, PCCP (2015)

Richardson (relative) diffusion in turbulence: $\langle l^2(t) \rangle \simeq t^3$



Brown

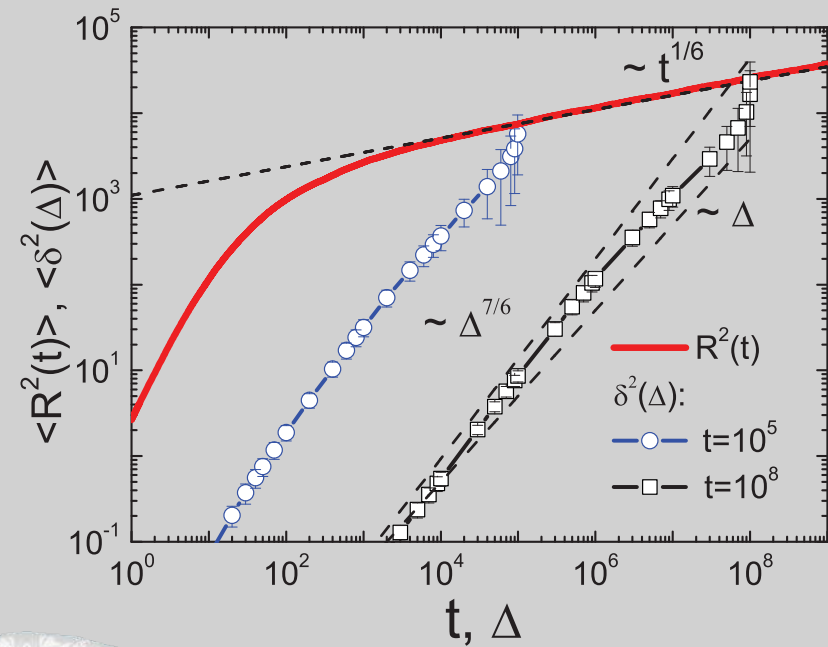
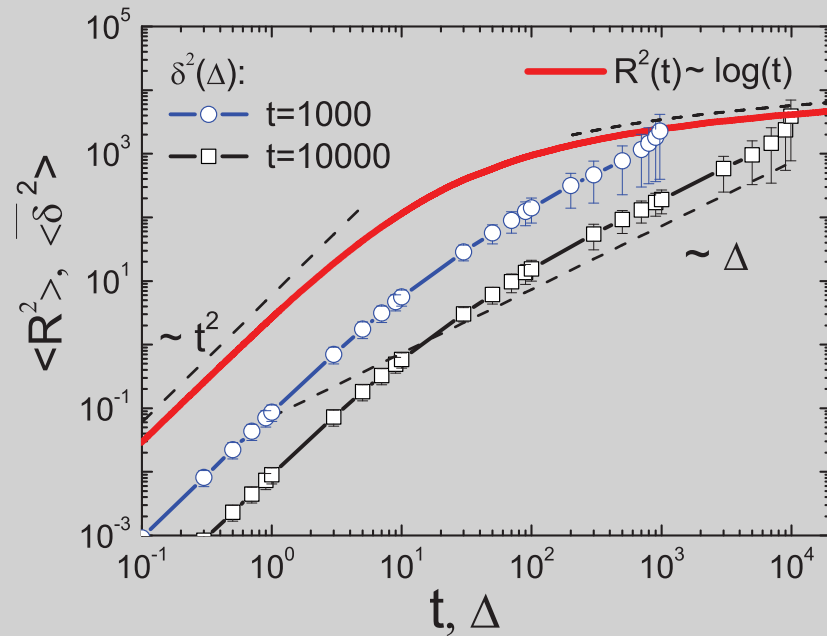
Turbulent spreading



$$\frac{\partial q}{\partial t} = \varepsilon \frac{\partial}{\partial l} \left(l^{4/3} \frac{\partial q}{\partial l} \right), \quad \varepsilon \sim 0.4 \text{ cm}^{2/3} \text{ sec}^{-1} \quad \text{Richardson}$$

$$\frac{\partial q}{\partial t} = \bar{\varepsilon} t^2 \frac{\partial^2}{\partial l^2} q(l, t) \quad \text{Batchelor}$$

WEB in granular gas & SBM as mean field theory



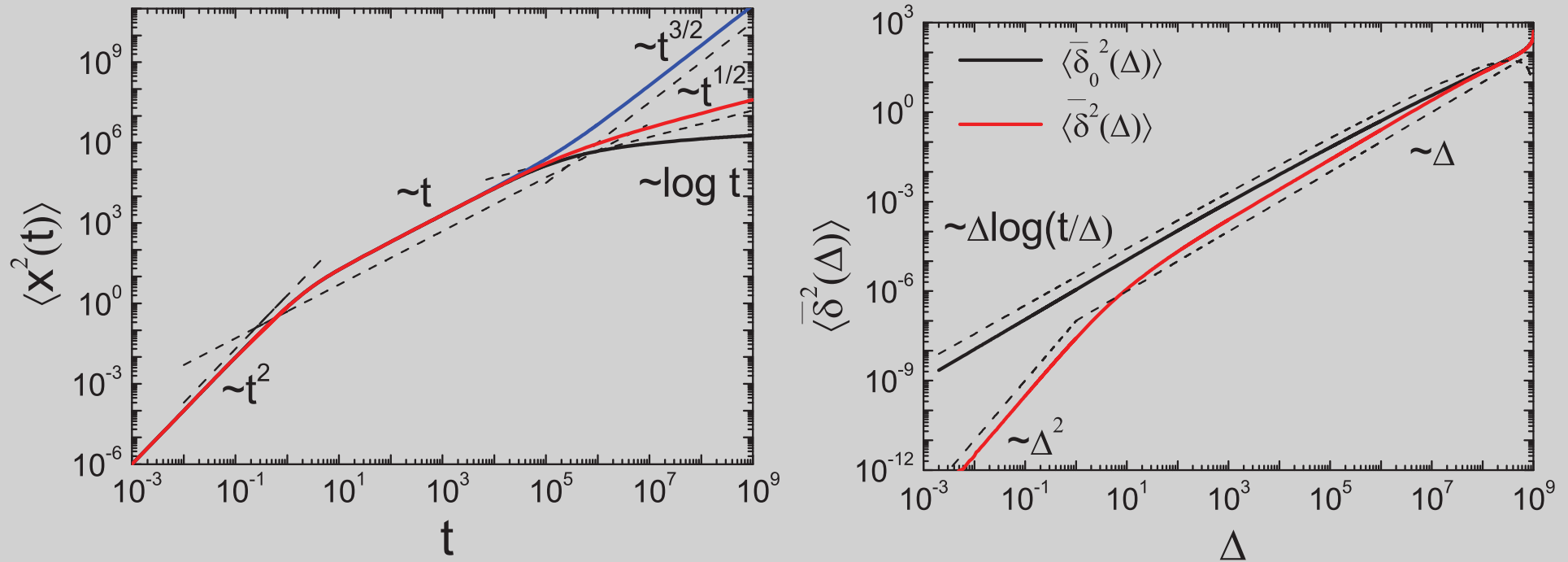
Haff's law: $\mathcal{I}(t) = \mathcal{I}_0 / (1 + t/\tau_0)^2$

$$\langle \mathbf{r}^2(t) \rangle \sim 6D_0\tau_0 \log(1 + t/\tau_0)$$

$$\langle \overline{\delta^2(\Delta)} \rangle \sim 6D_0\tau_0\Delta/T$$



Non-existence of the overdamped limit in slow SBM



Crossover from ballistic to overdamped motion no longer defined by time scale of inverse friction. For small α & ultraslow the SBM overdamped limit is never fulfilled

Ageing case: [H Safdari, A Bodrova, AV Chechkin, AG Cherstvy & RM, PRE (2017)]

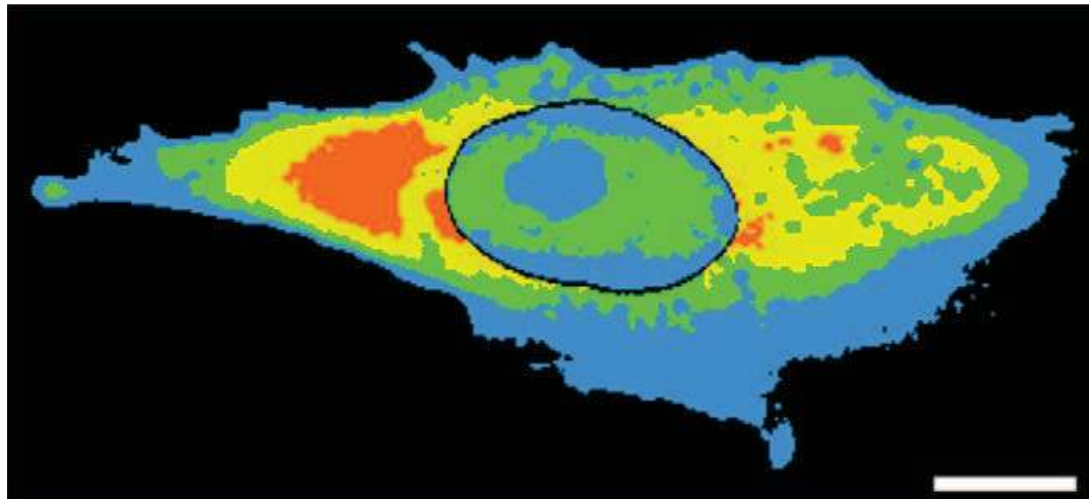
Position dependent diffusivity: heterogeneous diffusion

Markovian diffusion process with $K = K(x)$ à la Richardson:

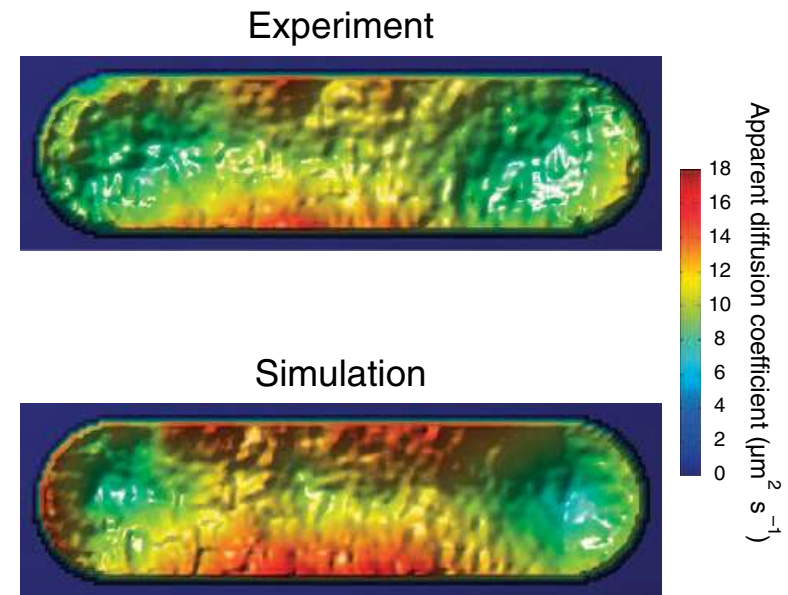
$$K(x) = K_0|x|^\alpha \implies \langle x^2(t) \rangle \simeq t^p$$

Multiplicative Langevin equation:

$$\frac{dx(t)}{dt} = \sqrt{D(x)}\xi(t)$$



Kuehn et al, PLoS ONE (2011)

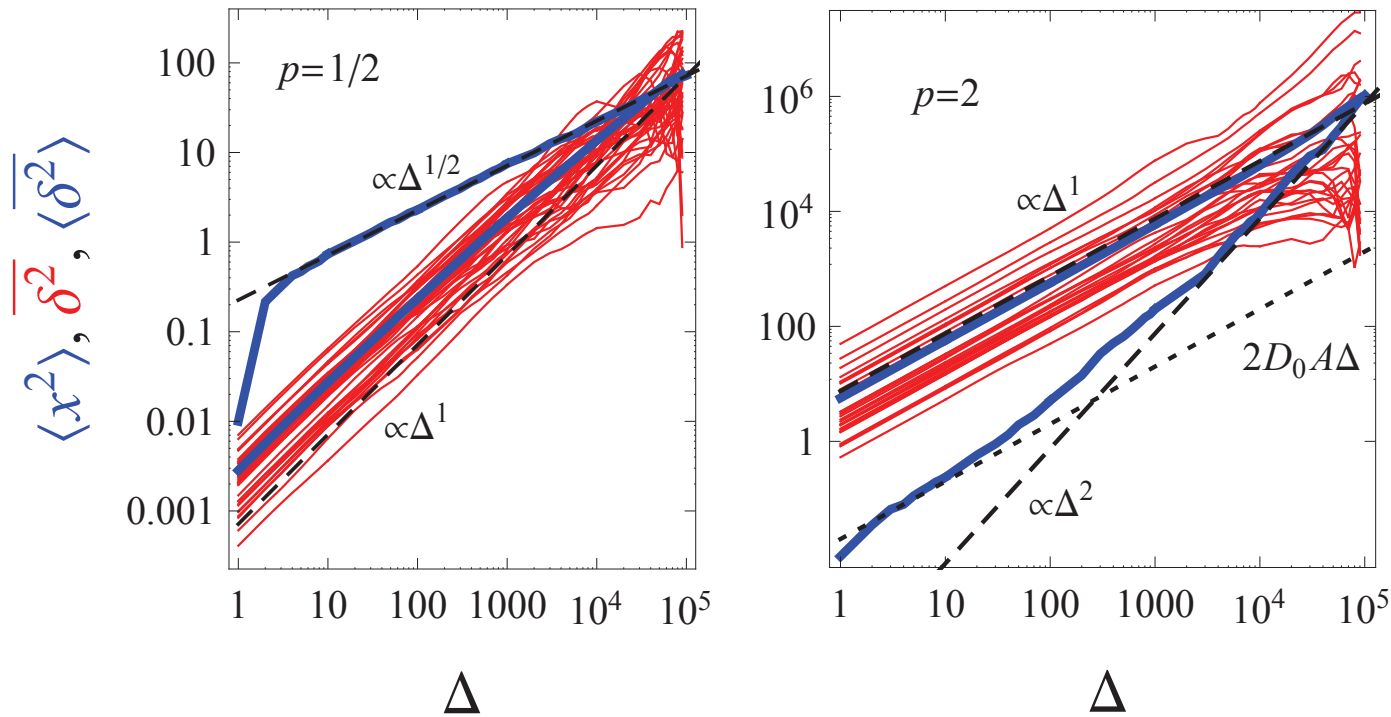


BP English et al, PNAS (2011)

Heterogeneous diffusion processes with $K(x)$

Markovian diffusion process with $K = K(x)$ à la Richardson:

$$K(x) = K_0|x|^\beta \implies \langle x^2(t) \rangle \simeq t^p \quad \langle \overline{\delta^2(\Delta)} \rangle \simeq \frac{\Delta}{T^{1-p}}, \quad p = \frac{2}{2-\beta}$$



Many common features with CTRW: PRE (2014), JPA (2014)

Ageing behaviour in anomalous diffusion processes

Continuous time random walk with $\psi(t) \simeq \tau^\alpha / t^{1+\alpha}$:

$$\left\langle \overline{\delta^2(\Delta)} \right\rangle_a = \Lambda_\alpha(t_a/T) \left\langle \overline{\delta^2(\Delta)} \right\rangle \quad \Leftrightarrow \quad \langle x^2(t) \rangle \simeq \begin{cases} t^\alpha, & t_a \ll t \\ t_a^{\alpha-1} t, & t_a \gg t \end{cases}$$

$$\therefore \Lambda_\alpha(z) = (1+z)^\alpha - z^\alpha \text{ [JHP Schulz, E Barkai \& RM, PRL (2013); PRX (2014)]}$$

Same functional forms for Scaled Brownian Motion with $K(t) = K_\alpha t^{\alpha-1}$

[H Safdari, AV Chechkin, GR Jafari & RM, PRE (2015)]

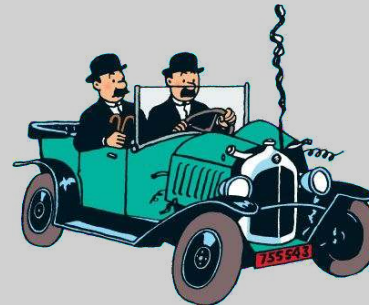
and for heterogeneous diffusion processes with $K(x) = K_0 |x|^\alpha$

[AG Cherstvy, AV Chechkin & RM, JPA (2014)]

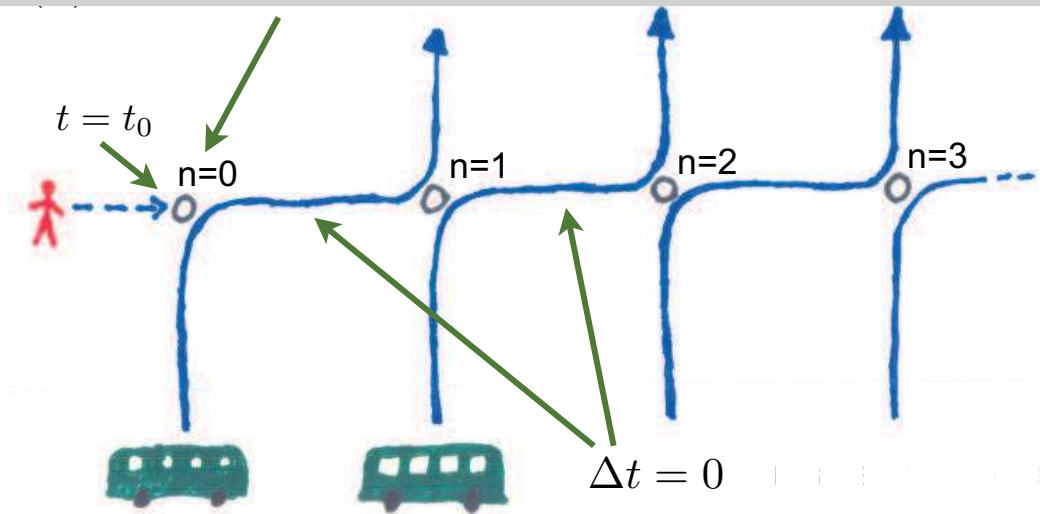
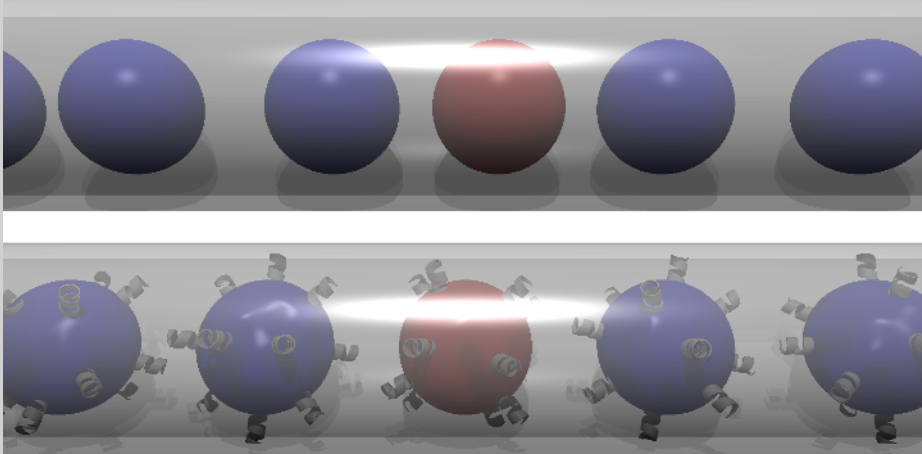
Fractional Langevin equation motion:

$$\left\langle \overline{\delta^2(\Delta)} \right\rangle = f_{\text{st}}(\Delta) + f_{\text{age}}(\Delta; t_a, T) \quad \therefore \quad f_{\text{age}} \simeq \left[v_0^2 - \langle v^2 \rangle_{\text{th}} \right] \Delta^2 t_a^{-2\alpha}, t_a \gg T$$

[J Kursawe, JHP Schulz & RM, PRE (2013)]



Ultralow dynamics in ageing many-particle systems

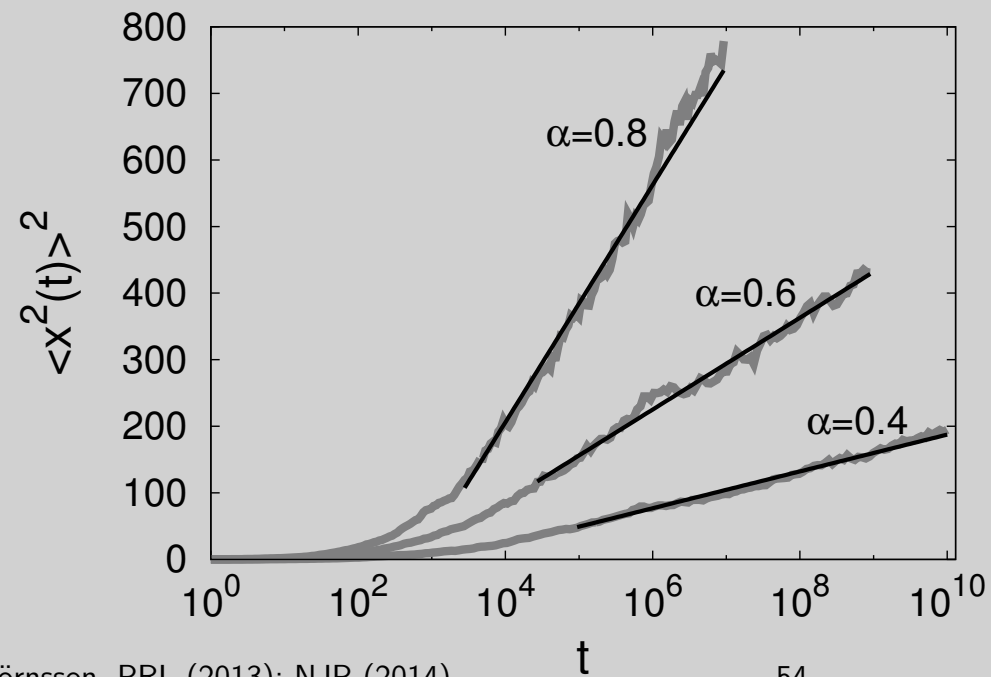


Brownian particles: Harris' law $\langle x^2(t) \rangle \simeq t^{1/2}$

Functionalised particles moving /w $\psi(t) \simeq t^{-1-\alpha}$
 $\hookrightarrow$ scaling argument: $n \rightarrow \log t$:

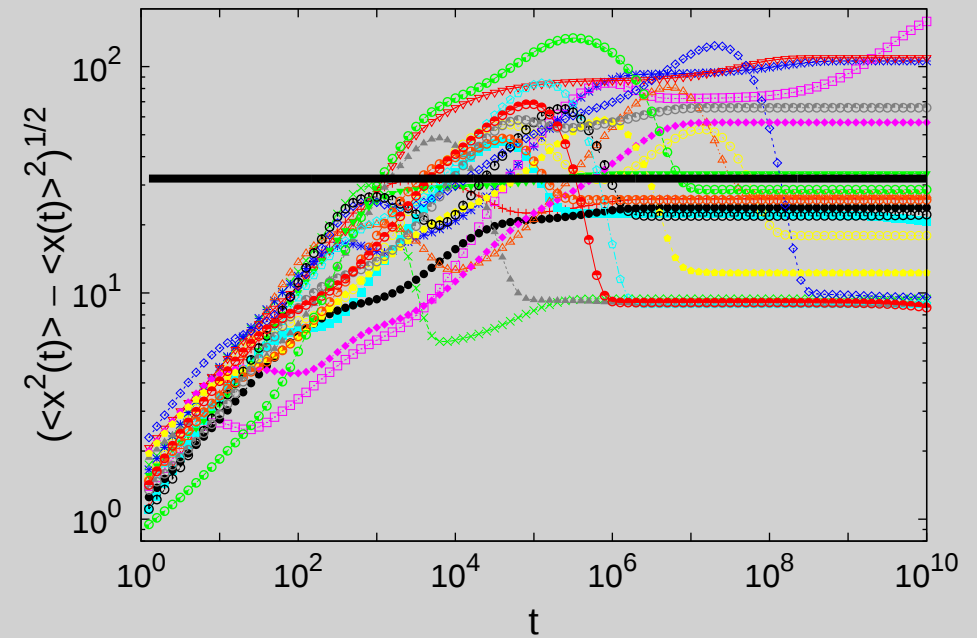
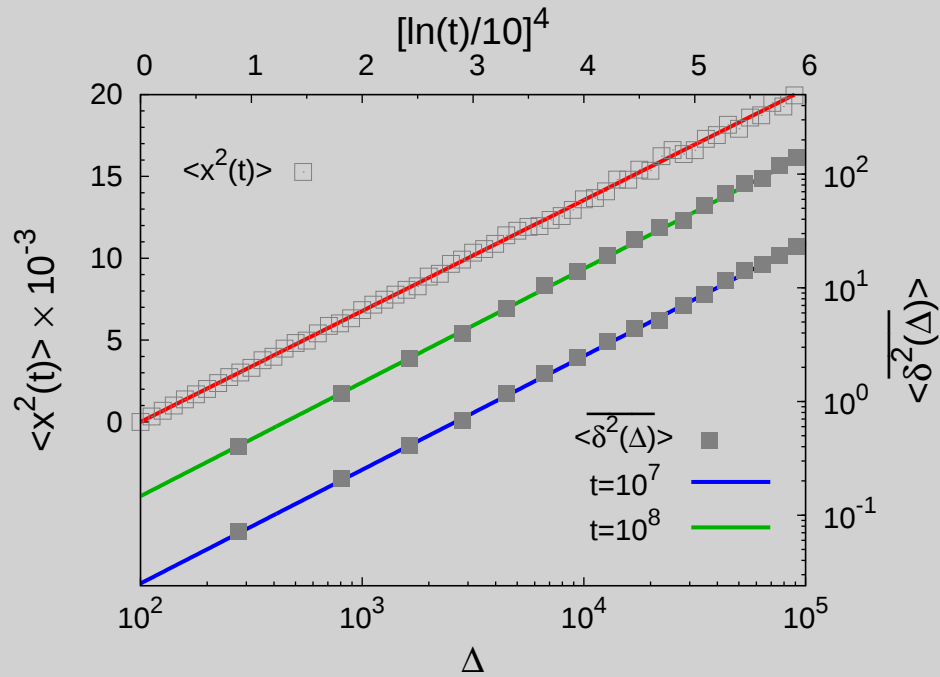
$$\langle x^2(t) \rangle \simeq \log^{1/2} \left(\frac{t}{t_0} \right)$$

WEB: $\overline{\delta^2(\Delta)} \simeq (\Delta/T) \log^{1/2} T$



[LP Sanders,] MA Lomholt, L Lizana, [K Fogelmark,] RM & T Ambjörnsson, PRL (2013); NJP (2014)

Logarithmic evolution: Sinai diffusion & mean field CTRW

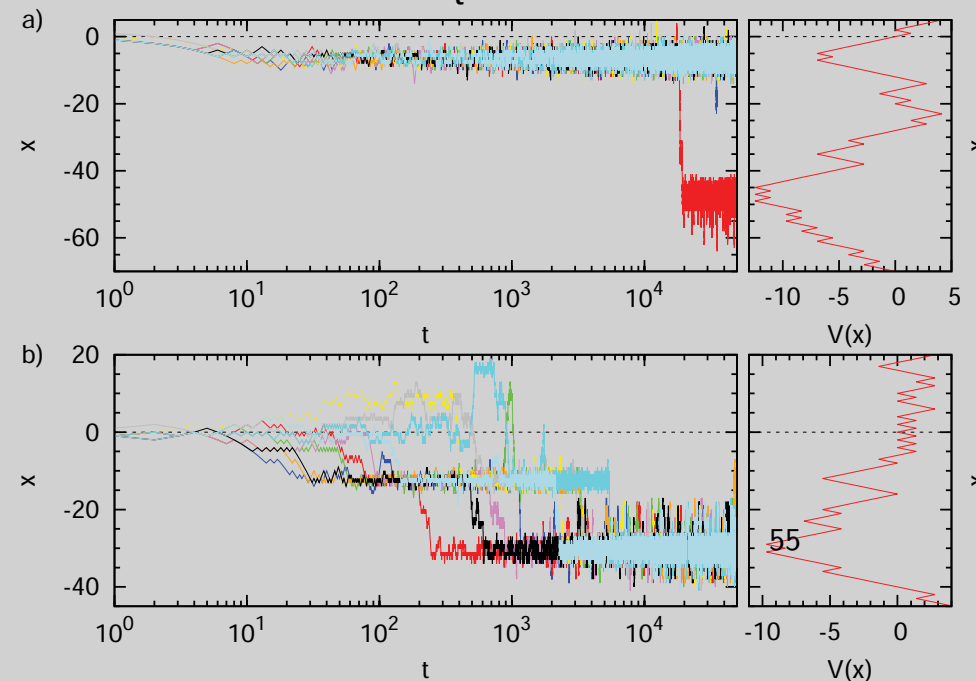


$$\langle \widetilde{x^2(t)} \rangle \simeq \log^4(t), \quad \langle \widetilde{\delta^2(\Delta)} \rangle \simeq \log^4(T) \frac{\Delta}{T}$$

Mean field approach (comp ultraslow maps):

$$\psi(t) \simeq \frac{1}{t \log^4 t}$$

AV Chechkin, H Kantz, A Godec, RM & E Barkai, JPA (2014)



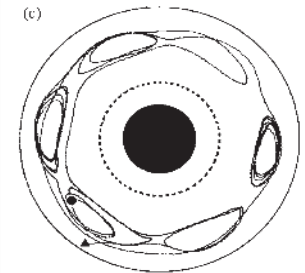
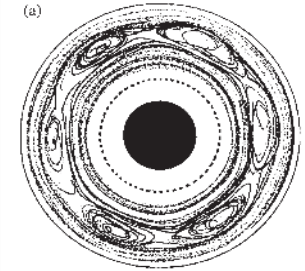
Ultraweak ergodicity breaking: chaos & marine predators

For superdiffusive Lévy walk with $\langle x^2(t) \rangle \simeq t^{3-\alpha}$ and $1 < \alpha < 2$:

$$\langle \overline{\delta^2(\Delta)} \rangle_T = \langle \overline{\delta x^2} \rangle_\infty + T^{3-\alpha} \left[\frac{\alpha-1}{3} \left(\frac{\Delta}{T} \right)^3 - \alpha \left(\frac{\Delta}{T} \right)^2 \right]$$

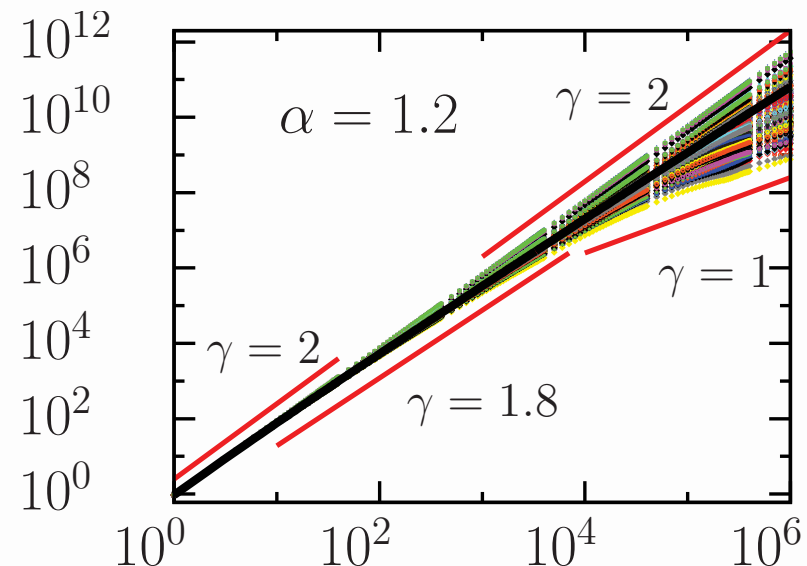
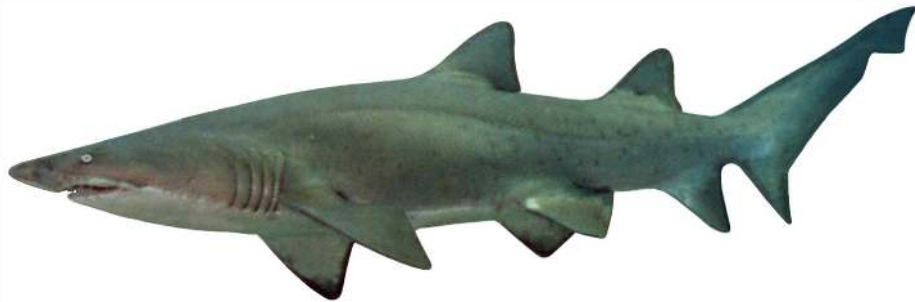
and

$$\langle \overline{\delta^2(\Delta)} \rangle_\infty = \lim_{T \rightarrow \infty} \overline{\delta^2(\Delta)} = 2 \left(\frac{(1+\Delta)^{3-\alpha} - 1}{(3-\alpha)(2-\alpha)} - \frac{\Delta}{2-\alpha} \right)$$



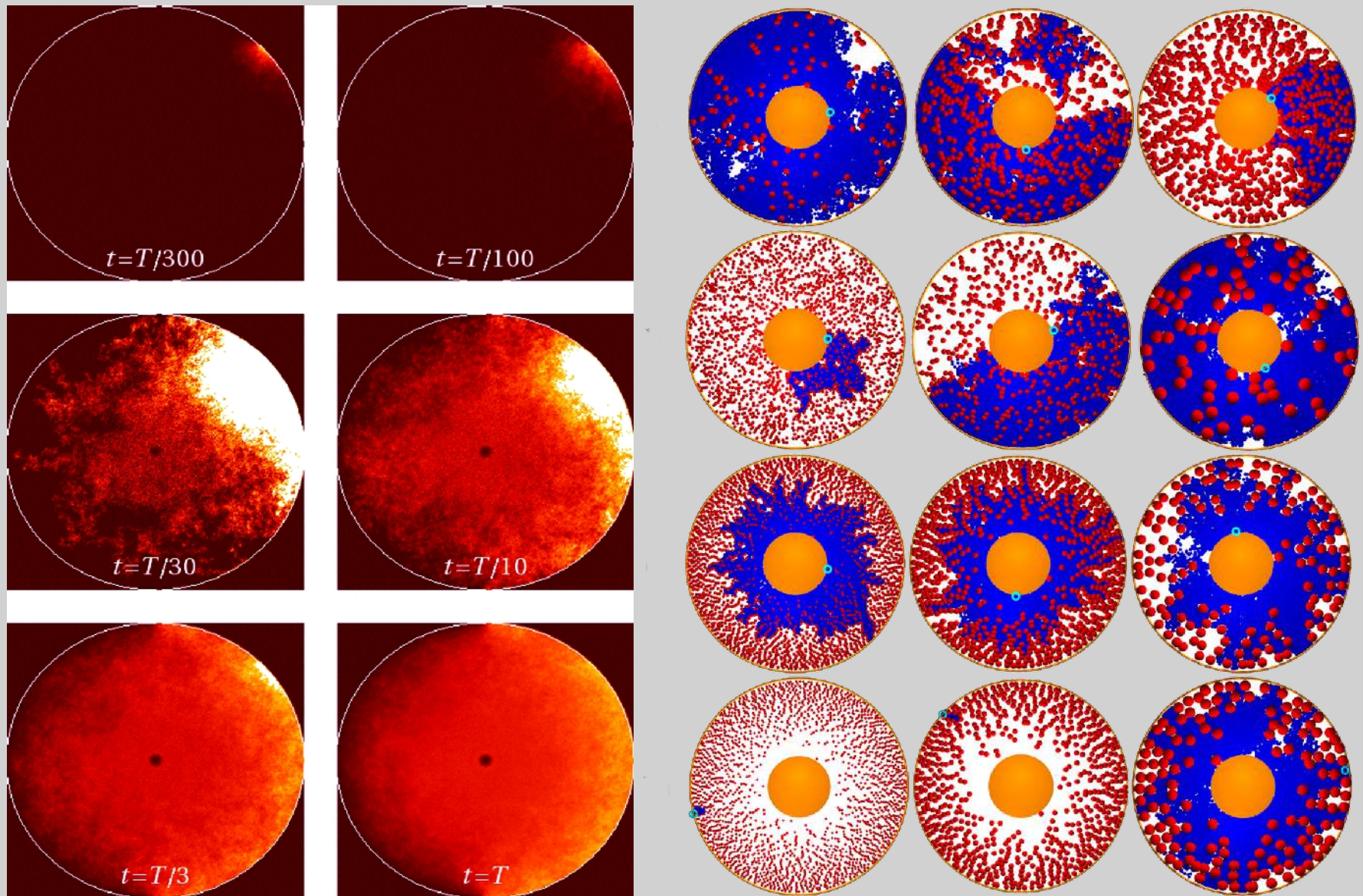
$\mathcal{EB}$ parameter ($T \rightarrow \infty$)

$$\mathcal{EB}(\tau) = \frac{\langle \overline{\delta^2(\Delta)} \rangle}{\langle x^2(\Delta) \rangle} \sim \frac{1}{\alpha - 1}$$

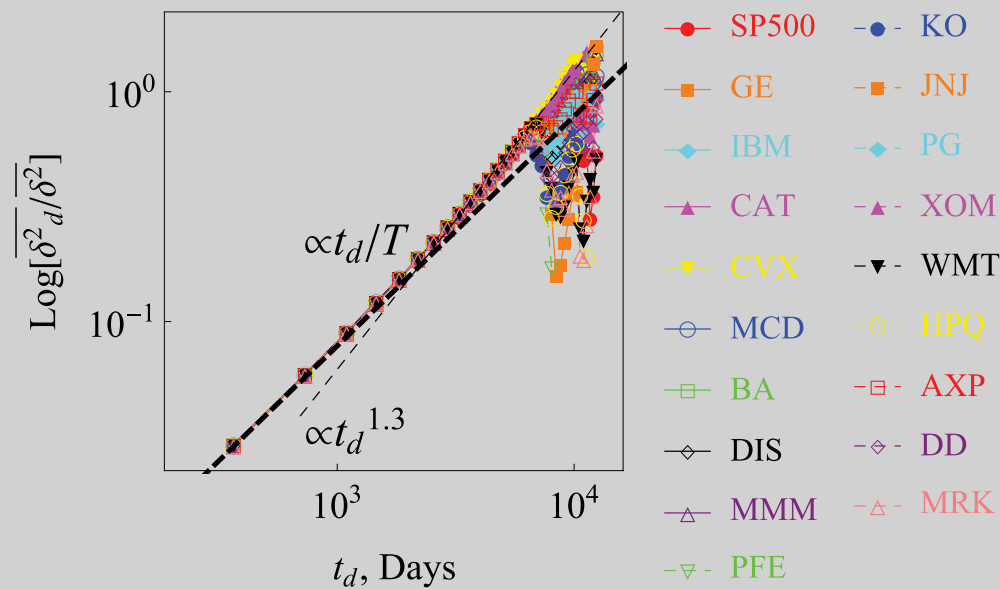
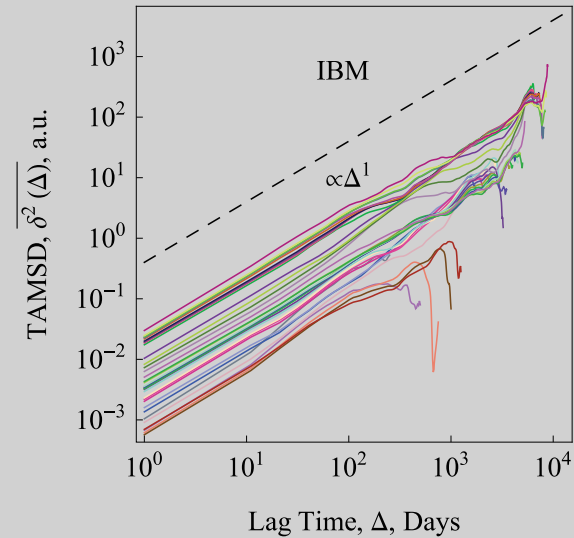
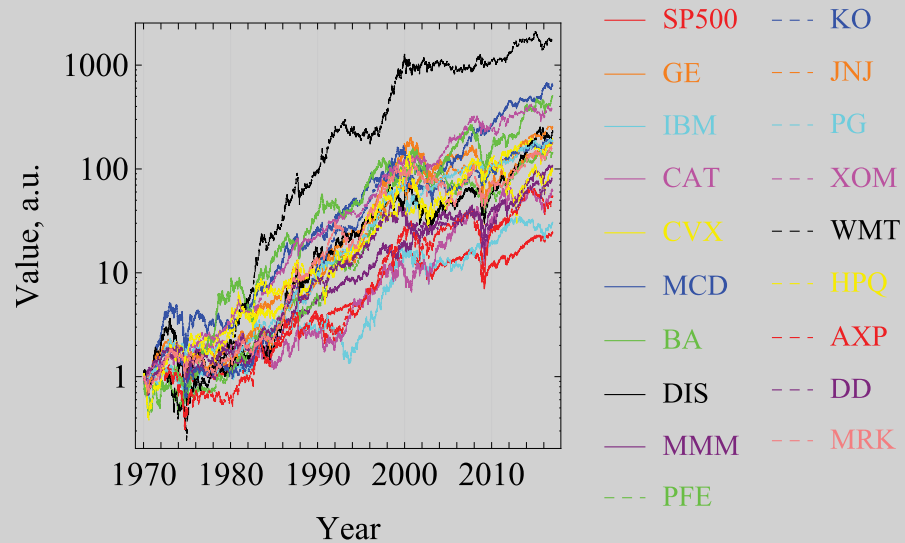


A Godec & RM, PRL (2013); D Froemberg & E Barkai, PRE (2013)

Tracer diffusion in disordered environments



Time averaged & delay time analyses of financial time series



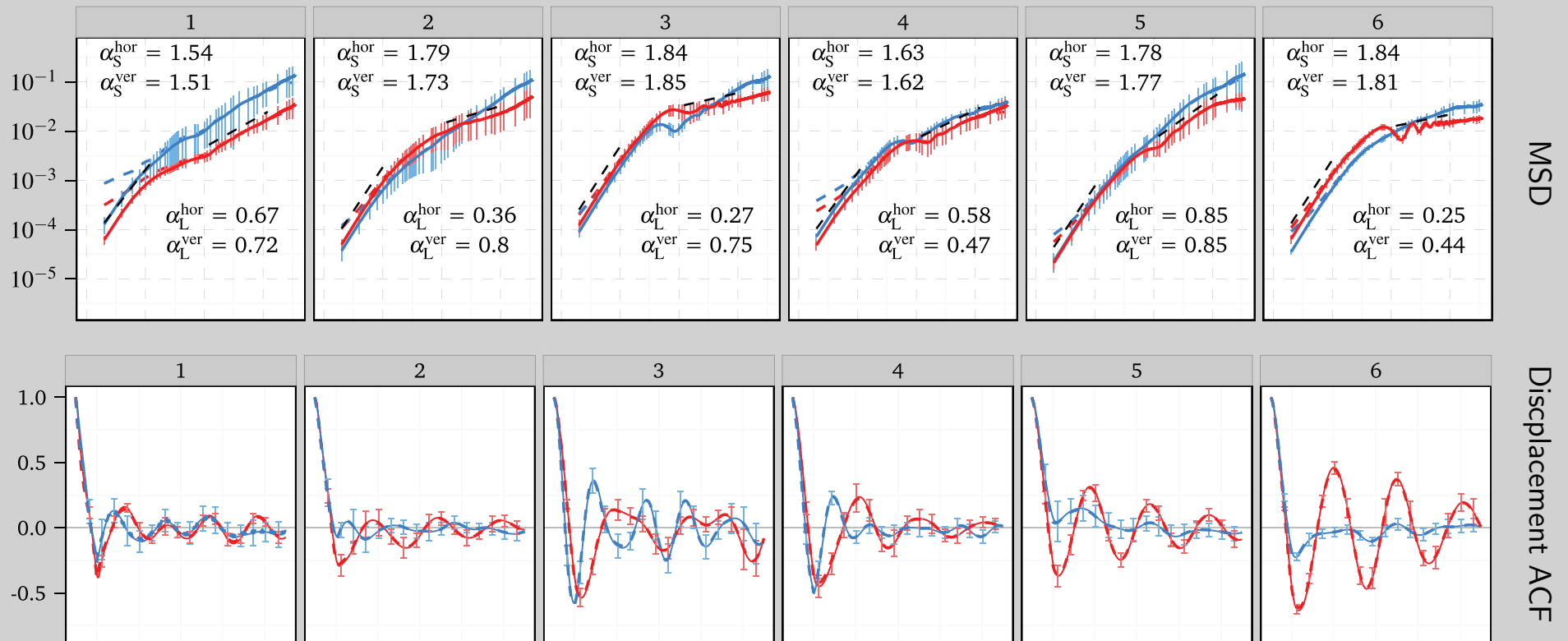
$$\text{GBM } dX(t) = \mu X(t)dt + \sigma X(t)dW(t)$$

$$\overline{\delta_d^2(\Delta)} = \frac{\int_{t_d}^{T-\Delta} [X(t+\Delta) - X(t)]^2 dt}{T - t_d - \Delta}$$

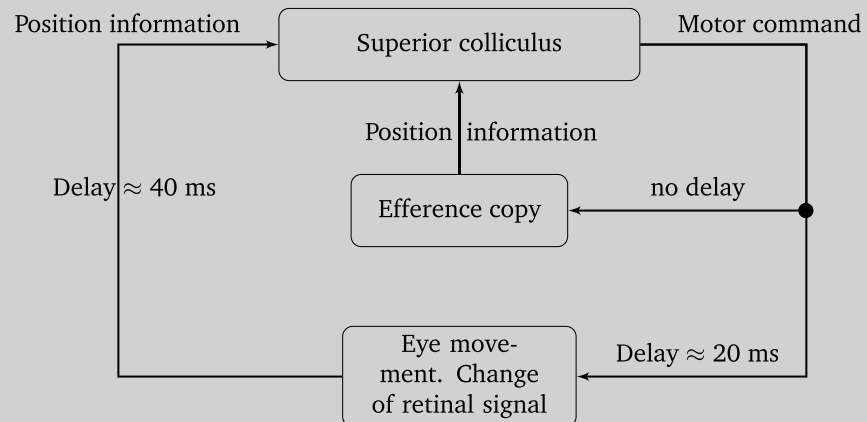
$$\sim \frac{\Delta}{T - t_d} X_0^2 \left(e^{\sigma^2 T} - e^{\sigma^2 t_d} \right)$$

$$\log \left[\frac{\langle \overline{\delta_d^2(\Delta, t_d)} \rangle}{\langle \overline{\delta^2(\Delta)} \rangle} \right] \sim t_d/T$$

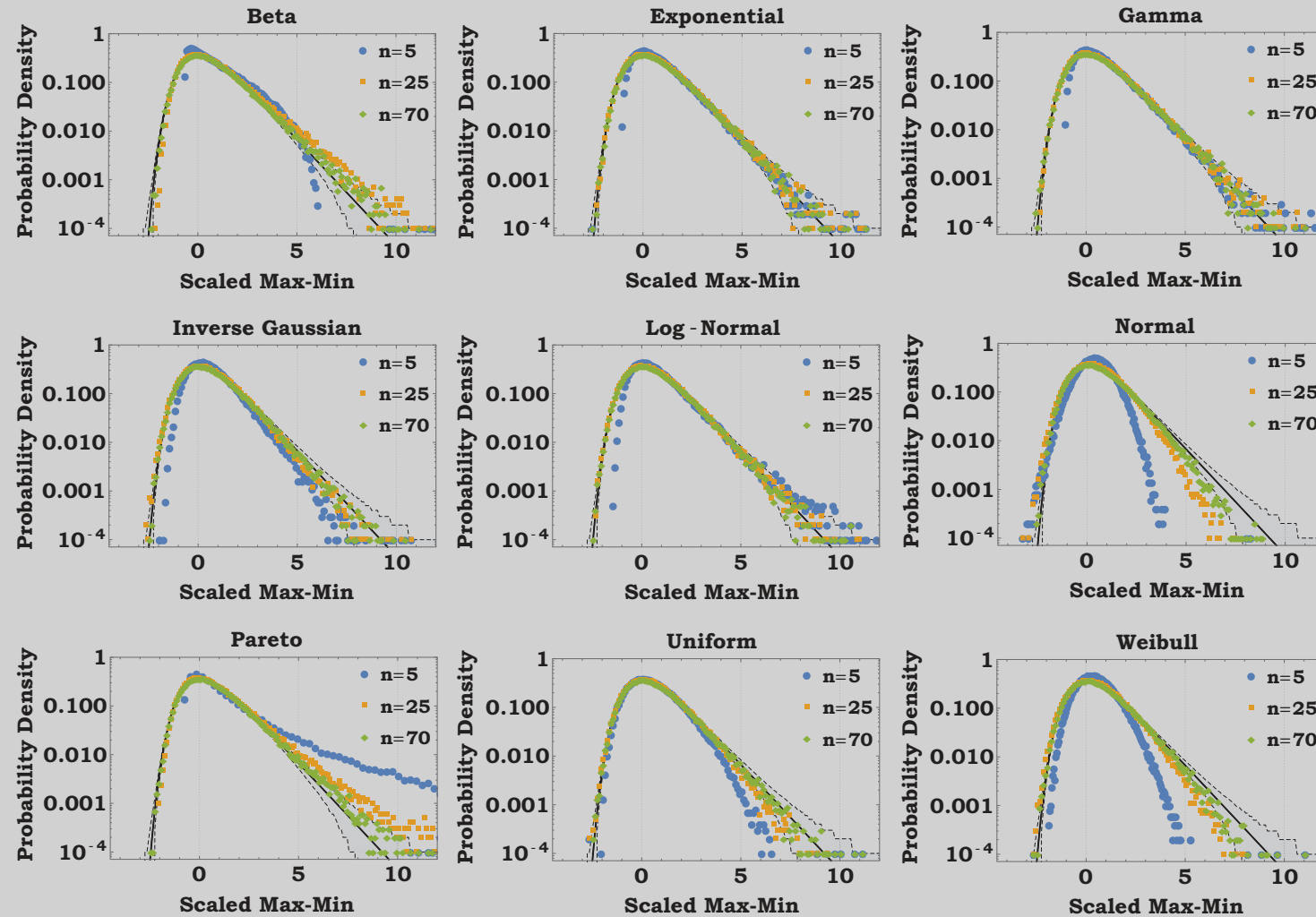
Stochasticity of fixational eye movements



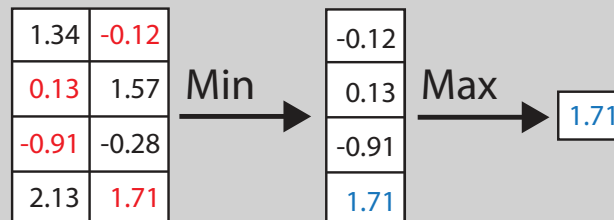
CJJ Herrmann, RM & R Engbert, Sci Rep (2017)



Gumbel universal limit law for Min-Max/Max-Min processes



I Eliazar, RM & S Reuveni, arXiv (2018)



Manifestations of WEB [RM, JH Jeon, AG Cherstvy & E Barkai, PCCP (2014)]

	$\langle x^2(t) \rangle$	$\langle \overline{\delta^2(\Delta)} \rangle$	
Correlated jump lengths	$\simeq t^3$	$\simeq \Delta^2 T$	V Tejedor & RM, JPA (2010)
$K(x) = K_0 x ^\alpha$	$\simeq t^{2/(2-\alpha)}$	$\simeq \Delta T^{2/(2-\alpha)-1}$	A Cherstvy et al, NJP (2013)
$K(t) \simeq K_0 t^{\alpha-1}$	$\simeq t^\alpha$	$\simeq \Delta T^{\alpha-1}$	JH Jeon et al, PCCP (2014)
Subdiffusive CTRW	$\simeq t^\alpha$	$\simeq \Delta T^{\alpha-1}$	Y He et al, PRL (2008)
Correlated waiting times	$\simeq t^{\alpha/(1+\alpha)}$	$\simeq \Delta T^{\alpha/(1+\alpha)-1}$	V Tejedor & RM, JPA (2010)
Ultralow CTRW	$\simeq \log^\alpha(t)$	$\simeq (\Delta/T) \log^\alpha T$	A Godec et al, JPA (2014)
Sinai (quenched)	$\simeq \log^4(t)$	$\simeq (\Delta/T) \log^4 T$	
Ageing CTRW	$\simeq \log(t)$	$\simeq \Delta \log(T)$	MA Lomholt et al, PRL (2013)
$K(x) = K_0 \exp(-2x)$	$\simeq \log^2(t)$	$\simeq (\Delta/T)^{1/2}$	A Cherstvy & RM, PCCP (2013)
Lévy walk	$\simeq A(\alpha) t^{3-\alpha}$	$\simeq \frac{A(\alpha)}{\alpha-1} \Delta^{3-\alpha}$	D Froemberg & E Barkai, PRE (2013) A Godec & RM, PRL/PRE (2013)

Overview articles

- I Single particle manipulation & tracking:
C Nørregaard, RM, CM Ritter, K Berg-Sørensen & LB Oddershede,
Chem Rev **117**, 4342 (2017)

- II Anomalous diffusion models, WEB & ageing:
RM, JH Jeon, AG Cherstvy & E Barkai, Phys Chem Chem Phys **16**,
24128 (2014)

- III Ageing renewal theory:
JHP Schulz, E Barkai & RM, Phys Rev X **4**, 011028 (2014)

- IIII Anomalous diffusion in membranes:
RM, JH Jeon & AG Cherstvy, Biochimica et Biophysica Acta - Biomem-
branes **1858**, 2451 (2016)

- IIII Polymer translocation across (membrane) nanopores:
V Palyulin, T Ala-Nissila & RM, Soft Matter **10**, 9016 (2014)

